

MT1: DYNAMIC MODE DECOMPOSITIONS AND KOOPMAN ANALYSIS

INTRODUCTION TO KOOPMAN ANALYSIS

MARKO BUDIŠIĆ

WITH J. NATHAN KUTZ (U WASH), M. HEMATI (U MINN)

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SIAM DS19 SNOWBIRD, 2019



- **Koopman operator:**
exact linear representation of (nonlinear) dynamics
- Koopman **eigenfunctions:**
generalize Lyapunov functions, isochrons,...
- Koopman **modes:**
in spirit, analogous to normal modes for linear PDEs
- data-driven (model-free) calculation enabled by (a family of) **Dynamic Mode Decomposition (DMD)** algorithm(s) (see J. N. Kutz's talk)
- can be **applied to** reduced-order modeling, global linearization, system ID, sensitivity, control... (see M. Hemati's talk)

RELATED SESSIONS

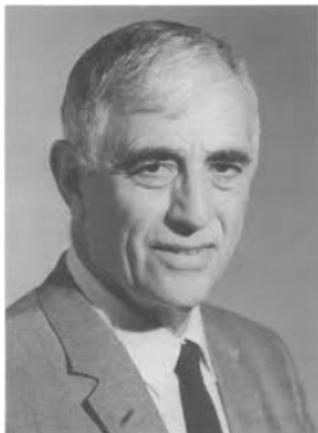
- MS48 Koopman Operator Techniques in Dynamical Systems: Theory
- MS61 Advanced Data-Driven Techniques and Numerical Methods in Koopman Operator Theory - Part I of II
- MS74 Advanced Data-Driven Techniques and Numerical Methods in Koopman Operator Theory - Part II of II
- MS86 Applications of Koopman Operator Theory in Dynamical Systems: From Fluids, through Machine Learning to Energy - Part I of II
- CP10 Data and Koopman Analysis
- MS97 Applications of Koopman Operator Theory in Dynamical Systems: From Fluids, through Machine Learning to Energy - Part II of II
- MS147 Control Techniques based on Koopman Operator Theory - Part I of II
- MS160 Control Techniques based on Koopman Operator Theory - Part II of II
- MS164 Theory and Application of Koopman Operator Methods in Molecular Simulation
- CP36 Koopman Analysis
- ... and a whole bunch of sessions on reduced-order models, data-driven nonlinear analysis, etc.

- 1 History and Present
- 2 Introduction
- 3 Koopman eigenfunctions
- 4 Koopman modes

Goals of this talk:

- explain the basics
- give some intuition
- point to active research areas

HISTORY AND PRESENT



Bernard Koopman

Vol. 17, 1931 *MATHEMATICS: B. O. KOOPMAN* 315
*HAMILTONIAN SYSTEMS AND TRANSFORMATIONS IN
HILBERT SPACE*
By B. O. KOOPMAN
DEPARTMENT OF MATHEMATICS, COLUMBIA UNIVERSITY



George Birkhoff

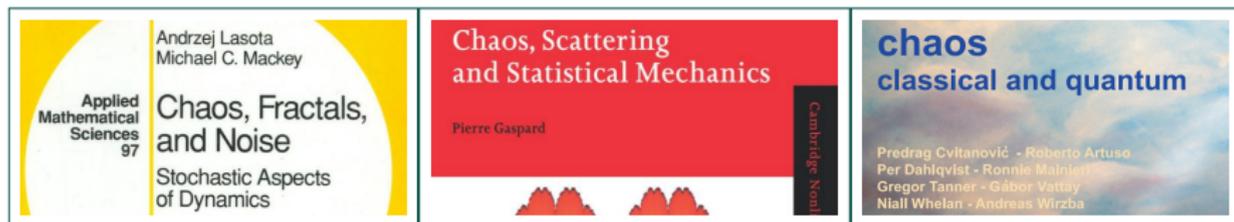
Vol. 18, 1932 *MATHEMATICS: BIRKHOFF AND KOOPMAN* 279
RECENT CONTRIBUTIONS TO THE ERGODIC THEORY
By G. D. BIRKHOFF AND B. O. KOOPMAN
DEPARTMENT OF MATHEMATICS, HARVARD UNIVERSITY



John von Neumann

*ZUR OPERATORENMETHODE IN DER KLASSISCHEN
MECHANIK*¹.
VON J. V. NEUMANN, PRINCETON.

