

# Periodic Orbits to Gross Pitaevskii with Vortices following Point Vortex Flow

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# The Gross-Pitaevskii Equations

Seek non-constant time periodic solutions to the Gross-Pitaevskii (GP) equations

$$iu_t(x, t) = \Delta u(x, t) + \frac{u(x, t)(1 - |u(x, t)|^2)}{\varepsilon^2}, \quad (x, t) \in \mathbb{D} \times \mathbb{R},$$

posed on the unit disc  $\mathbb{D}$ , subject to the Dirichlet boundary conditions (BC)

$$u(e^{i\theta}, t) = g_n(\theta) := e^{in\theta}, \quad \theta \in [0, 2\pi), t \in \mathbb{R}.$$

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$$\deg(g_n, \partial\mathbb{D}, 0) = n.$$

# Hamiltonian Structure

The flow (GP-BC) conserves the *Ginzburg-Landau Energy*,

$$E_\varepsilon(u) := \frac{1}{2} \int_{\mathbb{D}} |\nabla u|^2 + \frac{1}{2\varepsilon^2} (1 - |u|^2)^2 dx. \quad (0.1)$$

Here  $0 < \varepsilon \ll 1$ . Energetically, minimizers  $u_\varepsilon$  of GL prefer  $|u_\varepsilon| \approx 1$ . In the limit  $\varepsilon \rightarrow 0$ , topological restrictions from the boundary condition force *vortices*.

## Renormalized Energy *à la* Bethuel-Brezis-Helein

- The sequence of minimizers  $u_\varepsilon$  converge as  $\varepsilon \rightarrow 0$  to a *nice* function  $u_*$ , away from exactly  $n$  distinct points— **vortices**.

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- $u_*$  satisfies the Harmonic map PDE away from these vortices and has degree  $+1$  about each of these. Vortices are located at a global minimizer of the *re-normalized energy*  $W$ .
- More generally, for any positive number  $N \geq n$ , integers  $d_i, i = 1, \dots, N$  satisfying  $\sum d_i = n$ , and distinct points  $a_i, i = 1, \dots, N$ , and a boundary condition  $g$  taking values in  $\mathbb{S}^1$  with  $\deg(g, \partial\mathbb{D}, 0) = n$

$$W(a_1, \dots, a_N; d_1, d_2, \dots, d_N; g) := -\pi \sum_{i \neq j} d_i d_j \log |a_i - a_j|$$

+ boundary terms.

## Point Vortex Flow

The Hamiltonian system on  $\mathbb{C}^N$  associated to  $W$  :

$$d_j \left( \frac{da_j}{dt} \right)^\perp = -\frac{1}{\pi} \nabla_{a_j} W, \quad j = 1, \dots, N. \quad (PVF)$$

Arises in fluid mechanics as a singular limit of 2D incompressible Euler, (cf. Marchioro and Pulvirenti /Saffmann for more on this connection. )



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Crucial to us: (PVF) captures *effective dynamics* of vortices to GP as  $\varepsilon \rightarrow 0^+$ , *up to first collision time*.

Made rigorous by Colliander-Jerrard/Lin-Xin/Jerrard-Spirn.

Rigorous results on the hydrodynamic/mean field limit of GP:

Jerrard-Spirn/Serfaty.

## Main Question

Given a time periodic solution to (PVF), can we construct time-periodic solution to (GP), whose vortices follow the given periodic solution?

- Large time behavior for GP for  $\varepsilon > 0$  : given solutions to (PVF) with vortices that never collide, can we construct solutions to (GP) that *follow* these point vortices for all time, as  $\varepsilon \rightarrow 0^+$ ? Especially interesting when vortices of opposite degrees persist.

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Our (modest) contribution: in the very special case of (GP-BC), using variational and symmetry arguments, we show that for a very large class of time-periodic solutions, called relative equilibria to (PVF), there exist time-periodic solutions to (GP) following them.

# Relative Equilibria

- Definition: Uniformly rotating periodic solutions to the system (PVF).
- Obtained by pursuing the ansatz  $a_j(t) = a_j e^{i(-\tilde{\omega}t)}$ , where  $\tilde{\omega} \in \mathbb{R}$ .
- Results in nested rings of vortices, each with equal numbers of rings, and all vortices of a ring having the same degree.
- Different rings may be *aligned* or *staggered*.

