

# Reformulating spectral problems with the Krein matrix

Todd Kapitula

Department of Mathematics and Statistics  
Calvin College

SIAM Conference on Applications of Dynamical Systems  
(2015)

## My collaborators:

- Shamuel Auyeung, Eric Yu (Calvin undergraduate students)
- B. Deconinck (U. Washington), P. Miller (U. Michigan)

## The research has been partially supported by:

- Calvin College
  - Jack and Lois Kuipers Applied Mathematics Endowment
  - Calvin Research Fellowship
- NSF DMS-1108783

Goal is to understand the dynamics associated with perturbations of (small) spatially  $2\pi$ -periodic waves to Klein-Gordon-like (KG) equations:

$$\partial_t^2 u + \mathcal{M}u + f(u) = 0, \quad \mathcal{M} = \sum_{j=0}^N a_j \partial_x^{2j}, \quad \text{and } |f(u)| = \mathcal{O}(u^2).$$

Ideas also applicable to:

- KdV-like:  $\partial_t u + \partial_x \left( \mathcal{M}u + f(u, \partial_x u, \partial_x^2 u, \dots) \right) = 0$
- NLS-like

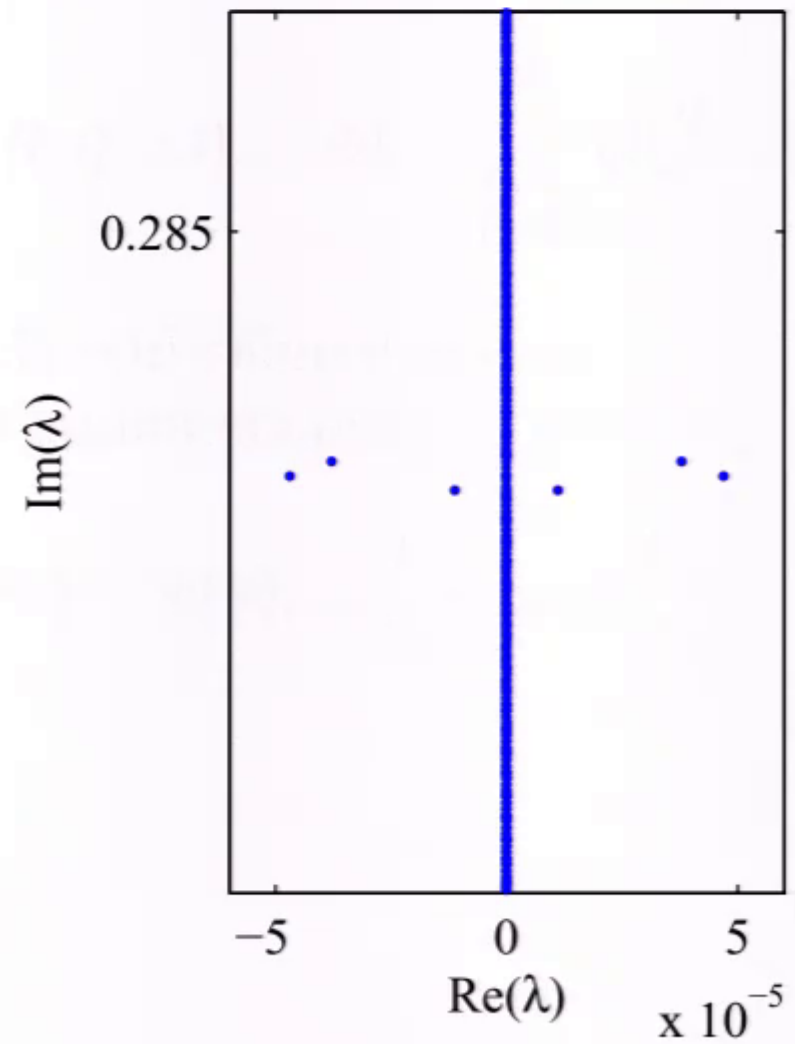
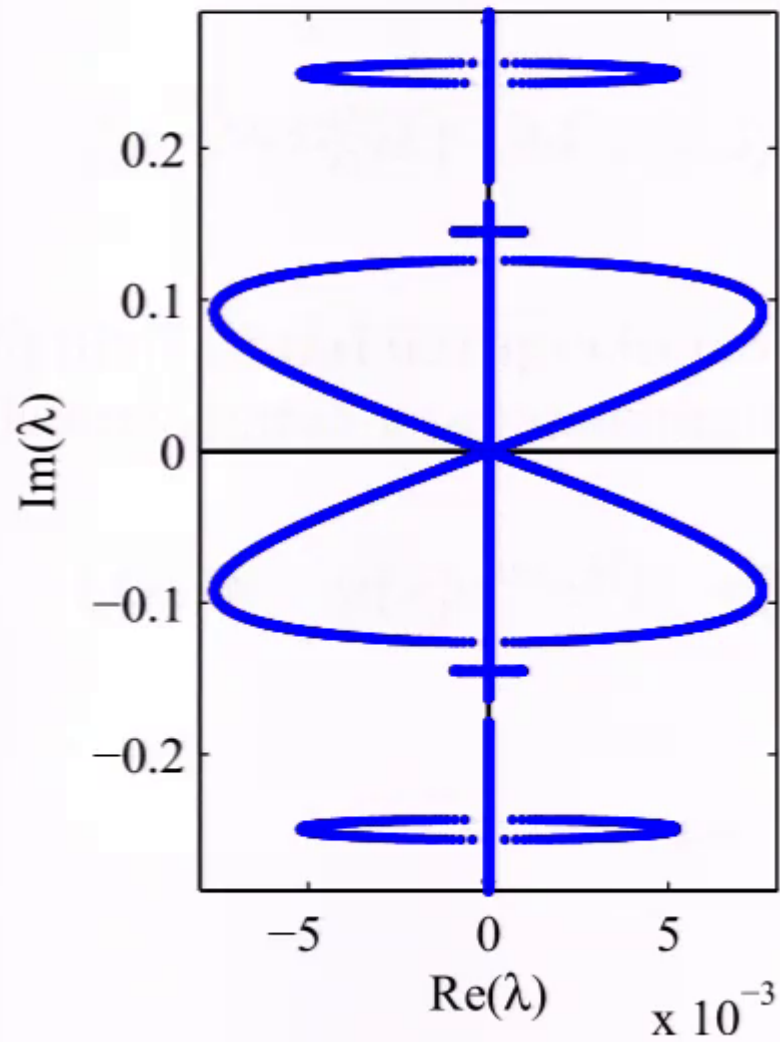
In traveling coordinates,  $z = x - ct$ , KG becomes

$$\partial_t^2 u - 2c\partial_{tz}^2 u + (\mathcal{M} + c^2\partial_z^2)u + f(u) = 0, \quad \mathcal{M} = \sum_{j=0}^N a_j \partial_z^{2j}.$$

