Spatiotemporal Intermittency and Chaos in a Ginzburg Landau System for Oscillatory Instabilities in Anisotropic Systems

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Prototype Example for Anisotropic Systems Electroconvection in Nematic Liquid Crystals

Symmetries in 2d axial anisotropic systems comprise:

Reflection invariances across and along a distinguished symmetry axis (director axis in case of nematic electroconvection)

In 2d (infinitely) extended anisotropic systems:

2 translation and 2 reflection symmetries: E(1) X E(1)

Since there is no Rotation Symmetry:

Instabilities of a uniform state involve only finitely many critical wave numbers

(in contrast to isotropic systems, where one has a circle of critical wave numbers)

Three Generic Possibilities for an Instability





Oblique Case: 4 critical wave numbers (±p_c, ±q_c) Normal Case: 2 critical wave numbers (±p_c,0)

Third Case: critical wave number (0,0) Linearized system admits two counterpropagating pairs of traveling wave solutions



Extend Study to Nonlinear System through Weakly Nonlinear Analysis Represent field variables through small and slowly varying amplitudes ϵA_1 , ϵA_2 , ϵA_3 , ϵA_4 :

$$U(t,x,y,z) = \mathcal{E}\{A_{1}U_{1}(z)e^{i(p_{c}x+q_{c}y)} + A_{2}U_{2}(z)e^{i(-p_{c}x+q_{c}y)} + A_{3}U_{3}(z)e^{i(-p_{c}x-q_{c}y)} + A_{4}U_{4}(z)e^{i(p_{c}x-q_{c}y)}\}e^{i\omega_{c}t} + cc + O(\varepsilon^{2})$$

where $A_1 = A_1(T, X_+, Y_+), A_2 = A_2(T, X_-, Y_+), A_3 = A_3(T, X_-, Y_-), A_4 = A_4(T, X_+, Y_-)$ with $X_{\pm} = \varepsilon(t \pm x/v_p), \quad Y_{\pm} = \varepsilon(t \pm y/v_q), \quad T = \varepsilon^2 t, \quad R - R_c = \varepsilon^2$ Represent field variables through small and slowly varying amplitudes ϵA_1 , ϵA_2 , ϵA_3 , ϵA_4 :

$$\begin{split} U(t,x,y,z) &= \varepsilon \{ \mathsf{A}_1 U_1(z) e^{i(p_c x + q_c y)} + \mathsf{A}_2 U_2(z) e^{i(-p_c x + q_c y)} + \\ &\quad \mathsf{A}_3 U_3(z) e^{i(-p_c x - q_c y)} + \mathsf{A}_4 U_4(z) e^{i(p_c x - q_c y)} \} e^{i\omega_c t} + cc + O(\varepsilon^2) \end{split}$$

where $A_1 &= A_1(T, X_+, Y_+), A_2 = A_2(T, X_-, Y_+), A_3 = A_3(T, X_-, Y_-), A_4 = A_4(T, X_+, Y_-)$
with $X_{\pm} = \varepsilon(t \pm x/v_p), \quad Y_{\pm} = \varepsilon(t \pm y/v_q), \quad T = \varepsilon^2 t, \quad R - R_c = \varepsilon^2$

System of Globally Coupled Ginzburg Landau Equations for the A_j:

$$\partial_{T}A_{1}(T, X, Y) = \left(a_{0} + D(\partial_{X}, \partial_{Y}) + a_{1} |A_{1}|^{2} + a_{2}\left\langle |A_{2}(\chi, Y)|^{2} \right\rangle + a_{3}\left\langle |A_{3}(X - \chi, Y - \chi)|^{2} \right\rangle + a_{4}\left\langle |A_{4}(X, \chi)|^{2} \right\rangle \right)A_{1} + a_{5}\left\langle A_{2}(\chi - X, Y)\overline{A}_{3}(\chi - X, \chi - Y)A_{4}(X, \chi - Y) \right\rangle$$

where

$$\langle g(\chi) \rangle = \lim_{L \to \infty} \frac{1}{2L} \int_{-L}^{L} g(\chi) d\chi = \text{average over } \chi$$

Numerically:
$$A_j(T, X, Y) = \sum_{m=-M}^{M} \sum_{n=-N}^{N} a_j(m, n, T) e^{imp_0 X + inq_0 Y}$$

Without spatial variations: $A_j = A_j(T) = a_j(0,0,T)$

Numerically:
$$A_j(T, X, Y) = \sum_{m=-M}^{M} \sum_{n=-N}^{N} a_j(m, n, T) e^{imp_0 X + inq_0 Y}$$

Without spatial variations: $A_j = A_j(T) = a_j(0,0,T) \rightarrow Normal Form for a Hopf Bifurcation with O(2) X O(2)-Symmetry:$

$$\frac{d}{dT}A_{1} = (a_{0} + a_{1}|A_{1}|^{2} + a_{2}|A_{2}|^{2} + a_{3}|A_{3}|^{2} + a_{4}|A_{4}|^{2})A_{1} + a_{5}A_{2}\overline{A}_{3}A_{4}$$

Numerically: $A_j(T, X, Y) = \sum_{m=-M}^{M} \sum_{n=-N}^{N} a_j(m, n, T) e^{imp_0 X + inq_0 Y}$

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Basic wave solutions:

- TW: Traveling waves(A,0,0,0)SW: Standing waves(A,0,A,0)TR_x: Traveling rectangles in x(A,0,0,A)
- TR_y: Traveling rectangles in y (A,A,0,0)
- SR: Standing rectangles (A,A,A,A)
- AW: Alternating waves

Other solutions:

Quasiperiodic waves in 4d Various heteroclinic cycles

(A, iA, A, iA)







Examples of Patterns shown by the GL System



NF has attracting heteroclinic cycle TR_x -SR-AW









Normal Form Chaos

For certain parameters, NF shows a chaotic attractor created through a period doubling cascade; no stable basic waves.



Iterates (variable of a suitable Poincare map) vs -a_{3r}

Chaotic NF-Attractors at $-a_{3r}=0.6442$: 4 symmetry-conjugated copies $(R_j=|A_j|)$



Nearby Symmetrized Attractors (2 copies)





R₂ versus R₃

Corresponding GL-Dynamics





Corresponding GL-Dynamics





Associated Pattern Dynamics



Low-dimensional model for switches (PhD-work of Zou) Perturbed NF: Breaking of y->y+φ

$$\frac{d}{dT}A_{1} = (a_{4} + a_{5})b^{2}A_{4} + (a_{0} + a_{1}|A_{1}|^{2} + a_{2}|A_{2}|^{2} + a_{3}|A_{3}|^{2} + a_{4}|A_{4}|^{2})A_{1} + a_{5}A_{2}\overline{A}_{3}A_{4}$$



Transition occurs at $b_c \approx 0.199$.

Mean switch time: $T_{switch} \sim |b-b_c|^{-\gamma}, \Upsilon \approx 3.9235$



Further increase of b:

Transition to a periodic orbit via In-Out-Intermittency at $b_c \approx 0.308$

Probability of mean time between bursts: $P(T_{mean}) \sim \alpha n^{-3/2} e^{-\beta n} + \Upsilon e^{-\delta n}$ (Ashwin's model), where $n=0.1T_{mean}$, $\alpha \approx 0.3321$, $\beta \approx 0.0524$, $\Upsilon \approx 0.0156$, $\delta \approx 0.0506$



Thank You