Although Koopmanism (then: *Koopmania*) was present in 90s monographs...



It's the development of data-driven algorithms...

<p>Nonlinear Dynamics (2005) 41: 309–325 © Springer 2005</p> <p>Spectral Properties of Dynamical Systems, Model Reduction and Decompositions</p> <p>IGOR MEZIĆ Department of Mechanical and Environmental Engineering and Department of Mathematics, University of California, Santa Barbara, CA 93106-5070, U.S.A. (e-mail: mezi@engineering.ucsb.edu; fax: +1-805-893-8551)</p> <p>(Received: 25 May 2004; accepted: 15 June 2004)</p>	<p><i>J. Fluid Mech.</i> (2010), vol. 656, pp. 5–28. © Cambridge University Press 2010 doi:10.1017/S0022112010000217</p> <p>Dynamic mode decomposition of numerical and experimental data</p> <p>PETER J. SCHMID† Laboratoire d'Hydrodynamique (LadHyX), CNRS-École Polytechnique, 91128 Palaiseau, France</p> <p>(Received 20 May 2009; revised 7 March 2010; accepted 7 March 2010; first published online 1 July 2010)</p>
<p><i>J. Fluid Mech.</i> (2009), vol. 641, pp. 115–127. © Cambridge University Press 2009 doi:10.1017/S0022112009992059</p> <p>Spectral analysis of nonlinear flows</p> <p>CLARENCE W. ROWLEY¹†, IGOR MEZIĆ², SHERVIN BAGHERI¹, PHILIPP SCHLATTER³ AND DAN S. HENNINGSON¹</p> <p>¹Department of Mechanical and Aerospace Engineering, Princeton University, NJ 08544, USA ²Department of Mechanical Engineering, University of California, Santa Barbara, CA 93106-5070, USA ³Linné Flow Centre, Department of Mechanics, Royal Institute of Technology (KTH), SE-10044 Stockholm, Sweden</p> <p>(Received 15 May 2009; revised 8 September 2009; accepted 9 September 2009; first published online 18 November 2009)</p>	

... that brought us here.

INTRODUCTION

Dynamics of states

Linear system $\underline{z}(t) = \begin{bmatrix} x(t) \\ y(t) \end{bmatrix}$

$$\dot{\underline{z}}(t) = \mathbf{A}\underline{z}(t), \underline{z}(0) = \underline{z}_0 \quad (\text{ODE})$$

$$\underline{z}(t) = \underbrace{\exp(\mathbf{A}t)}_{\Phi^t(\underline{z}_0)} \underline{z}_0 \quad (\text{Flow map})$$

Dynamics of measurements

Measurement (*observable*)

$h : \mathbb{R}^2 \rightarrow \mathbb{R}$, evolves along a trajectory according to:

$$h_t(\underline{z}_0) = h(\underline{z}(t)) = h(\Phi^t(\underline{z}_0))$$

KOOPMAN OPERATOR EVOLVES MEASUREMENTS.