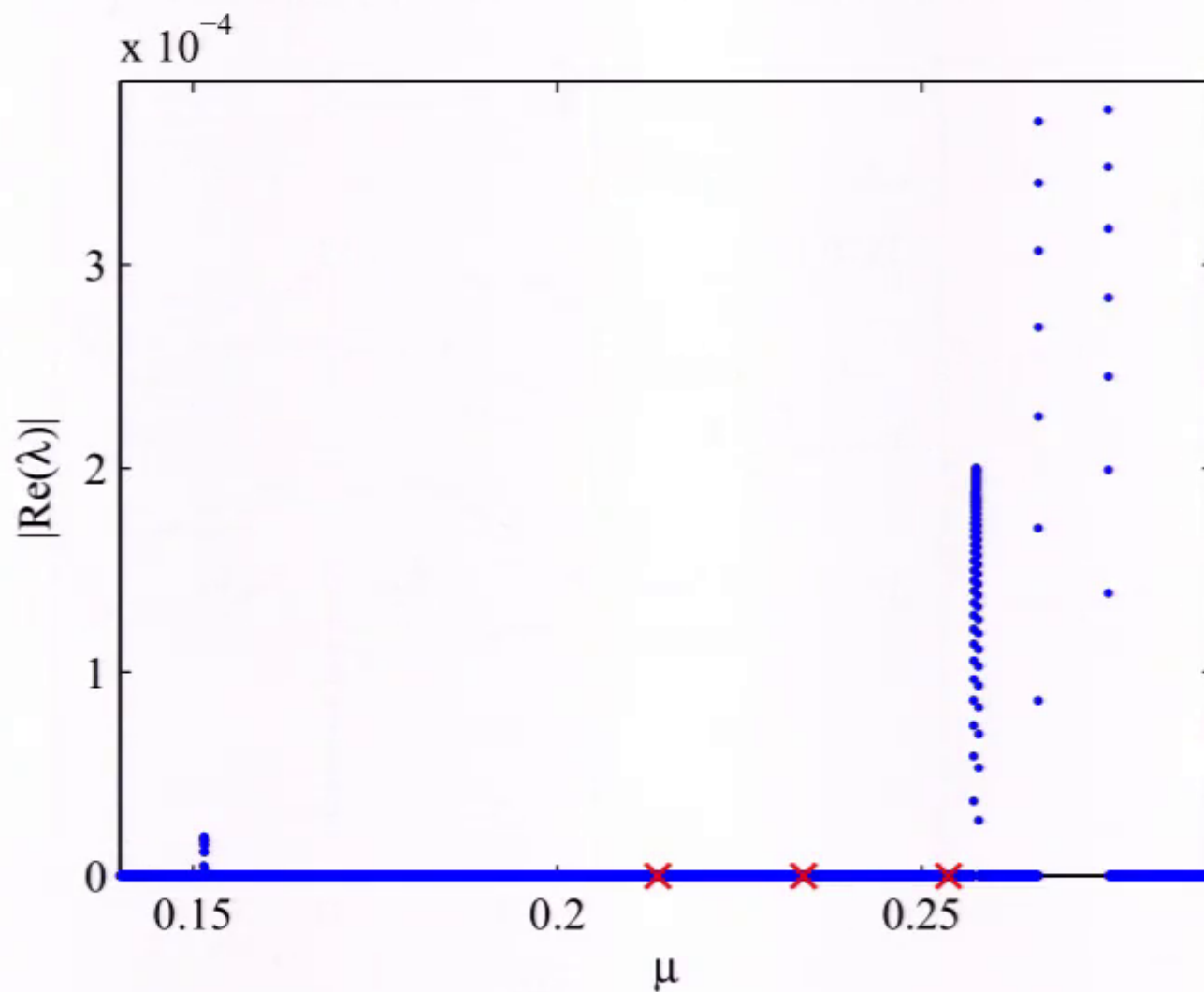
To understand the spectrum of the linearization use a Bloch-wave decomposition of the eigenfunctions:

$$u(z, t) = \psi(z)e^{i\mu z}e^{\lambda t}; \quad \psi(z + 2\pi) = \psi(z), \quad -\frac{1}{2} < \mu < \frac{1}{2}.$$

KG spectrum with wave amplitude  $\sim 0.16$ :



KG spectrum as a function of Bloch-wave parameter  $\mu$  with wave amplitude  $\sim 0.11$ :



In traveling coordinates,  $z = x - ct$ , KG becomes

$$\partial_t^2 u - 2c\partial_{tz}^2 u + (\mathcal{M} + c^2\partial_z^2)u + f(u) = 0, \quad \mathcal{M} = \sum_{j=0}^N a_j \partial_z^{2j}.$$

To understand the spectrum of the linearization use a Bloch-wave decomposition of the eigenfunctions:

$$u(z, t) = \psi(z)e^{i\mu z}e^{\lambda t}; \quad \psi(z + 2\pi) = \psi(z), \quad -\frac{1}{2} < \mu < \frac{1}{2}.$$

## My collaborators:

- Shamuel Auyeung, Eric Yu (Calvin undergraduate students)
- B. Deconinck (U. Washington), P. Miller (U. Michigan)

## The research has been partially supported by:

- Calvin College
  - Jack and Lois Kuipers Applied Mathematics Endowment
  - Calvin Research Fellowship
- NSF DMS-1108783



Goal is to understand the dynamics associated with perturbations of (small) spatially  $2\pi$ -periodic waves to Klein-Gordon-like (KG) equations:

$$\partial_t^2 u + \mathcal{M}u + f(u) = 0, \quad \mathcal{M} = \sum_{j=0}^N a_j \partial_x^{2j}, \quad \text{and } |f(u)| = \mathcal{O}(u^2).$$

Ideas also applicable to:

- KdV-like:  $\partial_t u + \partial_x \left( \mathcal{M}u + f(u, \partial_x u, \partial_x^2 u, \dots) \right) = 0$
- NLS-like

