

Mathematics of Extended Inversion for Wave Propagation

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Overview

Takeaway: unitary property of modeling operator \Rightarrow convergent gradient computation for extended FWI

Agenda:

- ▶ Transmission FWI
- ▶ Model extension and Variable Projection
- ▶ Accurate VPM gradient computation

Agenda

Transmission FWI

Model extension and Variable Projection

Accurate VPM gradient computation

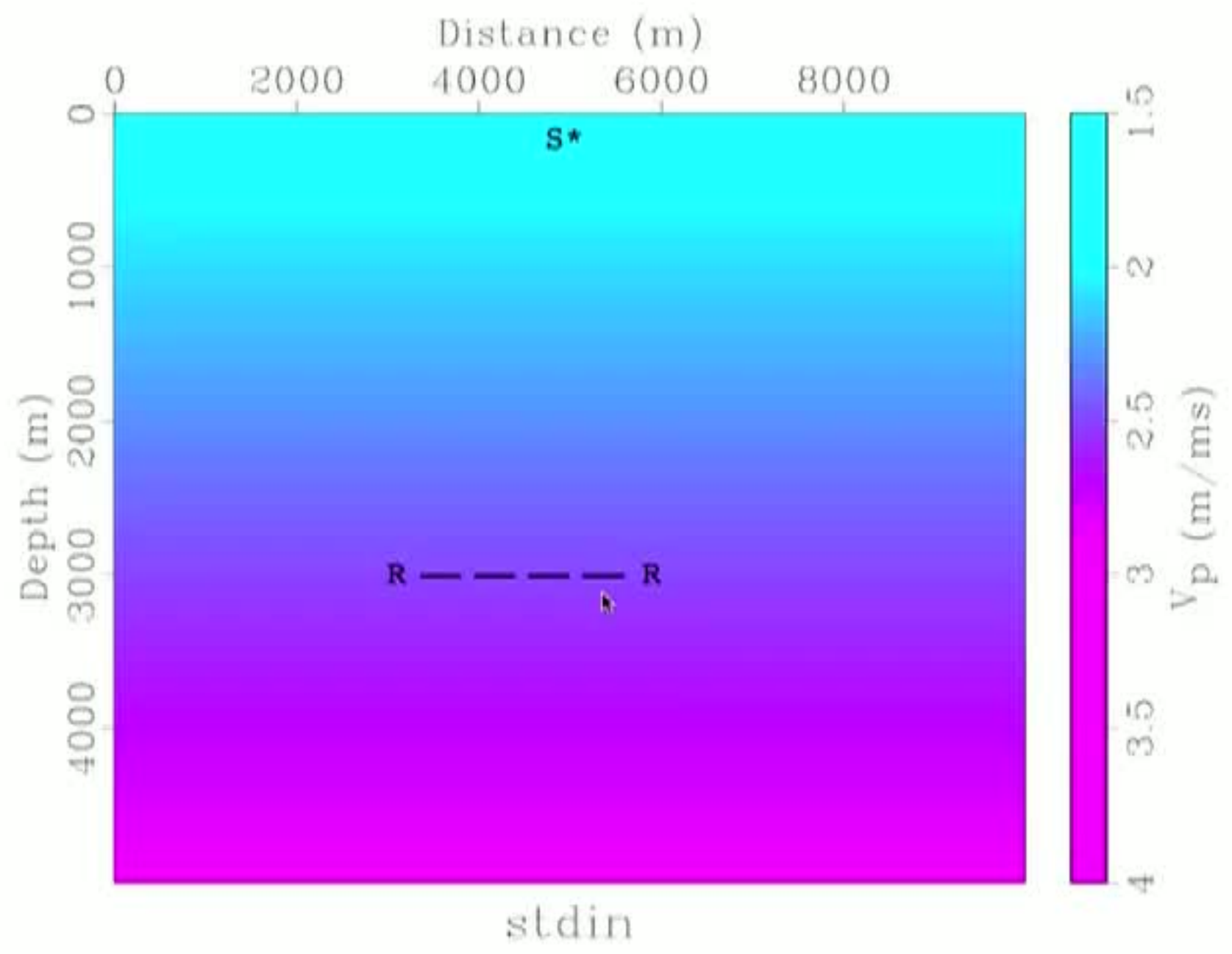
Simple example: transmission modeling, constant density acoustics with isotropic point radiator

$p(\mathbf{x}, t)$ = pressure, $v(\mathbf{x})$ = wave speed, $w(t)$ = source "wavelet"