General nonlinear systems:

$$\dot{\underline{z}}(t) = f(\underline{z}(t)), \underline{z}(0) = \underline{z}_0 \quad (\text{ODE})$$

$$\underline{z}(t) = \Phi^t(\underline{z}_0) \quad (\text{Flow map})$$

Measurement evolution:

Koopman operator

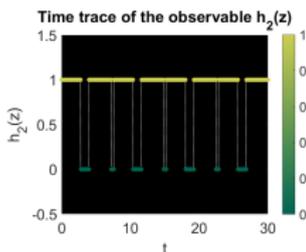
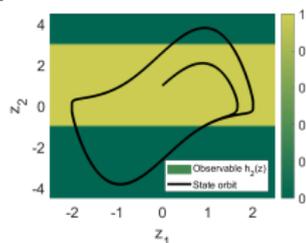
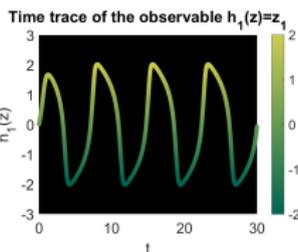
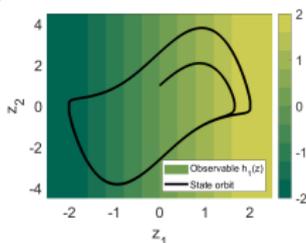
$$\mathbb{K}^t : \text{Fun} \rightarrow \text{Fun}$$

$$[\mathbb{K}^t h](\underline{z}) = h(\Phi^t(\underline{z}))$$

$$\mathbb{K}^t h = h \circ \Phi^t$$

a.k.a. pull-back by evolution/flow

a.k.a. composition operator



Recipe:

- Seed a grid of initial conditions z_k
- Compute a trajectory from each point
 $z_k \rightsquigarrow z_k(t)$
- Evaluate the (scalar) function
 $f(z_k(t)) =: \mathbb{K}^t f(z_k)$ at final point
- Plot color field $\mathbb{R}^2 \mapsto \mathbb{R}$

$$z_k \mapsto [\mathbb{K}^t f](z_k)$$

KOOPMAN OPERATOR IS LINEAR.

Koopman operator $\mathbb{K} : \text{Fun} \rightarrow \text{Fun}$ is
linear by construction

$$\underbrace{(\alpha f + \beta g) \circ \Phi^t}_{\mathbb{K}^t(\alpha f + \beta g)} = \alpha \underbrace{f \circ \Phi^t}_{\mathbb{K}^t f} + \beta \underbrace{g \circ \Phi^t}_{\mathbb{K}^t g}$$

No magic: this is **not linearity in state** variables

$$\mathbb{K} f(\alpha z + \beta w) \neq \mathbb{K} f(\alpha z) + \mathbb{K} f(\beta w)$$

Trade-off

Flow map Φ^t	Koopman o. \mathbb{K}^t
Non-linear	Linear
Finite dim.	∞ -dim.

No trade off if Φ^t is ∞ -dimensional itself!

- ergodic dynamics + L_2 space of observables = \mathbb{K} is unitary
- in this case, Koopman op. is adjoint to the Perron–Frobenius transfer operator

Spectral Decomposition of the Koopman Operator

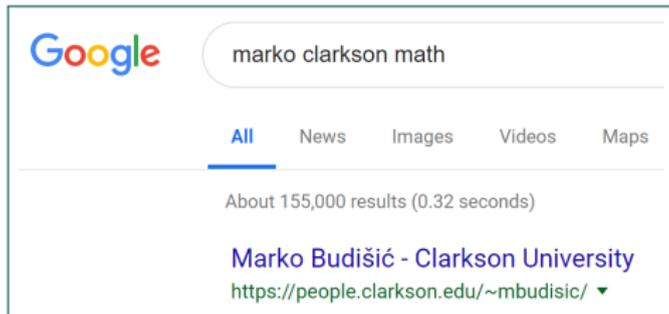
$$\mathbb{K}^n f = \int_{-\pi}^{\pi} e^{in\omega} d[\mathbb{E}(\omega)f] = \underbrace{\sum_k e^{in\omega_k} \mathbb{P}_k f}_{\text{atomic}} + \underbrace{\int_{-\pi}^{\pi} e^{in\omega} d[\mathbb{E}_c(\omega)f]}_{\text{continuous}}.$$

- **atomic spectrum** \rightarrow **eigenvalues** \rightarrow **(quasi)regular dynamics**
- a.c. spectrum \rightarrow density function \rightarrow mixing dynamics
- s.c. spectrum \rightarrow fractal \rightarrow “anomalous transport”

WHAT IS APPLIED KOOPMAN ANALYSIS?