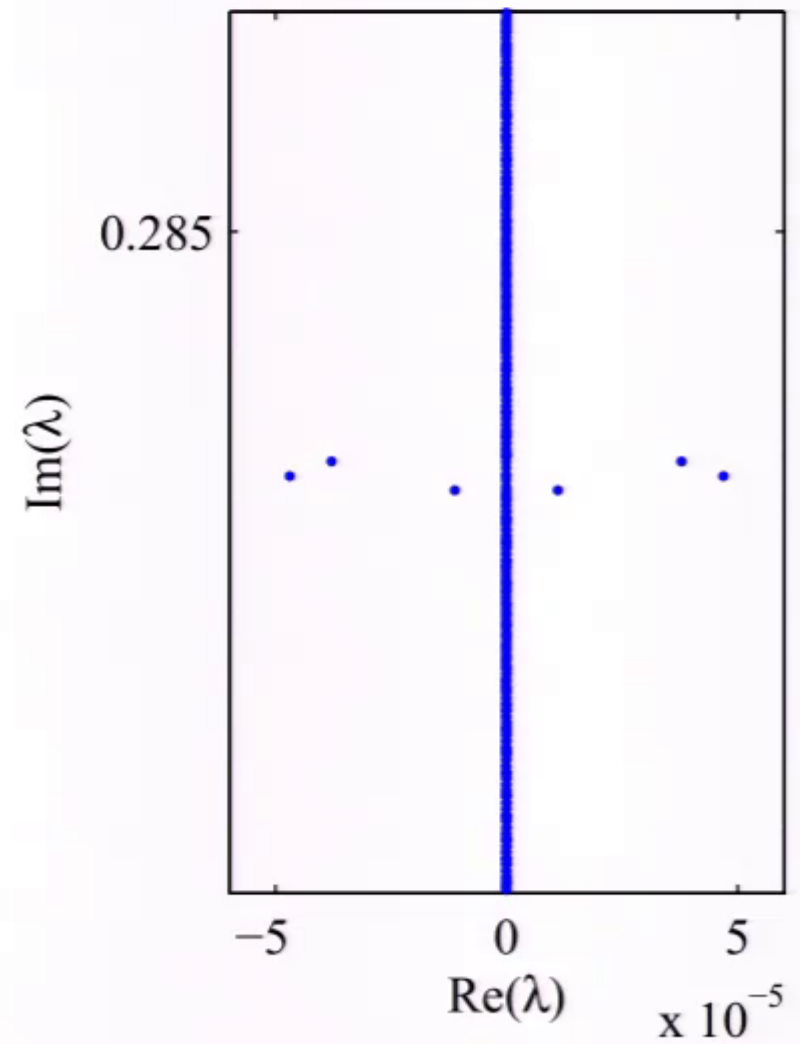
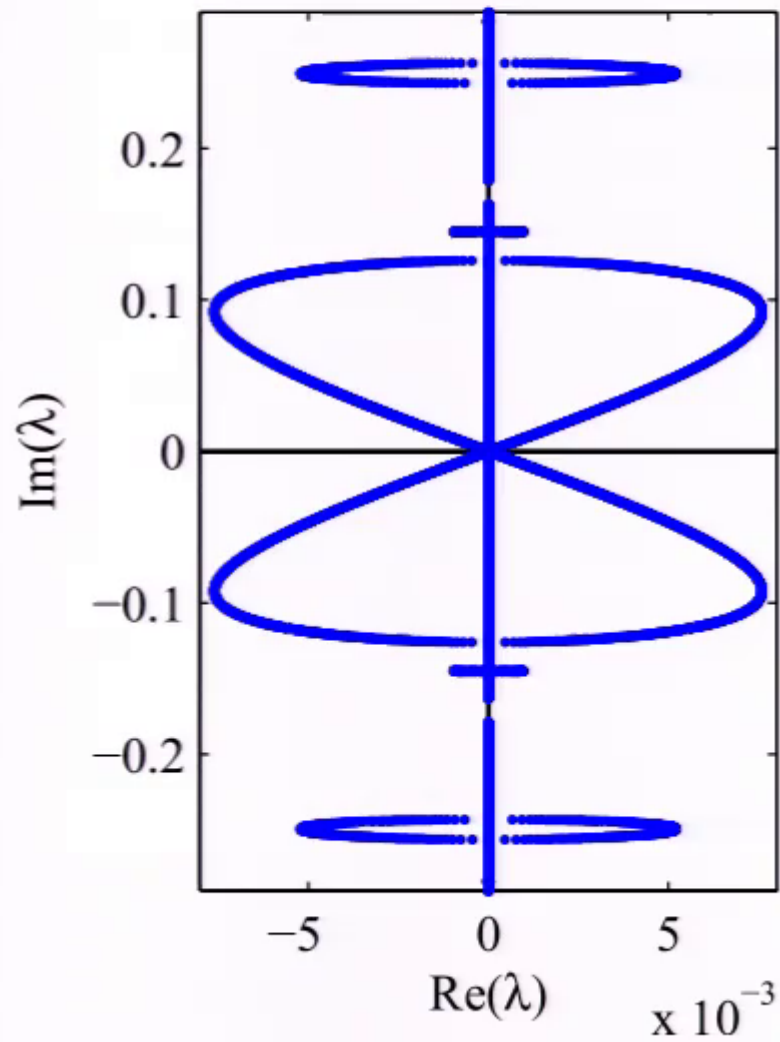
In traveling coordinates,  $z = x - ct$ , KG becomes

$$\partial_t^2 u - 2c\partial_{tz}^2 u + (\mathcal{M} + c^2\partial_z^2)u + f(u) = 0, \quad \mathcal{M} = \sum_{j=0}^N a_j \partial_z^{2j}.$$

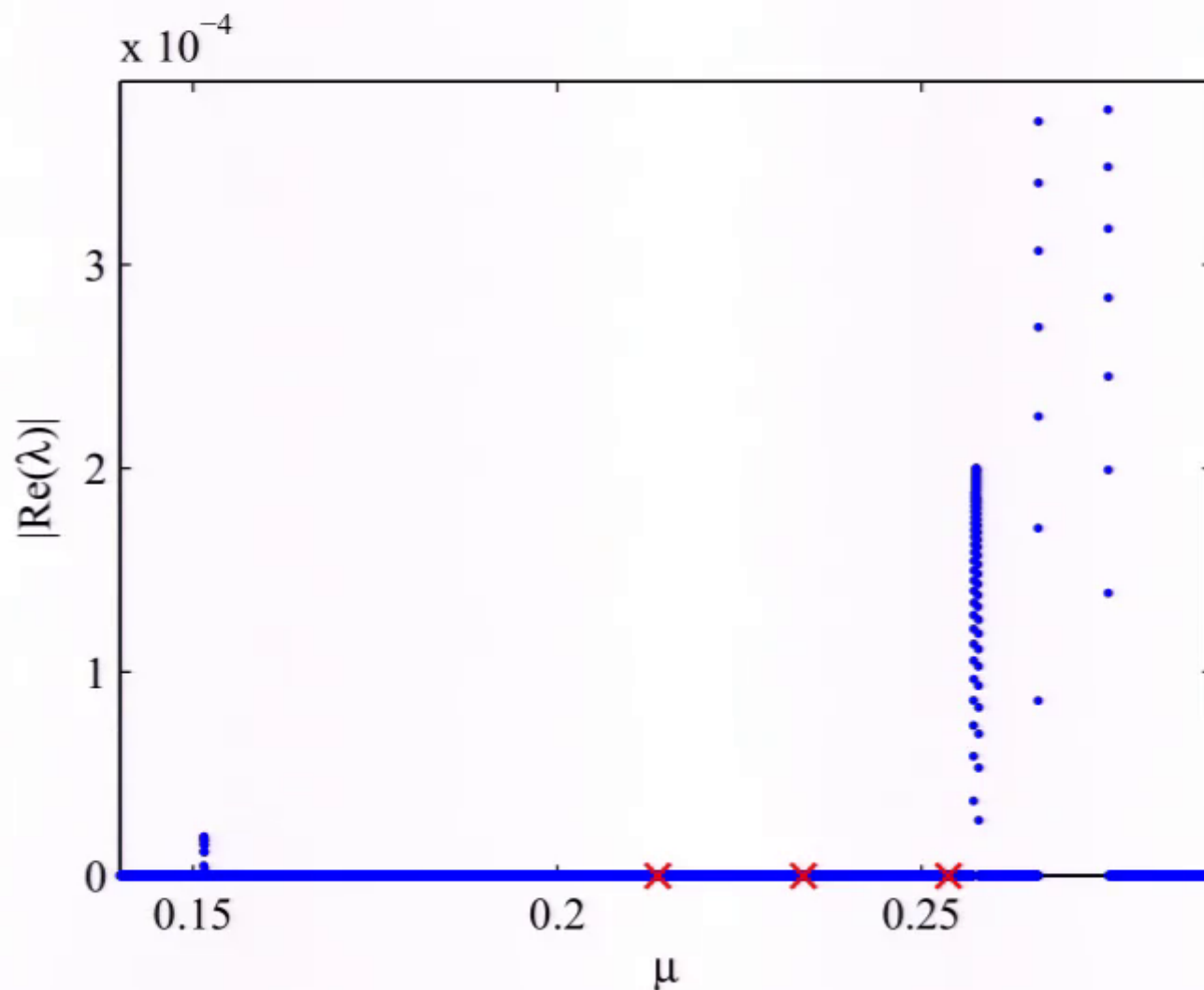
To understand the spectrum of the linearization use a Bloch-wave decomposition of the eigenfunctions:

$$u(z, t) = \psi(z)e^{i\mu z}e^{\lambda t}; \quad \psi(z + 2\pi) = \psi(z), \quad -\frac{1}{2} < \mu < \frac{1}{2}.$$

KG spectrum with wave amplitude  $\sim 0.16$ :



KG spectrum as a function of Bloch-wave parameter  $\mu$  with wave amplitude  $\sim 0.11$ :



In addition to the unstable spectra emanating from the origin, there are three bubbles of unstable spectra.

### Big question

Can an upper bound be placed on the total number of bubbles of unstable spectra for (at least) small waves? If so, under what conditions?

Miller/Marangell showed that superliminal waves have an *infinite* number of instability bubbles if the wave corresponds to an infinite-gap potential of a corresponding Hill's equation.

Under reasonable assumptions the spectrum of the polynomial operator,

$$\mathcal{P}_2(\lambda) = \mathcal{A}_0 + \lambda\mathcal{A}_1 + \lambda^2\mathcal{I},$$

satisfies:

- 1 the polynomial eigenvalues have the Hamiltonian spectral symmetry  $\{\lambda, -\bar{\lambda}\}$
- 2 point spectra only
- 3 each polynomial eigenvalue has finite multiplicity
- 4 the only possible accumulation point is infinity.

For our purposes, the coefficients are smooth functions of the Bloch-wave parameter,  $\mu$ .

We are first concerned with determining the number of (potentially) unstable polynomial eigenvalues in terms of the coefficient operators. Set

- 1  $k_r$ : total number of real positive polynomial eigenvalues
- 2  $k_c$ : total number of complex polynomial eigenvalues with positive real part and nonzero imaginary part.

The negative Krein index of a purely imaginary polynomial eigenvalue,  $\lambda_0$ , with associated eigenspace,  $\mathbb{E}_{\lambda_0}$ , is

$$k_i^-(\lambda_0) = n \left( -\lambda_0 \mathcal{P}'_2(\lambda_0) |_{\mathbb{E}_{\lambda_0}} \right).$$

The total negative Krein index,  $k_i^-$ , is the sum of  $k_i^-(\lambda_0)$  for each polynomial eigenvalue  $\lambda_0 \in i\mathbb{R}$ .

The Hamiltonian-Krein index is  $K_{\text{Ham}} = k_r + k_c + k_i^-$ .

The Hamiltonian-Krein index is  $K_{\text{Ham}} = k_r + k_c + k_i^-$ .

For  $\mathcal{P}_2(\lambda) = \mathcal{A}_0 + \lambda\mathcal{A}_1 + \lambda^2\mathcal{I}$  with  $\mathcal{A}_0$  invertible,

$$K_{\text{Ham}} = n(\mathcal{A}_0)$$

(Pelinovsky et al., Kapitula et al., Grillakis et al., others).

Under suitable (quite generic) assumptions, an underlying wave is orbitally stable if  $K_{\text{Ham}} = 0$  (Grillakis et al.).



We wish to graphically locate those purely imaginary polynomial eigenvalues with negative Krein signature. The Krein matrix for a polynomial operator,

$$\mathcal{P}_2(\lambda) = \mathcal{A}_0 + \lambda \mathcal{A}_1 + \lambda^2 \mathcal{I},$$

for  $\lambda = iz$  is:

$$\mathbf{K}_S(z) = -z \mathcal{P}_2(iz)|_S \cdots \\ \cdots + z \mathcal{P}_2(iz) P_{S^\perp} (P_{S^\perp} \mathcal{P}_2(iz) P_{S^\perp})^{-1} P_{S^\perp} \mathcal{P}_2(iz)|_S.$$

Here  $S$  is the finite-dimensional negative subspace for  $\mathcal{A}_0$  (or an approximation). For the problems at hand,

$$\dim[S] = K_{\text{Ham}}.$$

Polynomial eigenvalues are those values for which the Krein matrix is singular.

The Krein eigenvalues,  $r_j(z)$ , for  $j = 1, \dots, \dim[S]$ , are the eigenvalues of the Krein matrix. The Krein eigenvalues satisfy:

- if  $\lambda = iz_0$  is an eigenvalue, then for some  $j$ ,  $r_j(z_0) = 0$  with

$$r_j'(z_0) \begin{cases} < 0, & k_i^-(iz_0) = 1 \\ > 0, & k_i^-(iz_0) = 0 \end{cases}$$

- $r_j'(z) > 0$  near poles (eigenvalues of  $P_{S^\perp} \mathcal{P}_2(iz) P_{S^\perp}$ )
- each pole of the Krein matrix is a removable singularity for all but one of the Krein eigenvalues.

We wish to graphically locate those purely imaginary polynomial eigenvalues with negative Krein signature. The Krein matrix for a polynomial operator,

$$\mathcal{P}_2(\lambda) = \mathcal{A}_0 + \lambda \mathcal{A}_1 + \lambda^2 \mathcal{I},$$

for  $\lambda = iz$  is:

$$\mathbf{K}_S(z) = -z \mathcal{P}_2(iz)|_S \cdots \\ \cdots + z \mathcal{P}_2(iz) P_{S^\perp} (P_{S^\perp} \mathcal{P}_2(iz) P_{S^\perp})^{-1} P_{S^\perp} \mathcal{P}_2(iz)|_S.$$

Here  $S$  is the finite-dimensional negative subspace for  $\mathcal{A}_0$  (or an approximation). For the problems at hand,

$$\dim[S] = K_{\text{Ham}}.$$

Polynomial eigenvalues are those values for which the Krein matrix is singular.