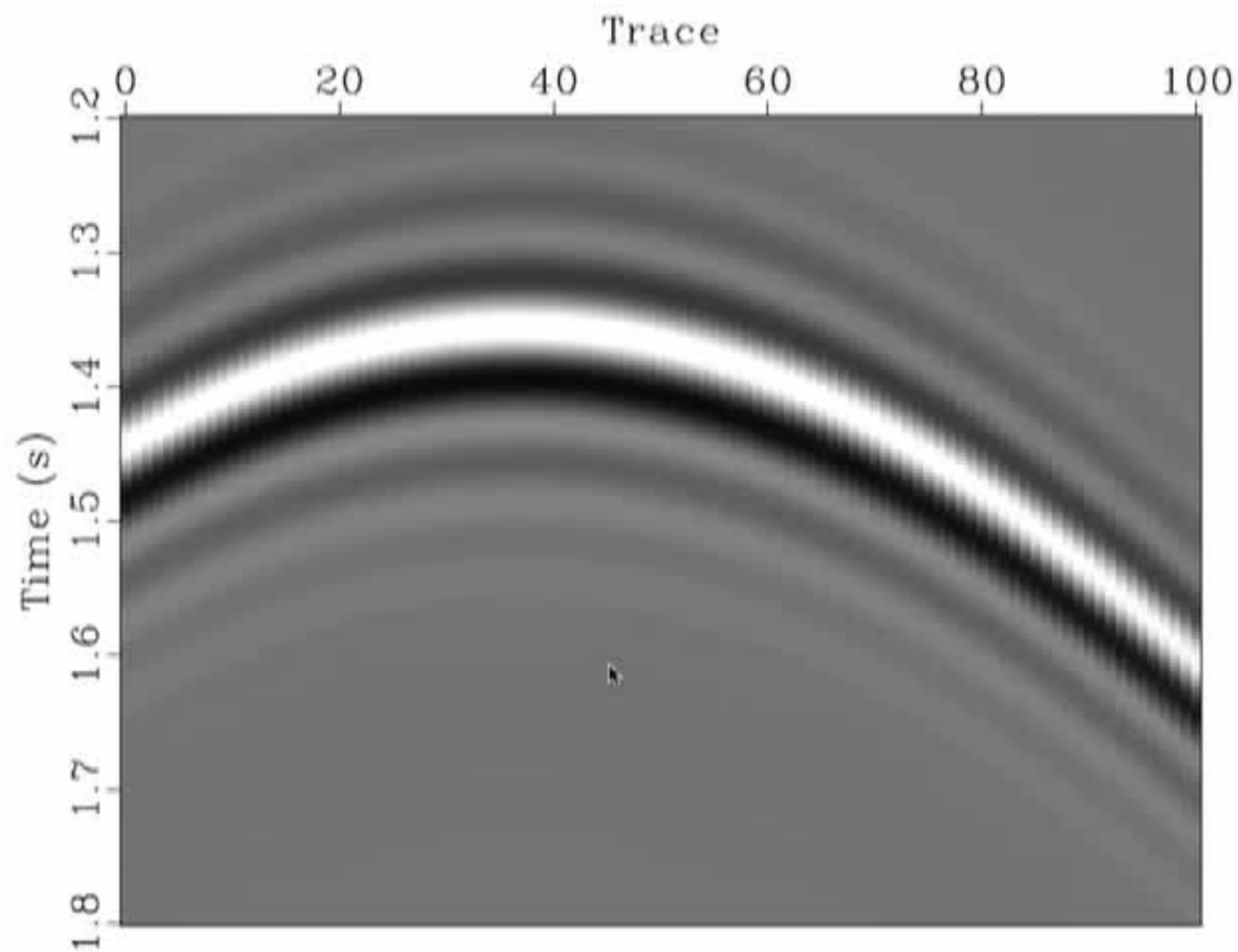
$$\left(\frac{1}{v(\mathbf{x})^2} \frac{\partial^2}{\partial t^2} - \nabla_{\mathbf{x}}^2 \right) p(\mathbf{x}, t; \mathbf{x}_s) = \delta(\mathbf{x} - \mathbf{x}_s) w(t);$$

$$p \equiv 0, t \ll 0$$

Predicted data for source \mathbf{x}_s , receiver $\mathbf{x}_r = p(\mathbf{x}_r, t; \mathbf{x}_s)$ (solve wave equation, sample pressure field)



Target velocity field v for horizontal well survey ("cross-well on its side")



"Observed" data for $\mathbf{x}_s = (5100\text{m}, 10\text{m})$:
 v as in previous slide, $w = [2, 5, 15, 20]$ Hz trapezoidal bandpass,
10 m \times 10 m grid, 2-8 centered FD

Informal problem setting:

M = set of admissible models ("model space") = mechanical parameter fields (bulk modulus, density, $C_{ijkl}(\mathbf{x})$, energy sources...) - in this case, $v \in C^\infty(\mathbf{R}^3)$ and $w \in L^2(\mathbf{R})$ - causal

$(D, \langle \cdot, \cdot \rangle)$ = Hilbert space of observed data = $d \in L^2(\mathbf{R}_r^2 \times \mathbf{R}_t \times \mathbf{R}_s^2)$

$F : M \rightarrow D$ modeling (data prediction) operator = solve wave equations for pressure, displacement, ..., sample at \mathbf{x}_r, t (RHS = function of \mathbf{x}_s)

This map is *separable*: $M = V \times W$, F linear in W : $F[(v, w)] = S[v]w$,
 $S : V \rightarrow \mathcal{B}(W, D)$

$$F[(v, w)] = S[v]w = \{p(\mathbf{x}_r, t; \mathbf{x}_s) : (\mathbf{x}_r, \mathbf{x}_s) \in \Sigma, t \in [0, T]\}$$