# Relative Equilibria

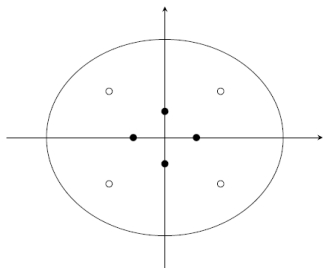


Figure 2: An staggered configuration. The solid and hollow bullets indicate possibly different degrees.  $k = 4$ .  
Not to scale.

Figure: A staggered configuration

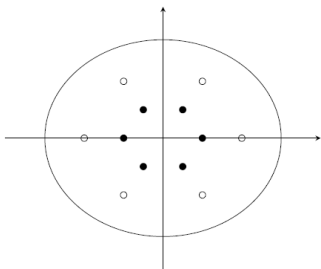


Figure 1: An aligned configuration. The solid and hollow bullets indicate possibly different degrees.  $k = 6$ . Not to scale.

Figure: An aligned configuration

# Rotational Ansatz

Starting Point: Make a rotational ansatz:

$$u(x, t) = R(-k\omega t)v\left(R\left(\frac{\omega}{m}t\right)x\right),$$



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$$u(x, t) = R(-k\omega t)v\left(R\left(\frac{\omega}{m}t\right)x\right),$$

Here:  $R(\beta)$  is the counterclockwise rotation matrix by an angle  $\beta$ ;  $k, m$  are integers and  $\omega \in \mathbb{R}$ .

Thanks to Bob Jerrard for suggesting this ansatz in the case  $n = 1$  of a single vortex.

# An Elliptic PDE: Variational Formulations

Plugging in the ansatz into (GP) yields an elliptic PDE.

$$\Delta v(y) + \frac{v}{\varepsilon^2}(1 - |v|^2)(y) = \omega \left( kv + \frac{1}{m}y^\perp \cdot \nabla v^\perp \right) (y), \quad y \in \mathbb{D},$$
$$v(e^{i\theta}) = e^{in\theta}.$$

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The boundary condition is *compatible* with the rotating frame ansatz above, iff  $k|n$  and  $m = \frac{n}{k}$ . Inspired by relative equilibria, look for  $v$  with  $k$ -fold symmetry.

In case  $n = 0$ , use the ansatz with  $k = 0$  and  $m$  an arbitrary integer, reflecting  $m$ -fold symmetry.

# Conserved Quantities

- Hamiltonians:

Gross-Pitaevskii: Ginzburg-Landau Energy:

$$E_\epsilon(u) := \int_{\mathbb{D}} \frac{|\nabla u|^2}{2} + \frac{(1 - |u|^2)^2}{4\epsilon^2} dx.$$

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- Momenta:

Gross-Pitaevskii:

$$J(v) := -\frac{1}{2} \int_{\mathbb{D}} k|v|^2 + \frac{1}{m} v \cdot (x^\perp \cdot \nabla) v^\perp dx.$$

Point Vortex Flow:

$$J_0(\mathbf{b}, d) = -\frac{1}{2} \sum_{j=1}^N d_j |b_j|^2$$

## Constrained Minimization Approach

The elliptic PDE above has a variational formulation based on momentum-constrained minimization. Since this is a minimization procedure, can only yield  $+1$  vortices.

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**General Idea:** Fix a relative equilibrium, to (PVF) whose vortices are aligned rather than staggered. Then consider the problem

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**General Idea:** Fix a relative equilibrium, to (PVF) whose vortices are aligned rather than staggered. Then consider the problem

$$\min_{u \in A} E_\varepsilon(u)$$

- $A$ : admissible set reflecting symmetry of the chosen relative equilibrium, and
- $J(u) \approx J_0(a_1, \dots, a_n)$ .

$\omega = \omega_\varepsilon$  arises as a Lagrange multiplier. This approach follows work by Gelantalis and Sternberg.



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- Assuming this, in the case of a single ring, the momentum constraint value determines the position of the vortices up to a rotation.
- Complete the proof using the vortex balls construction and the Jacobian estimate.