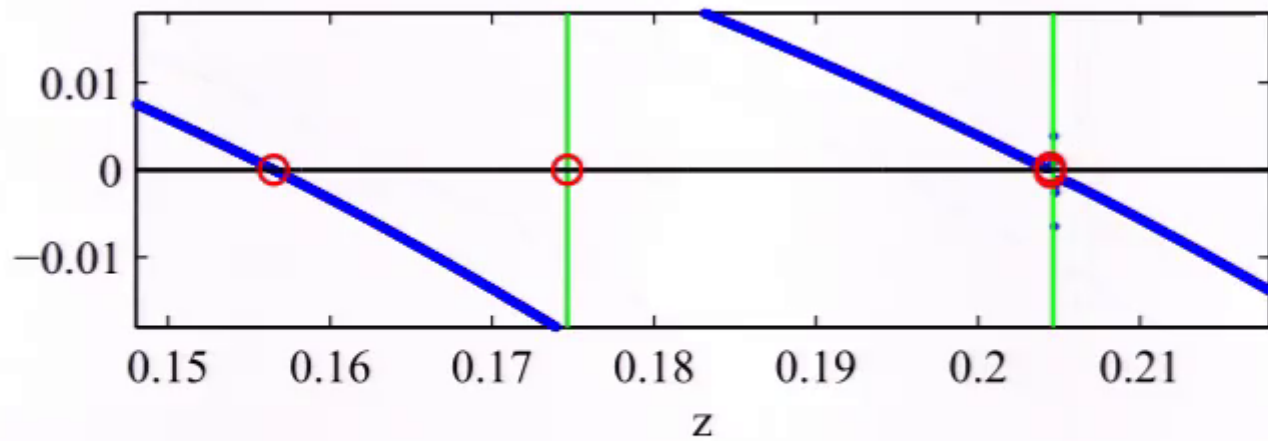
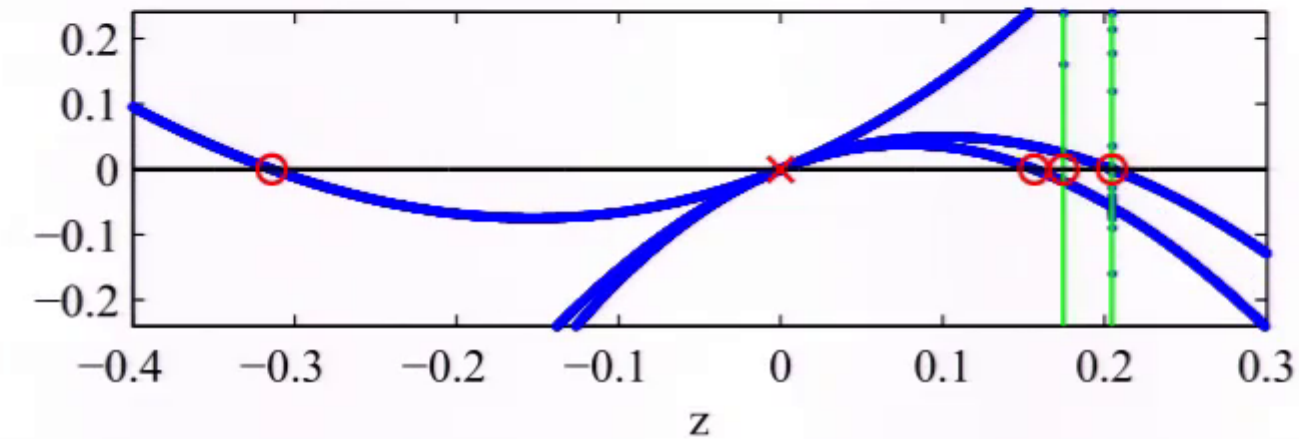
The Krein eigenvalues,  $r_j(z)$ , for  $j = 1, \dots, \dim[S]$ , are the eigenvalues of the Krein matrix. The Krein eigenvalues satisfy:

- if  $\lambda = iz_0$  is an eigenvalue, then for some  $j$ ,  $r_j(z_0) = 0$  with

$$r_j'(z_0) \begin{cases} < 0, & k_i^-(iz_0) = 1 \\ > 0, & k_i^-(iz_0) = 0 \end{cases}$$

- $r_j'(z) > 0$  near poles (eigenvalues of  $P_{S^\perp} \mathcal{P}_2(iz) P_{S^\perp}$ )
- each pole of the Krein matrix is a removable singularity for all but one of the Krein eigenvalues.

Krein eigenvalue for small waves ( $\sim 0.05$ ) to KG:



$$\mu = 0.2435$$

## Conclusion

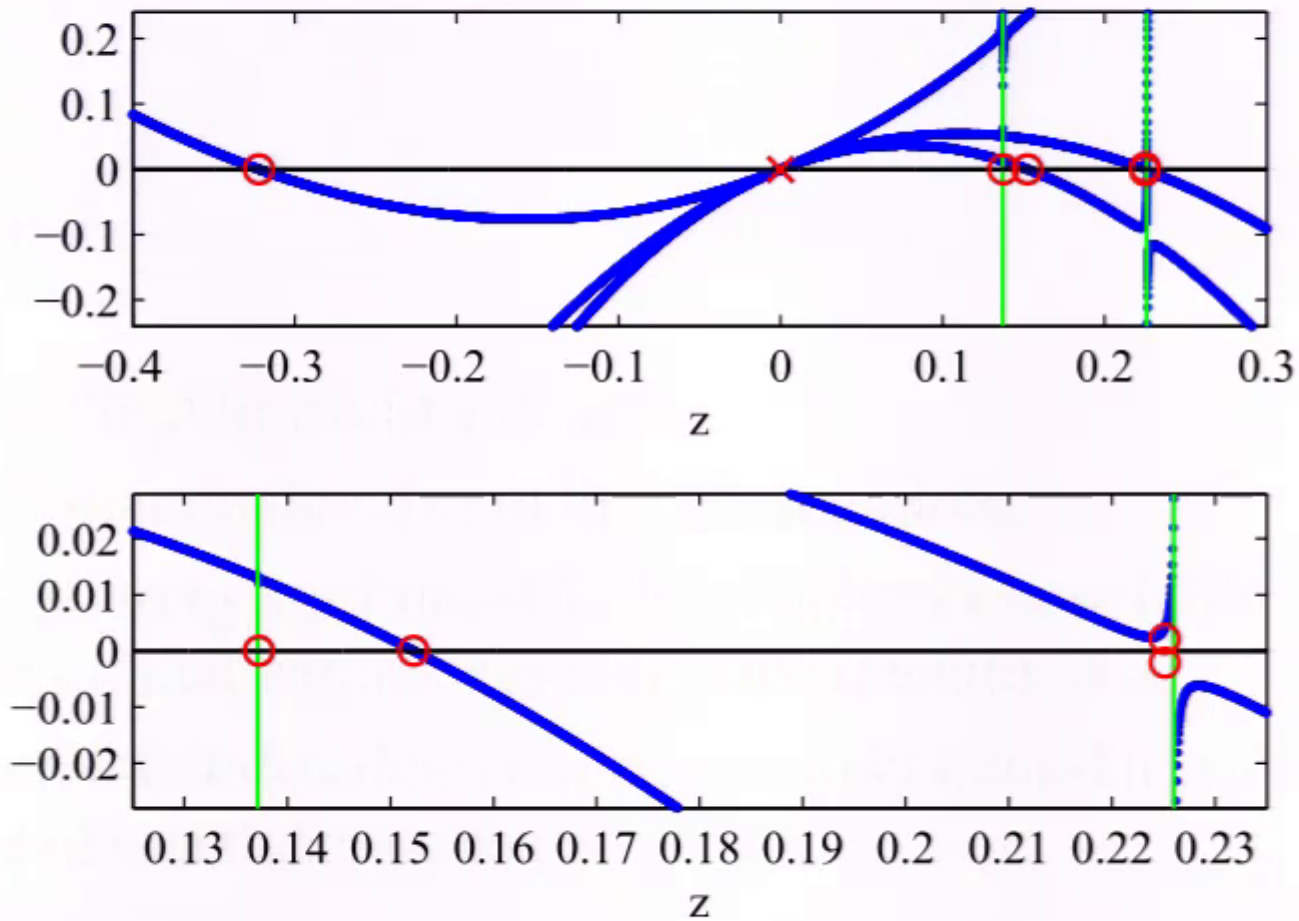
The number of instability bubbles for small waves is bounded above by the number of (Krein eigenvalue) zero/pole collisions for the unperturbed problem.