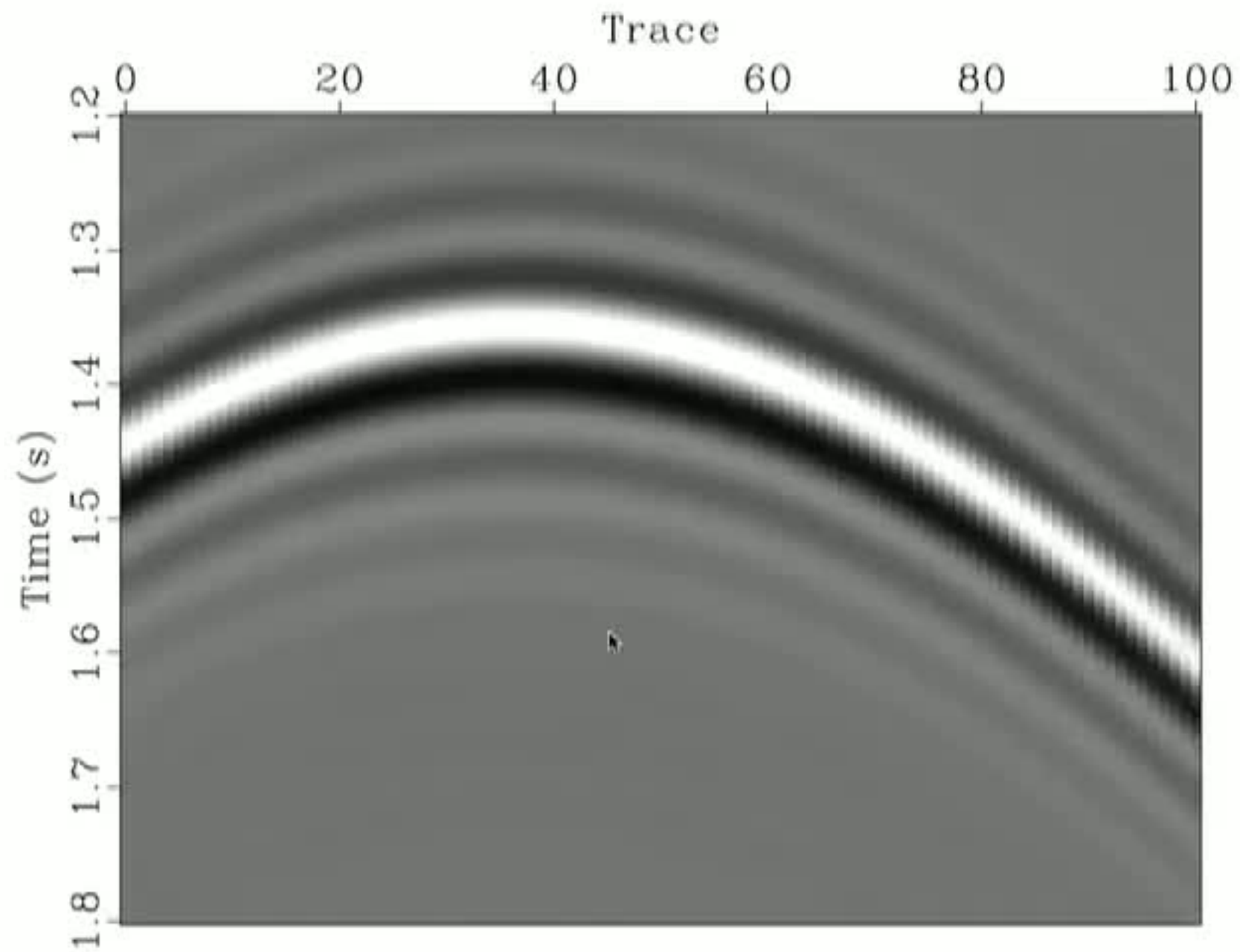
Inverse problem via Least Squares (“Full Waveform Inversion”):

given $d \in D$, find $m = (v, w) \in M$ so that

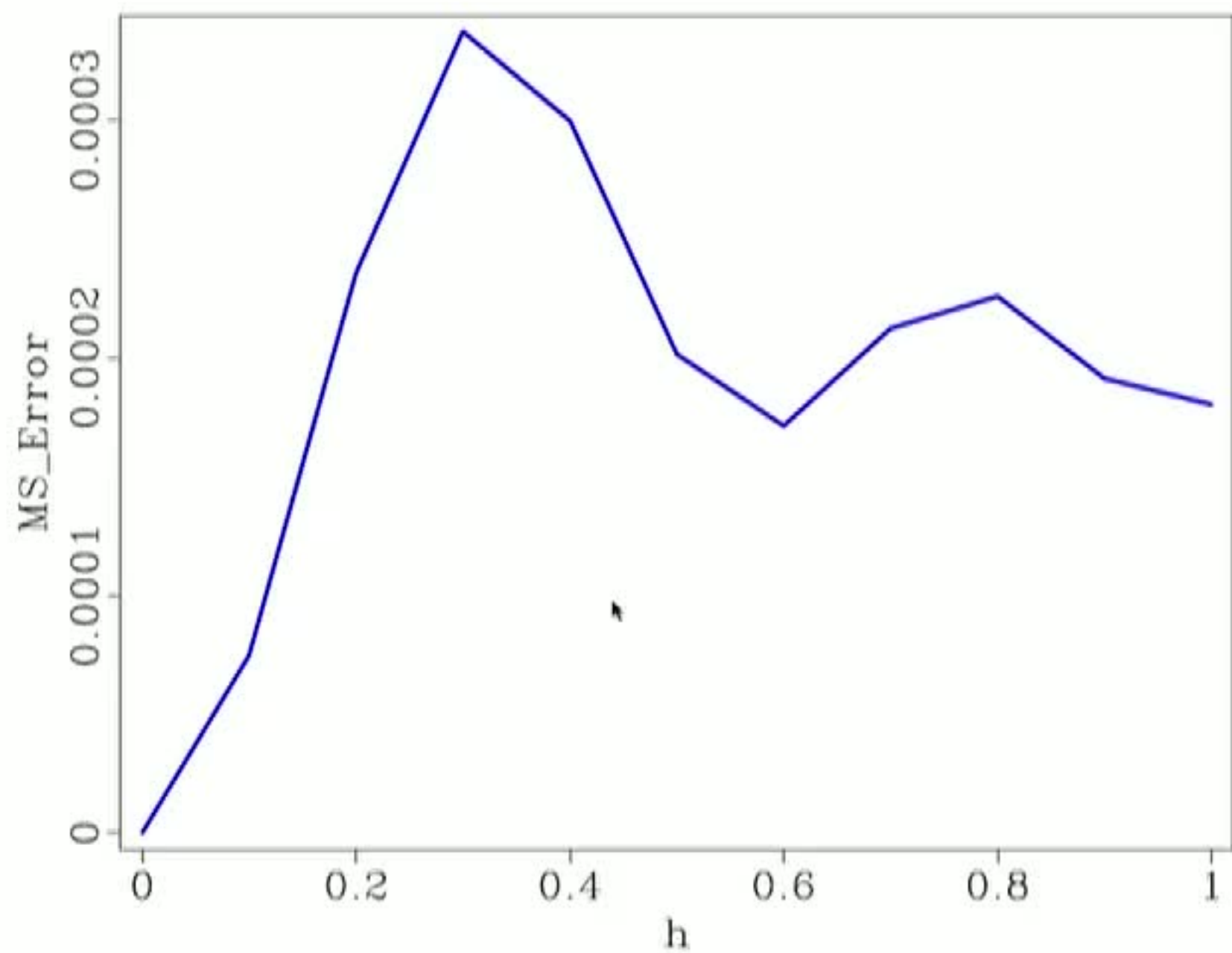
$$S[v]w \simeq d$$

by minimizing over v, w

$$J_{\text{FWI}}[v, w; d] = \frac{1}{2} \|S[v]w - d\|^2 [+ \text{regularizing terms}]$$