- approximation of the Koopman operator from data
- computational spectral analysis
- interpretation and application of results

The PDF of this talk online (soon) on my website:



KOOPMAN EIGENFUNCTIONS

EIGENFUNCTIONS FOR LINEAR DYNAMICS (SADDLE)

$$\dot{z}(t) = \mathbf{A}z(t), z(0) = z_0$$

$$z(t) = \exp(\mathbf{A}t)z_0$$

Given left eigenvector at $\lambda \in \mathbb{R}$:

$$v^* \mathbf{A} = \lambda v^*, \quad v^* \exp(\mathbf{A}t) = e^{\lambda t} v^*$$

Functions

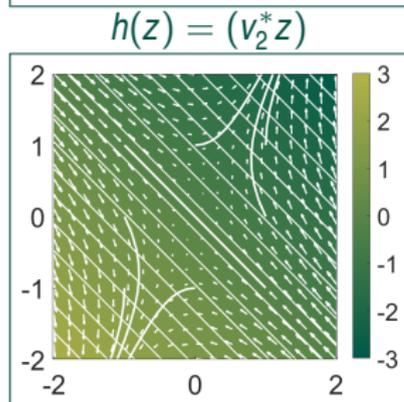
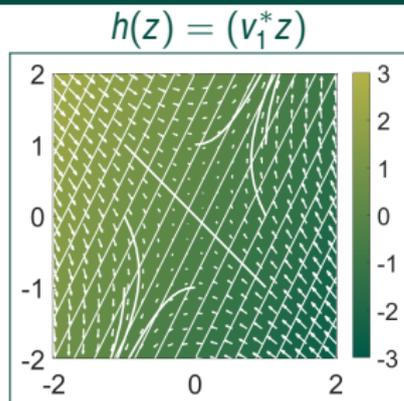
$$h(z) = v^* z$$

are **eigenfunctions** at $e^{\sigma t}$

$$h(z) = (v^* z)$$

$$\mathbb{K}^t h(z) = v^* e^{\mathbf{A}t} z = e^{\lambda t} \underbrace{v^* z}_{h(z)}$$

Structure of level sets stays constant in time. Values grow/decay according to eigenvalues.



EIGENFUNCTIONS FOR LINEAR DYNAMICS (FOCUS)

$$\dot{z}(t) = \mathbf{A}z(t), z(0) = z_0$$

$$z(t) = \exp(\mathbf{A}t)z_0$$

Given left e.-vectors $v_+, v_- = \bar{v}$ at $\sigma \pm i\omega$:

$$v_{\pm}^* \mathbf{A} = \lambda v_{\pm}^*, \quad v_{\pm}^* \exp(\mathbf{A}t) = e^{\sigma t \pm i\omega t} v_{\pm}^*$$

Functions

$$h(z) = (v^* z)(\bar{v}^* z)$$

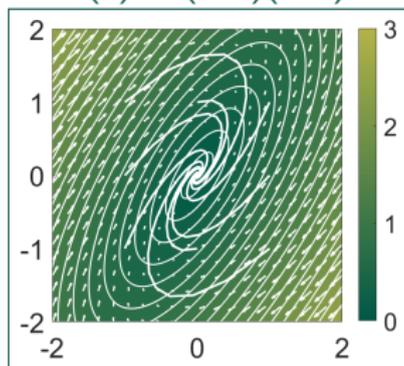
are **eigenfunctions** at $e^{\sigma t}$