The Hamiltonian-Krein index:

- 1 determines the size of the Krein matrix
- 2 gives an upper bound for the number of unstable polynomial eigenvalues for a *fixed* value of  $\mu$ .

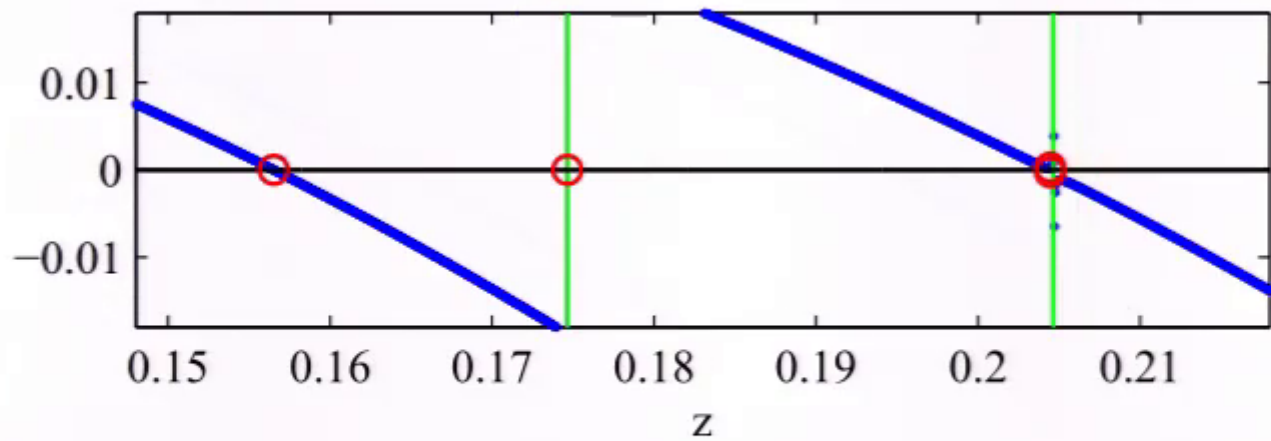
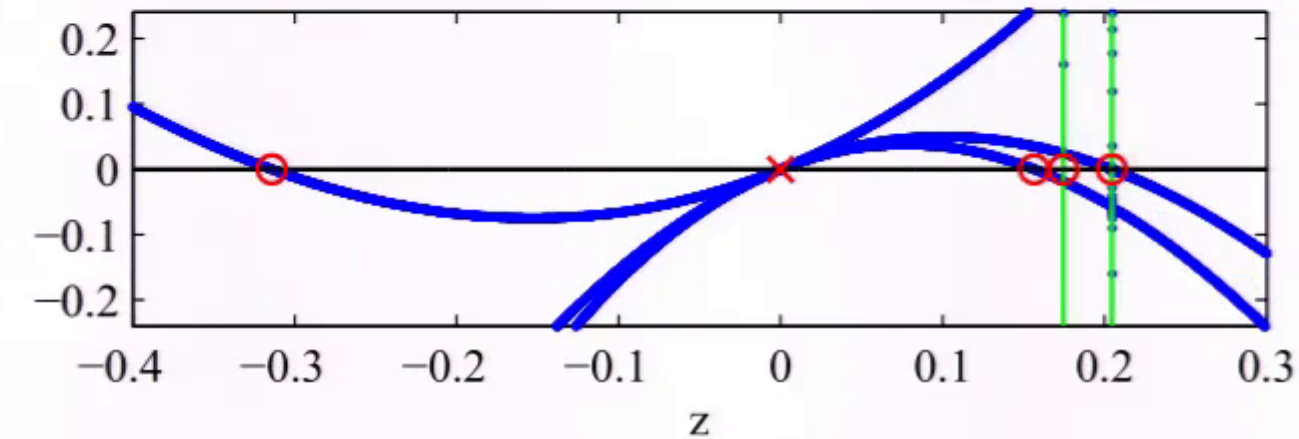
However, the index does not necessarily bound the total number of instability bubbles.

Krein eigenvalue for small waves ( $\sim 0.11$ ) to KG:



$$\mu = 0.2725$$

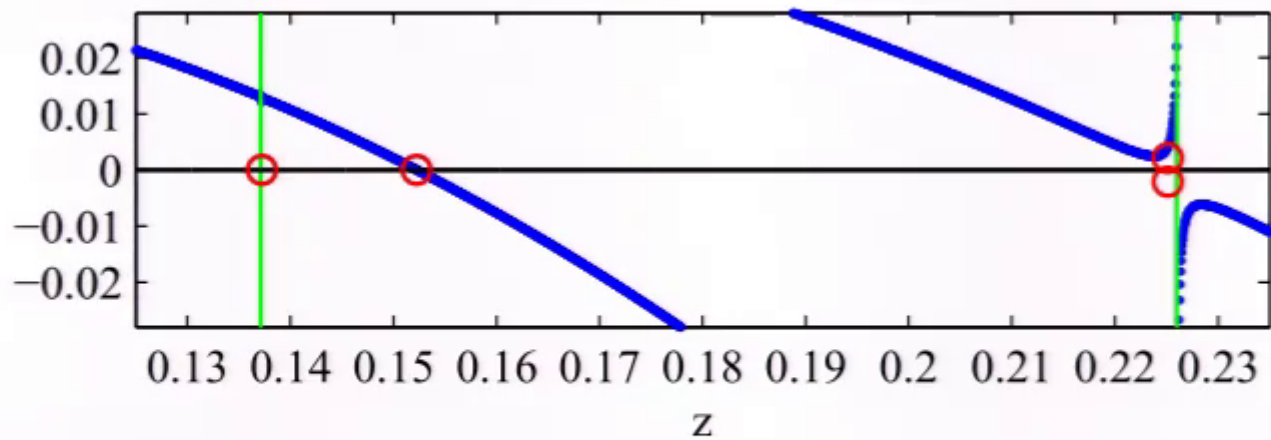
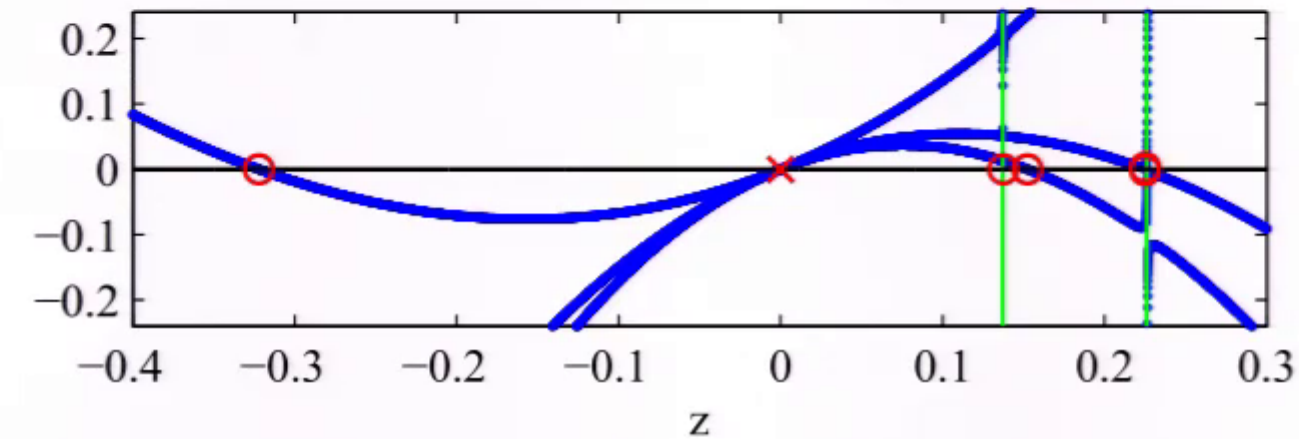
Krein eigenvalue for small waves ( $\sim 0.05$ ) to KG:



$$\mu = 0.2435$$



Krein eigenvalue for small waves ( $\sim 0.11$ ) to KG:



$$\mu = 0.2725$$

## Conclusion

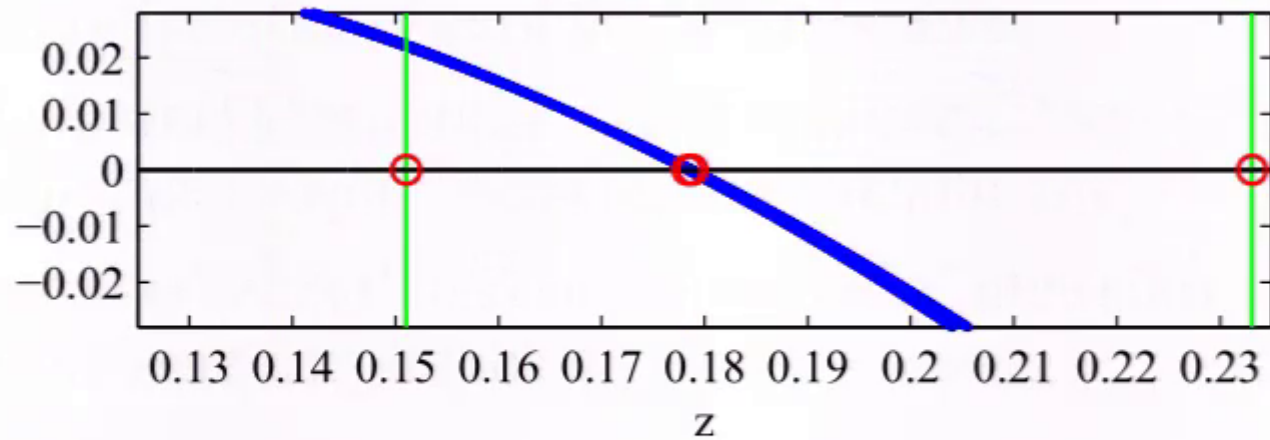
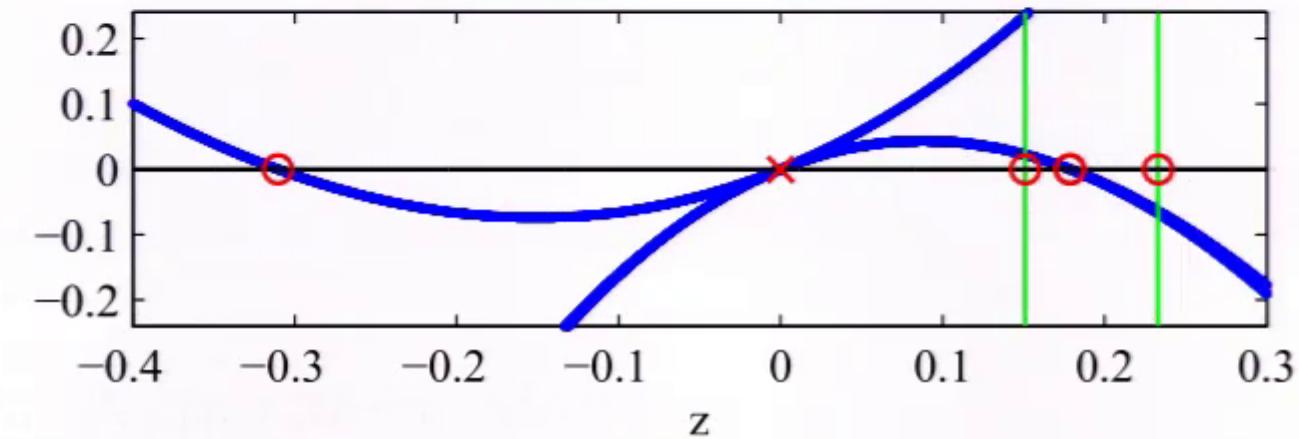
The number of instability bubbles for small waves is bounded above by the number of (Krein eigenvalue) zero/pole collisions for the unperturbed problem.

The Hamiltonian-Krein index:

- 1 determines the size of the Krein matrix
- 2 gives an upper bound for the number of unstable polynomial eigenvalues for a *fixed* value of  $\mu$ .