"Observed" data for $\mathbf{x}_s = (5100\text{m}, 10\text{m})$:
target ("exact") v



$$J_{\text{FWI}}[(1 - h) * v + h * v_0, w; d]$$

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Model extension and Variable Projection

Accurate VPM gradient computation

Inference: local optimization of J_{FWI} (descent, Newton-like) unlikely to succeed in finding ν , starting at ν_0

Reason: data fit error $\approx 100\%$ unless ν nearly correct

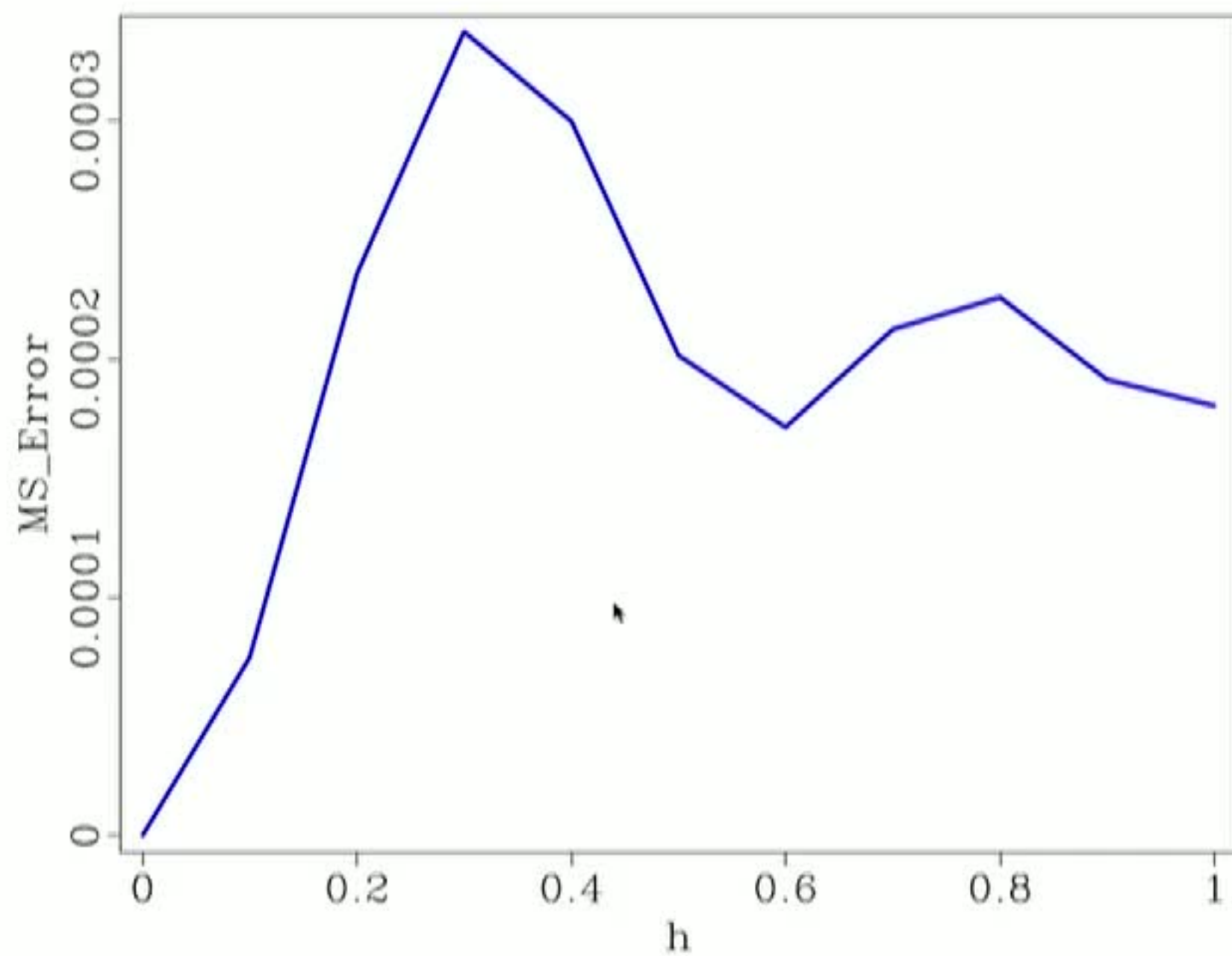
One (of many) possible fixes: enlarge set of models (model extension) to permit data fit for *any* ν , penalize additional degrees of freedom via *annihilator* = operator vanishing on original ("physical") models