$$h(z) = (v^* z)(\bar{v}^* z) = |v^* z|^2$$

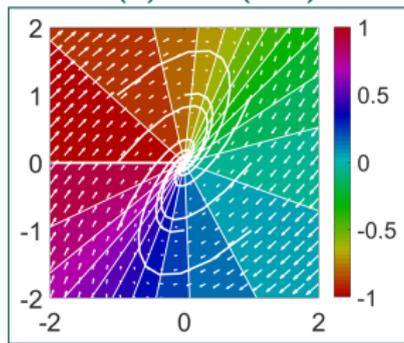
$$\mathbb{K}^t h(z) = (v^* e^{\mathbf{A}t} z)(\bar{v}^* e^{\mathbf{A}t} z) = e^{\sigma t} \underbrace{(v^* z)(\bar{v}^* z)}_{h(z)}$$

Level sets correspond to a **Lyapunov function** $|v^* z|^2$ and **isochrons** $\angle(v^* z)$.

$$h(z) = (v^* z)(\bar{v}^* z)$$



$$h(z) = \angle(v^* z)$$



HOW TO CALCULATE EIGENFUNCTIONS

Trajectory averages $z_{n+1} = \Phi(z_n)$ project **any** observable onto an eigenfunction:

$$h_\omega(z_0) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \underbrace{\lambda^n}_{\text{E-value.}} \overbrace{h(z_n)}^{[\mathbb{K}^n h](z_0)}$$

- **Ergodic average** – $\lambda = 0$ – invariant functions
- **Harmonic average** – $\lambda = e^{i\omega}$ – (quasi)periodic functions

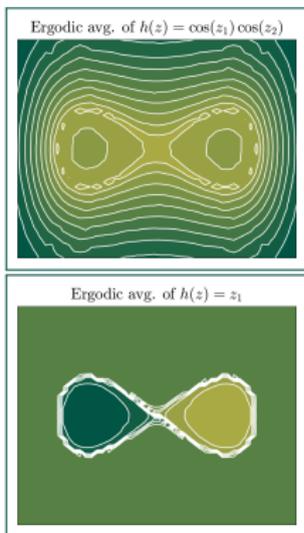
Recipe for eigenfunctions

1. Seed a grid of initial conditions z_k .
2. Simulate (long) trajectories $z_k \rightsquigarrow z_k(t)$.
3. Choose an observable and a frequency ω .
4. Compute the harmonic average and visualize $z_k \mapsto h_\omega(z_k)$.
(If you choose a non-eigenvalue $e^{i\omega}$ you'll get $h_\omega \equiv 0$.)

EIGENFUNCTIONS FOR NONLINEAR DYNAMICS

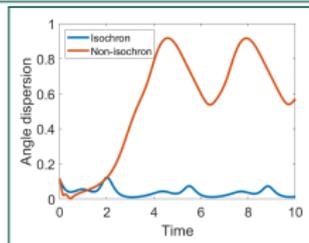
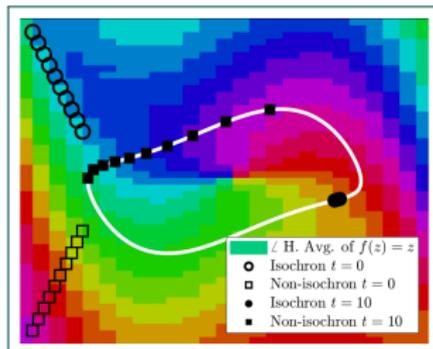
Ergodic Average (double-well oscillator)

- Ergodic averages of observables are conserved quantities.
- Level sets of ergodic averages are invariant sets.



Harmonic Average (van der Pol oscillator)

- Angle of harmonic average of $h(z) = z$ at frequency of limit cycle.
- Level curves are isochrons.