# Limitations of the constrained minimization approach

- Unable to treat multiple ring solutions/staggered ring solutions.
- Unable to show  $\omega_\varepsilon \rightarrow \omega$  where  $\omega$  is the speed corresponding to the limit.

## Alternative approach: Linking

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**Good news:** no need to control  $\omega_\varepsilon$  any more.

**Bad news:** we lose the constrained minimization formulation from above: can't specify constraint value and Lagrange multiplier!



# Linking

**Definition:** Fix a Banach space  $V$ , a closed subset  $S \subset V$  and a submanifold  $Q$ , and denote its relative boundary by  $\partial Q$ . The sets  $S$  and  $\partial Q$  are said to link if

- $S \cap \partial Q = \emptyset$
- For any continuous map  $h : V \rightarrow V$  such that  $h|_{\partial Q} = id$ , there holds  $h(Q) \cap S \neq \emptyset$ .

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In the context of Ginzburg Landau, a linking method was used by F-H. Lin to construct critical points of GL *near* critical points of  $W$ .

# Main Theorem

## Theorem (V., '16)

Let  $(\mathbf{a}, d)$  be a relative equilibrium, with speed  $\omega_0$ . Write  $\mathbf{a}(t) := \mathbf{a}e^{i\omega_0 t}$ . For each  $\varepsilon > 0$  sufficiently small (depending on  $\mathbf{a}$ ), there exists a non-trivial time periodic solution  $u_\varepsilon$  to (GP-BC), with the same period of rotation as the given relative equilibrium, such that the Jacobian

$$Ju_\varepsilon(\cdot, t) \rightarrow \pi \sum_{i=1}^N d_i \delta_{a_i}(t)$$

as  $\varepsilon \rightarrow 0$ , in  $W^{-1,1}(\mathbb{D})$ , for each time  $t \in \mathbb{R}$ .

Here, one can think of  $Ju_\varepsilon(\cdot, t) := \det(\nabla u_\varepsilon(\cdot, t)) dx$ .

## Some Details in the Proof of the Main Theorem

**Main Goal:** Find a critical point of the functional  $\mathcal{E}_\varepsilon := E_\varepsilon - \omega_0 J$  near a given critical point  $(\mathbf{a}, d)$  of  $\mathcal{H}^{n, \omega_0}(\mathbf{a}, d) := \frac{1}{\pi} W(\mathbf{a}, d) - \omega_0 J_0(\mathbf{a}, d)$

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- **Critical value of  $\mathcal{E}_\varepsilon$ :** Using this family, we can give an  $\inf - \sup$  characterization of the critical value, which, upto multiples of  $\pi \log \frac{1}{\varepsilon}$  and  $O(1)$  terms, is the energy  $\mathcal{H}^{n,\omega_0}(\mathbf{a}, d)$ .

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- **Conclusion of Critical Value Step:**  $\mathcal{E}_\varepsilon$  is Palais-Smale, so pass to the large time limit, holding  $\varepsilon$  fixed. Obtain a critical point  $v_\varepsilon$  satisfying

$$\left| \mathcal{E}_\varepsilon(v_\varepsilon) - N\pi \log \frac{1}{\varepsilon} - \mathcal{H}^{n,\omega_0}(\mathbf{a}) - N\gamma \right| = o_\varepsilon(1)$$

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- Above argument only says energies are close. We need  $v_\varepsilon$  to have zeroes close to the given critical point of  $\mathcal{H}$ . Follows from Pohazaev-type identities and letting  $\varepsilon \rightarrow 0^+$ .

## Afterthought: some examples

- For each integer  $k$ , there exists a solution to Gross-Pitaevskii with boundary condition  $g \equiv 1$ , and zeroes on the vertices of concentric  $k$ -gons, one with  $+1$  vortices, and the other, staggered, with  $-1$  vortices.



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- Given  $n$ , corresponding to the degree of the b.c., fix a divisor  $k$  of  $n$ . Then there exists a periodic orbit to (GP-BC) containing  $\frac{n}{k}$  aligned rings, each with  $k + 1$  vortices.
- When  $n$  is a prime, there's only one relative equilibrium to (GP-BC) with all  $+1$  vortices. Stability??