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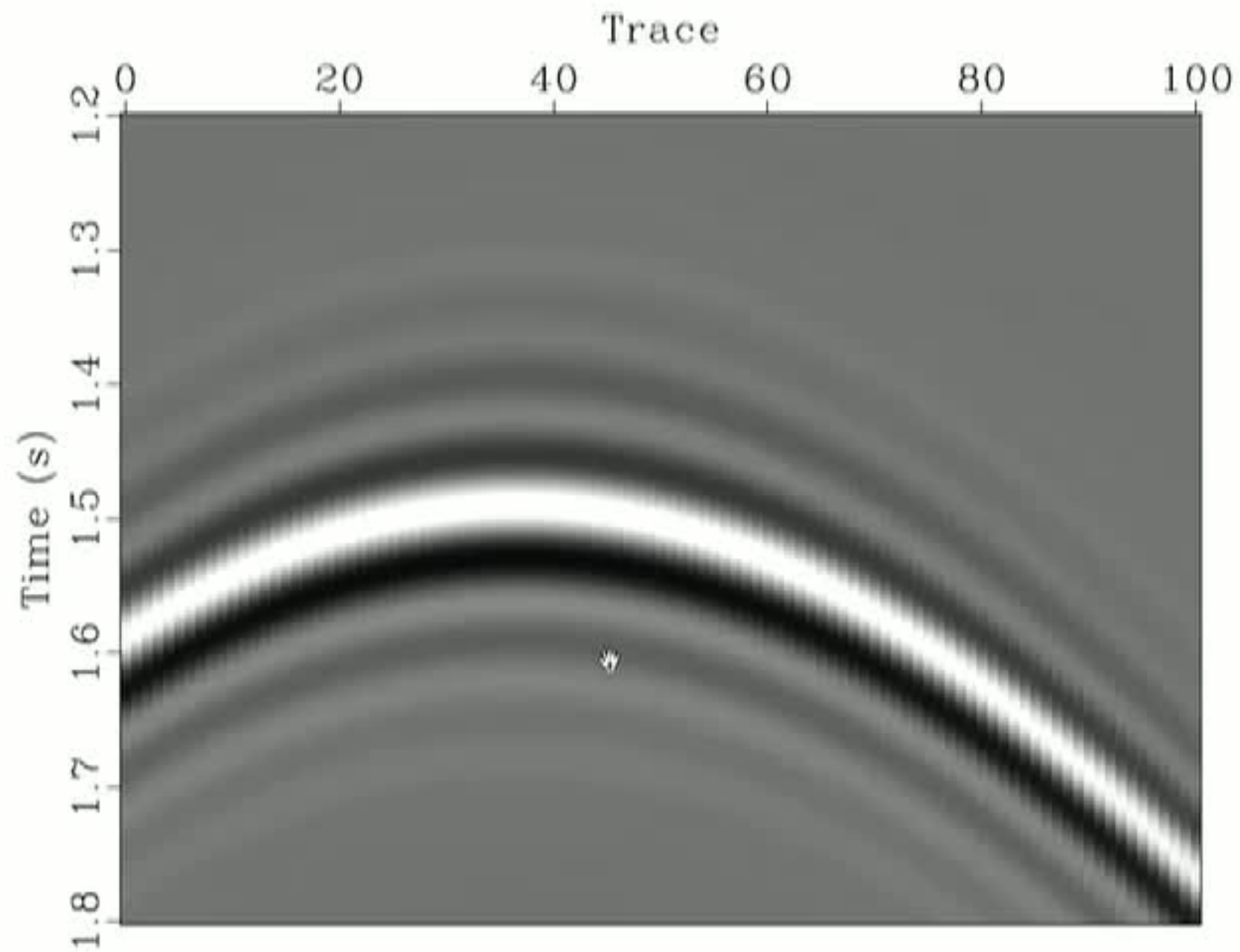
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"Predicted" data for $\mathbf{x}_s = (5100\text{m}, 10\text{m})$:
 $v = v_0 = 2 \text{ km/s}$ - possible initial velocity estimate