ACTIVE RESEARCH TOPICS

- spectral theory for transient (non-steady-state) dynamical systems
- global stability and global linearization based on Koopman eigenfunctions
- control and system identification based on Koopman spectral analysis
- numerical methods for approximations and analysis of Koopman operator
- investigations into the spectral measure and non-atomic Koopman spectrum
- rigorous extensions of the Koopman theory for PDEs and SPDEs
- Koopman theory in reproducing kernel Hilbert spaces

IPAM Operator Theoretic Methods in Dynamic Data Analysis and Control, Feb 2019, (link with videos)

Speakers: Nelida Črnjarić-Žic (University of Rijeka) Zlatko Drmač (University of Zagreb) Maria Fonoberova (AIMdyn) Gary Froyland (University of New South Wales) Dimitris Giannakis (New York University, Courant Institute of Mathematical Sciences) Didier Henrion (Centre National de la Recherche Scientifique (CNRS), Laboratoire d'Analyse et d'Architecture des Systemes (LAAS)) Maria Infusino (Universität Konstanz) Oliver Junge (Technical University of Munich) Milan Korda (Centre National de la Recherche Scientifique (CNRS)) J. Nathan Kutz (University of Washington, Applied Mathematics) Jean Lasserre (Université de Toulouse III (Paul Sabatier), LAAS-CNRS) Yuri Latushkin (University of Missouri-Columbia) Senka Maćešić (University of Rijeka) Krithika Manohar (California Institute of Technology, Computing and Mathematical Sciences) Alexandre Mauroy (Université de Namur) Igor Mezic (University of California, Santa Barbara (UCSB), Mechanical Engineering) Ryan Mohr (University of California, Santa Barbara (UCSB)) Nader Motee (Lehigh University, Mechanical Engineering and Mechanics) Hiroya Nakao (Tokyo Institute of Technology) Frank Noe (Freie Universität Berlin) Mihai Putinar (University of California, Santa Barbara (UCSB), Mathematics) Peter Schmid (Imperial College, Mathematics) Amit Surana (United Technologies Research Center) Umesh Vaidya (Iowa State University, Mechanical Engineering) Irène Waldspurger (Université de Paris IX (Paris-Dauphine)) Tillmann Weisser (Los Alamos National Laboratory) Enoch Yeung (University of California, Santa Barbara (UCSB))

KOOPMAN MODES

LINEAR PDEs: NORMAL MODE ANALYSIS

- **States:** Simulated 16 linear oscillators
- **Observables:** Displacement between them (polynomial interpolation)
- Observables are **indexed** by x – there is a continuum of them.
 - Interpretation: discretized wave equation.
 - **Normal modes:** x –spatial profiles oscillating at isolated frequencies (standing waves).
 - **Modes do not depend on state representation** (could have used spectral instead of FD solver...)

MANY OBSERVABLES THROUGH LENS OF 1 EIGENFUNCTION

If $f \in \text{span}(\text{eigenfunctions})$ we can **decompose it into eigenfunctions**.

$$f(z) = \sum_k \underbrace{m_k}_{\text{Coefficients}} \underbrace{h_k(z)}_{\text{E-fun.}} \quad (\text{Observable})$$

$$\mathbb{K}^t f(z) = \sum_k \underbrace{m_k}_{\text{Coefficients}} \underbrace{\lambda_k^t}_{\text{E-val.}} h_k(z) \quad (\text{Evolution})$$

Now, consider *many* observables $f_x(z)$ for $x \in \mathcal{S}$

$$\mathbb{K}^t f_x(z) = \sum_k \underbrace{m_k(x)}_{\text{Koopman mode}} \lambda_k^t h_k(z)$$

Koopman mode $m_k(x)$

Demonstrates importance of eigenvalue λ_k across a measurements indexed by x .