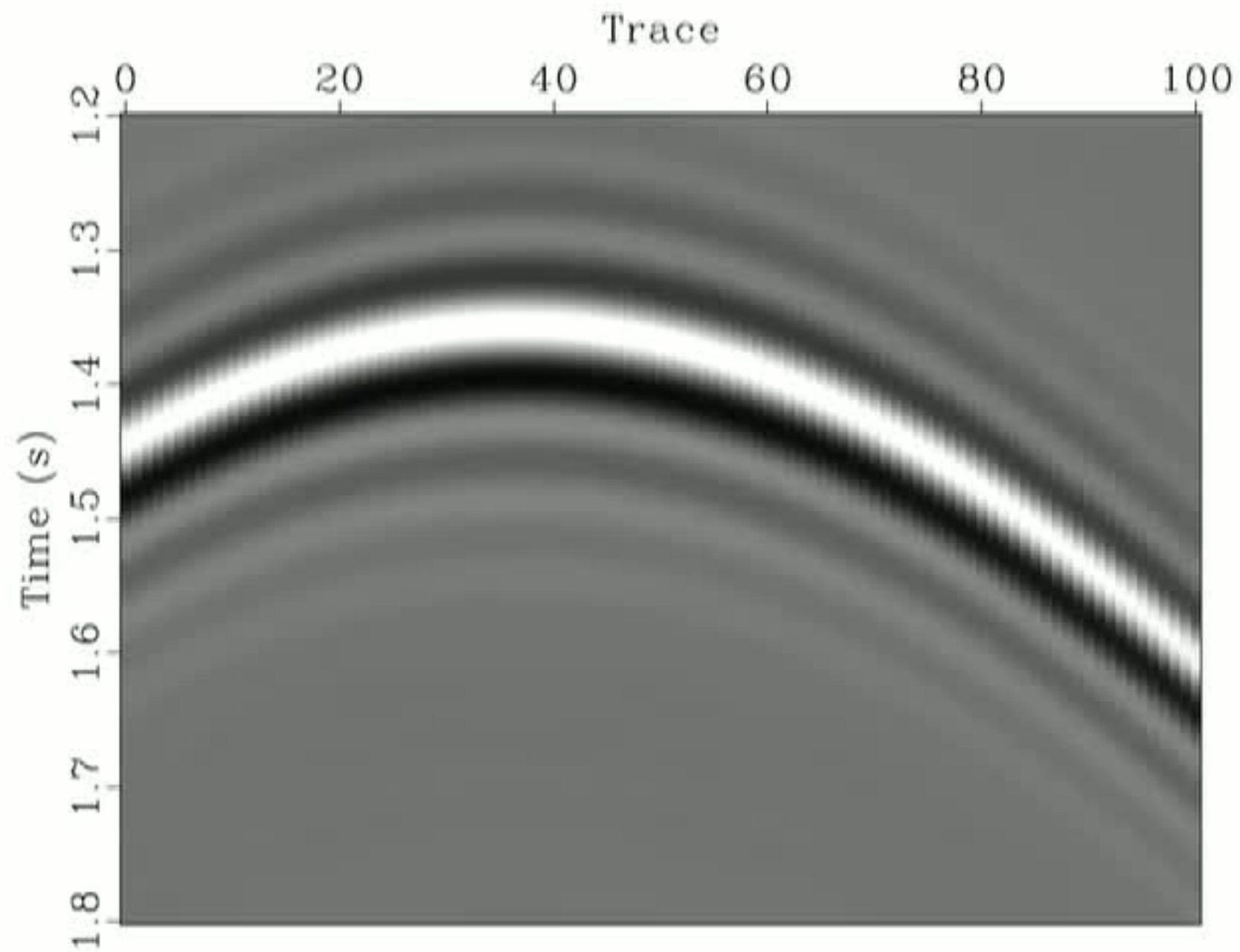
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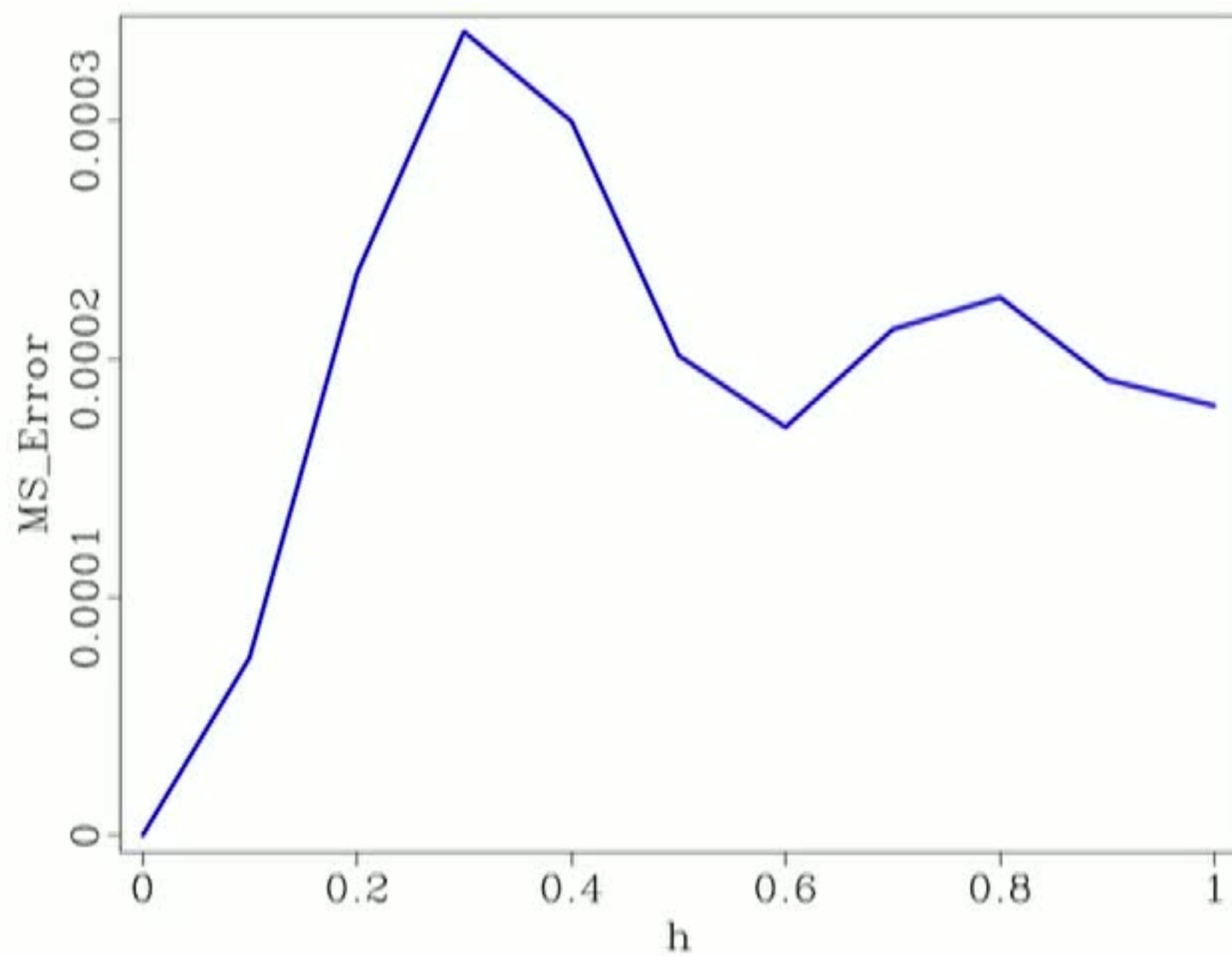
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Example of model extension: source-receiver extension - independent source wavelet $w(\mathbf{x}_r, t; \mathbf{x}_s)$ for each data trace

\bar{M} = set of extended models ("model space") = $v \in C^\infty(\mathbf{R}^3)$ and $\bar{w} \in L^2(\mathbf{R}_r^2 \times \mathbf{R}_t \times \mathbf{R}_s^2)$ (causal)

$(D, \langle \cdot, \cdot \rangle)$ = Hilbert space of observed data = $d \in L^2(\mathbf{R}_r^2 \times \mathbf{R}_t \times \mathbf{R}_s^2)$ (same!)

$\bar{F} : \bar{M} \rightarrow D$ modeling (data prediction) operator = solve wave equations for pressure, displacement, ..., sample at \mathbf{x}_r, t (RHS = function of $\mathbf{x}_r, \mathbf{x}_s$)

$$\bar{F}[(v, \bar{w})] = \bar{S}[v]\bar{w} = \{p(\mathbf{x}_r, t; \mathbf{x}_s) : (\mathbf{x}_r, \mathbf{x}_s) \in \Sigma, t \in [0, T]\}$$

Asymptotics (Hadamard, 3D): absent conjugate points,

$$\bar{S}[v]\bar{w}(\mathbf{x}_r, t; \mathbf{x}_s) \approx a(\mathbf{x}_r, \mathbf{x}_s)\bar{w}(\mathbf{x}_r, t - \tau[v](\mathbf{x}_r, \mathbf{x}_s); \mathbf{x}_s)$$

τ = travel time, a = geometric amplitude

\Rightarrow can fit any data with small error - take

$$\bar{w}(\mathbf{x}_r, t; \mathbf{x}_s) = d(\mathbf{x}_r, t + \tau[v](\mathbf{x}_r, \mathbf{x}_s); \mathbf{x}_s) / a(\mathbf{x}_r, \mathbf{x}_s)$$

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\Rightarrow fit error does not constrain v - restore constraint by penalizing extended degrees of freedom (dependence of \bar{w} on $\mathbf{x}_r, \mathbf{x}_s$)

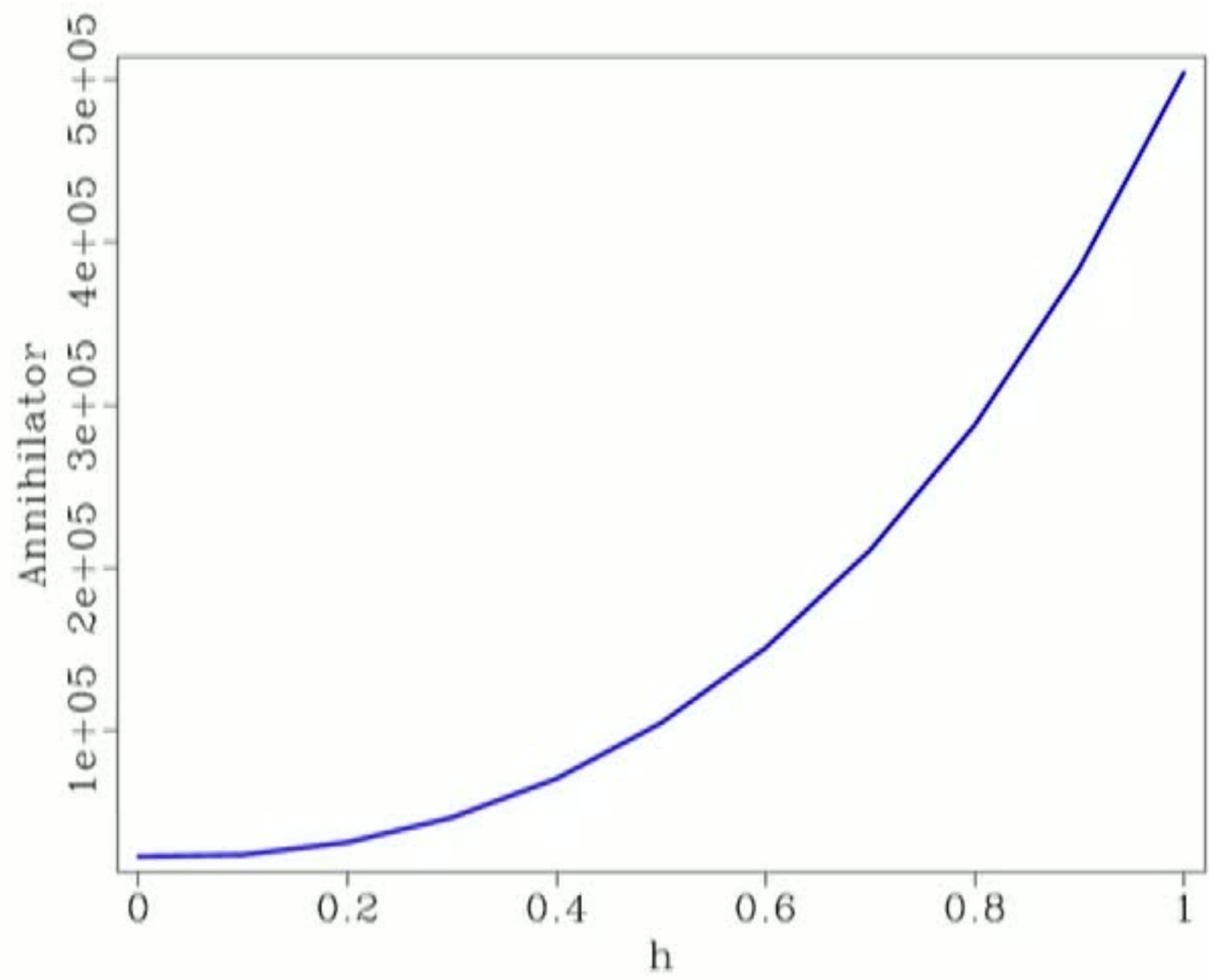
Many choices - eg. presume correct w known, pre-process (deconvolve) so that $w(t) \approx \delta(t)$. Popular choice of annihilator (Plessix et al. 00, Warner & Guatsch 14): $A\bar{w}(\mathbf{x}_r, t; \mathbf{x}_s) \equiv t\bar{w}(\mathbf{x}_r, t; \mathbf{x}_s) \approx 0$ if $\bar{w}(\dots, t) \approx \delta(t)$

Penalized least squares: minimize

$$J_{\text{ELS}}[v, \bar{w}] = \frac{1}{2}(\|\bar{S}[v]\bar{w} - d\|^2 + \alpha^2\|A\bar{w}\|^2)$$

Reduced objective (Variable Projection Method, Golub & Pereyra 73, 03)

$$\bar{w}[v] = \operatorname{argmin}_{\bar{w}} J_{\text{ELS}}[v, \bar{w}]; \quad J_{\text{ELS}}^{\text{red}}[v] = J_{\text{ELS}}[v, \bar{w}[v]]$$



$$J_{\text{ELS}}^{\text{red}}[(1 - h) * v + h * v_0; d]$$

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The Problem:

(1) nice VPM gradient formula

$$\nabla_v J_{\text{ELS}}^{\text{red}}[v; d] = D_v(\bar{S}[v]\bar{w})^T (\bar{S}[v]\bar{w} - d)_{\bar{w}=\bar{w}[v]}$$

requires $D_v\bar{S}$ - one order less regular than \bar{S} :

$$\begin{aligned} D_v\bar{S}[v]\bar{w}(\dots, t) &\approx D_v(a[v]\bar{w}(\dots, t - \tau[v]) + \dots) \\ &= -(a[v]D_v\tau[v])\frac{\partial\bar{w}}{\partial t} + \dots \end{aligned}$$

(2) $\bar{w}[v]$ estimated iteratively = $\lim_{a \rightarrow \infty} \bar{w}_a$ in L^2

\Rightarrow obvious gradient approximation

$$\nabla_v J_{\text{ELS}}^{\text{red}}[v; d] \approx D_v(\bar{S}[v]\bar{w})^T (\bar{S}[v]\bar{w} - d)_{\bar{w}=\bar{w}_a}$$

may not converge (Kern & S. 94, Yin Zhang thesis 16)

A Solution:

Part of VPM gradient/deriv construction is harmless:

$$D_\nu J_{\text{ELS}}^{\text{red}}[\nu; d] = \langle D\bar{S}[\nu]\bar{w}, \bar{S}[\nu]\bar{w} - d \rangle_{\bar{w}=\bar{w}[\nu]}$$

$$= \mathbf{1} + \mathbf{2}$$

$$\mathbf{1} = \langle D_\nu \bar{S}[\nu]\bar{w}, \bar{S}[\nu]\bar{w} \rangle_{\bar{w}=\bar{w}[\nu]} = \left(D_\nu \frac{1}{2} \langle \bar{w}, \bar{S}[\nu]^T \bar{S}[\nu]\bar{w} \rangle \right)_{\bar{w}=\bar{w}[\nu]}$$

$\bar{S}[\nu]^T \bar{S}[\nu]$ is Ψ DO w symbol smooth in $\nu \Rightarrow$

$$\mathbf{1} = \lim_{a \rightarrow \infty} \langle D_\nu \bar{S}[\nu]\bar{w}, \bar{S}[\nu]\bar{w} \rangle_{\bar{w}=\bar{w}_a}$$

$\bar{S}[v]$ is approx. unitary in appropriate weighted L^2 norm:

$$\bar{S}[v]^T \bar{S}[v] = a[v]^2 I \text{ modulo } OPS^{-1}$$

Define

$$\langle \bar{w}_1, \bar{w}_2 \rangle_v = \langle \bar{w}_1, a[v]^2 \bar{w}_2 \rangle_{L^2}$$

denote adjoint of $\bar{S}[v]$ wrt $\langle \cdot, \cdot \rangle_v$ by $\bar{S}[v]^\dagger = a[v]^{-2} \bar{S}[v]^T$

$$\Rightarrow \bar{S}[v]^\dagger \bar{S}[v] = I + K[v], \quad K[v] \in OPS^{-1}$$

Weighted normal residual for \bar{w}_a :

$$g_a = (\bar{S}[v]^\dagger \bar{S}[v] + \alpha^2 A^\dagger A)(\bar{w}_a - \bar{w}[v])$$

$$= (I + \alpha^2 A^\dagger A)(\bar{w}_a - \bar{w}[v]) + K[v](\bar{w}_a - \bar{w}[v]), \quad K[v] \in OPS^{-1}$$

\Rightarrow

$$\bar{w}_a - \bar{w}[v] = (I + \alpha^2 A^\dagger A)^{-1}(K[v](\bar{w}_a - \bar{w}[v]) - g_a)$$

\Rightarrow

$$\tilde{w}_a \equiv \bar{w}_a + (I + \alpha^2 A^\dagger A)^{-1}g_a = \bar{w}[v] \text{ modulo error } \rightarrow 0 \text{ in } H_t^1$$

\Rightarrow

$$\mathbf{2} = - \lim_{a \rightarrow \infty} \langle D\bar{S}[v]\bar{w}, d \rangle_{\bar{w}=\tilde{w}_a}$$

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\Rightarrow

$$\mathbf{2} = - \lim_{a \rightarrow \infty} \langle D\bar{S}[v]\bar{w}, d \rangle_{\bar{w}=\tilde{w}_a}$$

Upshot: $\bar{w}_a \rightarrow \bar{w}[v] \Rightarrow$

$$\nabla_v J_{\text{ELS}}^{\text{red}}[v; d] = \lim_{a \rightarrow \infty} D_v(\bar{S}[v]\bar{w})^T (\bar{S}[v]\bar{w} - d)_{\bar{w}=\tilde{w}_a}$$

where

$$\tilde{w}_a \equiv \bar{w}_a + (I + \alpha^2 A^\dagger A)^{-1} g_a$$

Exercise for audience: accounting for v -dep of weighted norms does not change conclusion

Computability:

- ▶ g_a, \tilde{w}_a computable *in principle* - in practice, feasible computation \Rightarrow additional $O(\lambda)$ error
- ▶ doesn't matter because $\|D\bar{S}\|_{L^2}$ actually $= O(\lambda^{-1})$

And so on...

Conjugate points \Rightarrow source-receiver extension loses unitary property (Song & S. 94, Huang & S. 17)

Surface source extension (Huang & S. 18) recovers unitary property - but only microlocally

- ▶ reliable gradient comp requires (computable) Ψ DO cutoff

Reflection inversion - subsurface offset extension (Stolk & de Hoop 01, Shen et al. 03,...) microlocally unitary at physical subspace (ten Kroode 12, Hou & S. 15)

- ▶ additional error absorbed in frequency continuation scheme (Fu & S. 17)