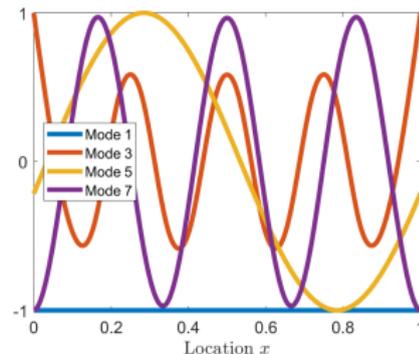
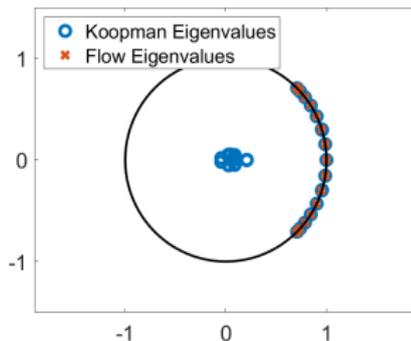
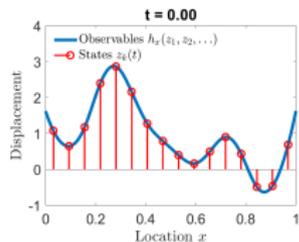
Normal mode analysis

- model-dependent
- analytic
- works for select (non)linear systems

Koopman mode analysis

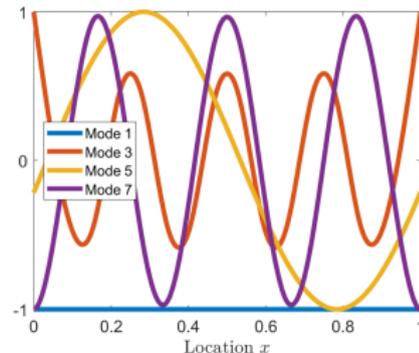
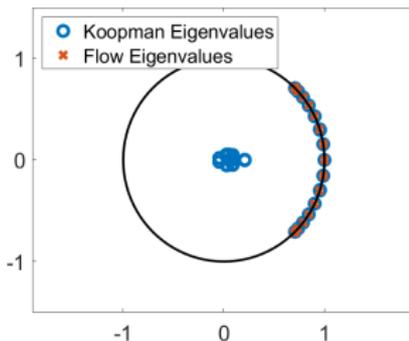
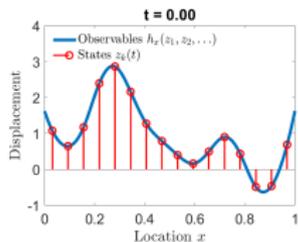
- model-free (data-driven)
- computational
- works for all (non)linear systems (they are after all just \mathbb{K})

KOOPMAN MODES FOR LINEAR VIBRATIONS



- Based on a single trajectory of dynamics
- Eigenvalues match the linear analysis
- Modes correctly capture the expected standing waves
- ...but don't they look a bit funny at ends?
- Different trajectory (initial condition) could excite different modes

KOOPMAN MODES FOR LINEAR VIBRATIONS



Based on a single trajectory of dynamics

Computation (Dynamic Mode Decomposition)

Simulation data \Rightarrow SVD \Rightarrow EIG \Rightarrow done.

... (much) more on this in the coming parts.

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KOOPMAN MODE ANALYSIS EXTENDS NORMAL MODES TO NONLINEAR DYNAMICS.

J. Fluid Mech. (2009), vol. 641, pp. 115–127. © Cambridge University Press 2009
doi:10.1017/S0022112009992059

Spectral analysis of nonlinear flows

CLARENCE W. ROWLEY¹†, IGOR MEZIĆ²,
SHERVIN BAGHERI³, PHILIPP SCHLATTER³
AND DAN S. HENNINGSON³

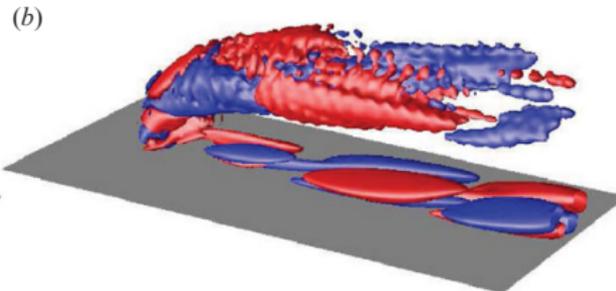
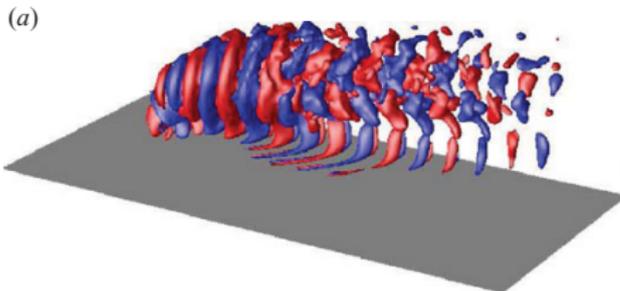
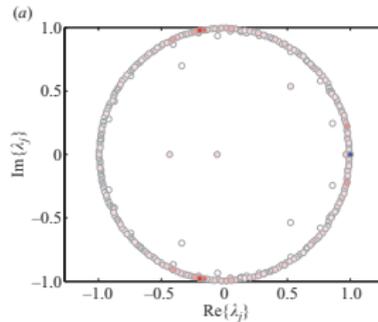
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18 November 2009)

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WHERE TO START READING (A NON-EXHAUSTIVE LIST)

Papers:



Applied Koopmanism

Marko Budišić, Ryan Mohr, and Igor Mezić

Citation: *Chaos: An Interdisciplinary Journal of Nonlinear Science* **22**, 047510 (2012); doi: 10.1063/1.4772195

Annu. Rev. Fluid Mech. 2013, 45:317–74
First published online as a Review in Advance on October 5, 2012

The Annual Review of Fluid Mechanics is online at fluid.mech.berkeley.edu

This article's doi:
10.1146/annurev-fluid-011212-140052

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Analysis of Fluid Flows via Spectral Properties of the Koopman Operator

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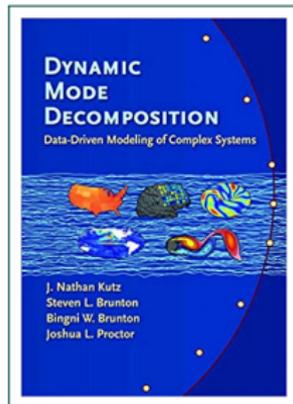
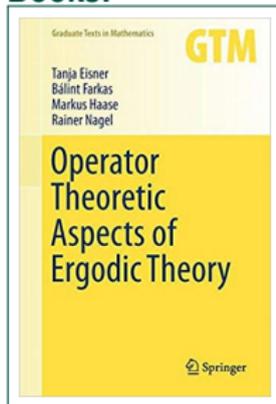
J Nonlinear Sci (2012) 22:887–915
DOI 10.1007/s00332-012-9130-9

Journal of
**Nonlinear
Science**

Variants of Dynamic Mode Decomposition: Boundary Condition, Koopman, and Fourier Analyses

Kevin K. Chen · Jonathan H. Tu ·
Clarence W. Rowley

Books:



Upcoming books:

- I. Mezić on spectral analysis of dynamical systems
- S. Brunton et al. on data-driven methods in dynamics

- **Koopman operator:**
exact linear representation of (nonlinear) dynamics
- Koopman **eigenfunctions:**
generalize Lyapunov functions, isochrons,...
- Koopman **modes:**
in spirit, analogous to normal modes for linear PDEs
- data-driven (model-free) calculation enabled by (a family of) **Dynamic Mode Decomposition (DMD)** algorithm(s) (see J. N. Kutz's talk)
- can be **applied to** reduced-order modeling, global linearization, system ID, sensitivity, control... (see M. Hemati's talk)