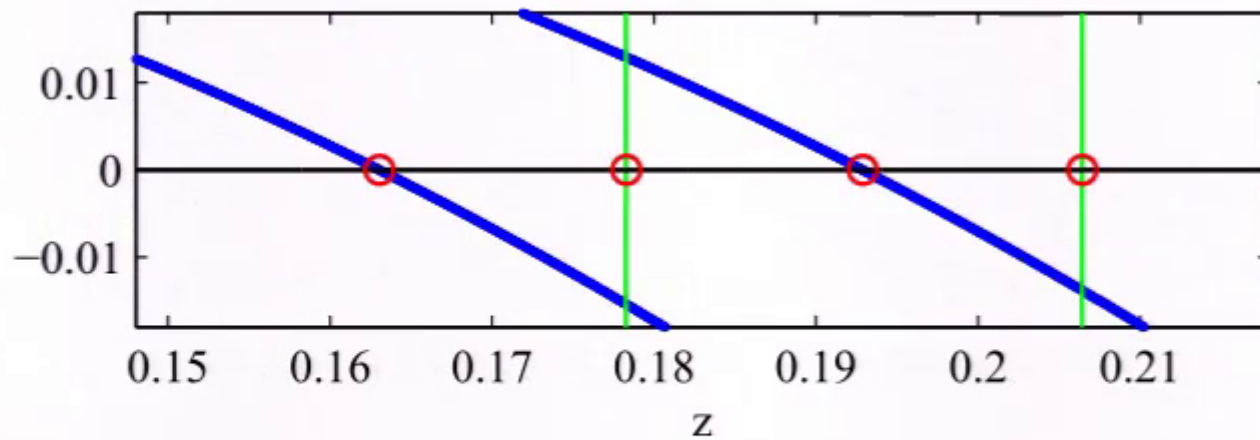
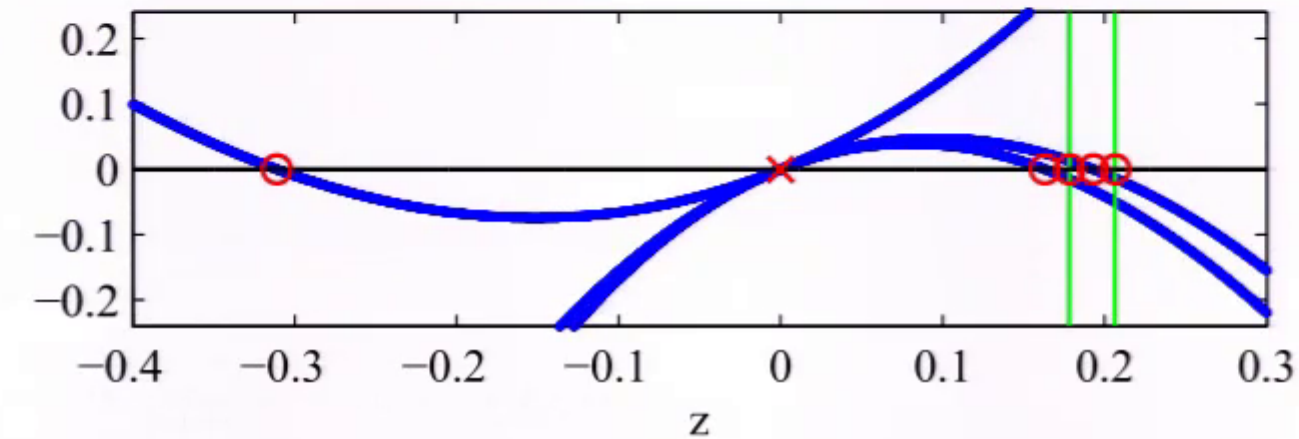
However, the index does not necessarily bound the total number of instability bubbles.

Krein eigenvalue for unperturbed KG:



$$\mu = 0.2725$$

Krein eigenvalue for unperturbed KG ( $K_{\text{Ham}} = 3$ ):



$$\mu = 0.2435$$

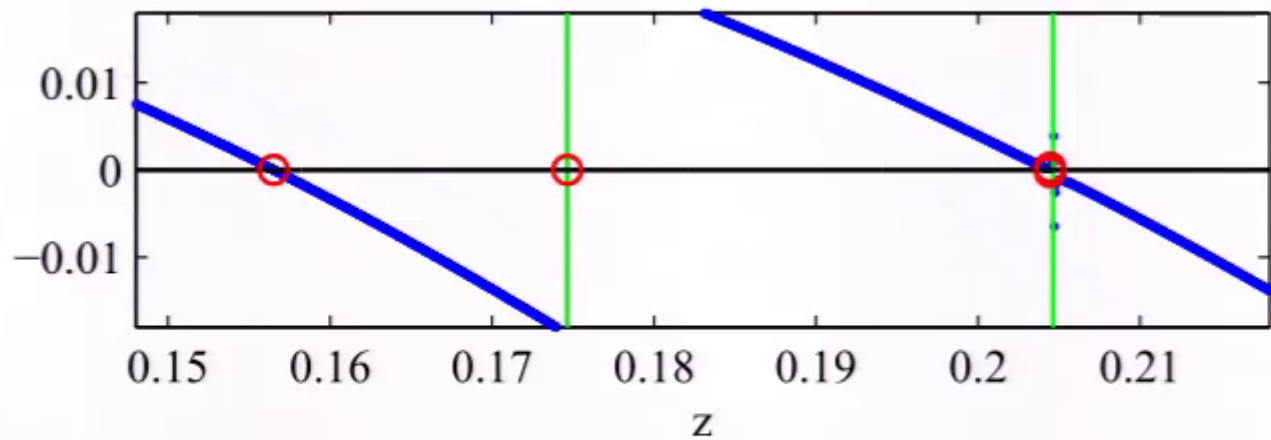
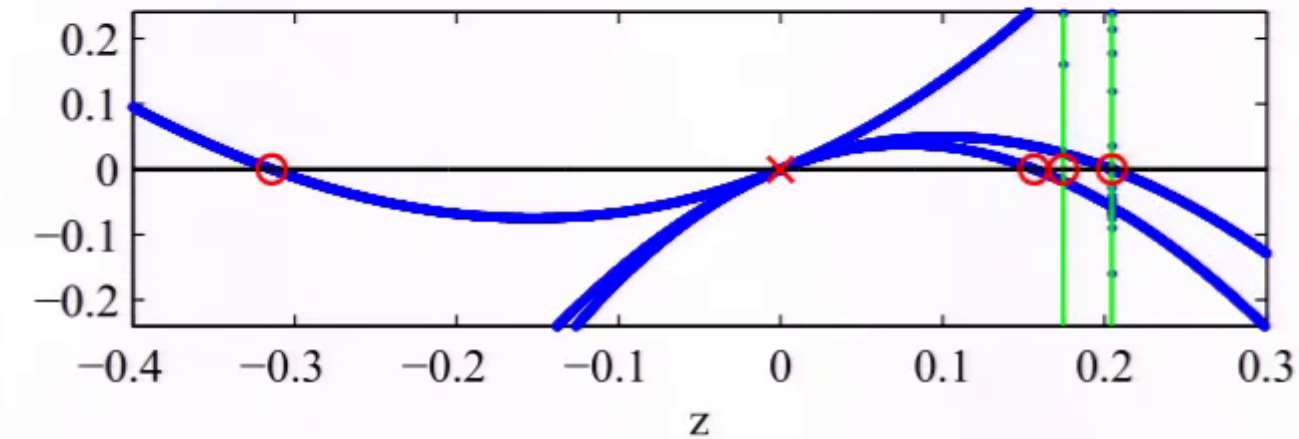
The Krein eigenvalues,  $r_j(z)$ , for  $j = 1, \dots, \dim[S]$ , are the eigenvalues of the Krein matrix. The Krein eigenvalues satisfy:

- if  $\lambda = iz_0$  is an eigenvalue, then for some  $j$ ,  $r_j(z_0) = 0$  with

$$r_j'(z_0) \begin{cases} < 0, & k_i^-(iz_0) = 1 \\ > 0, & k_i^-(iz_0) = 0 \end{cases}$$

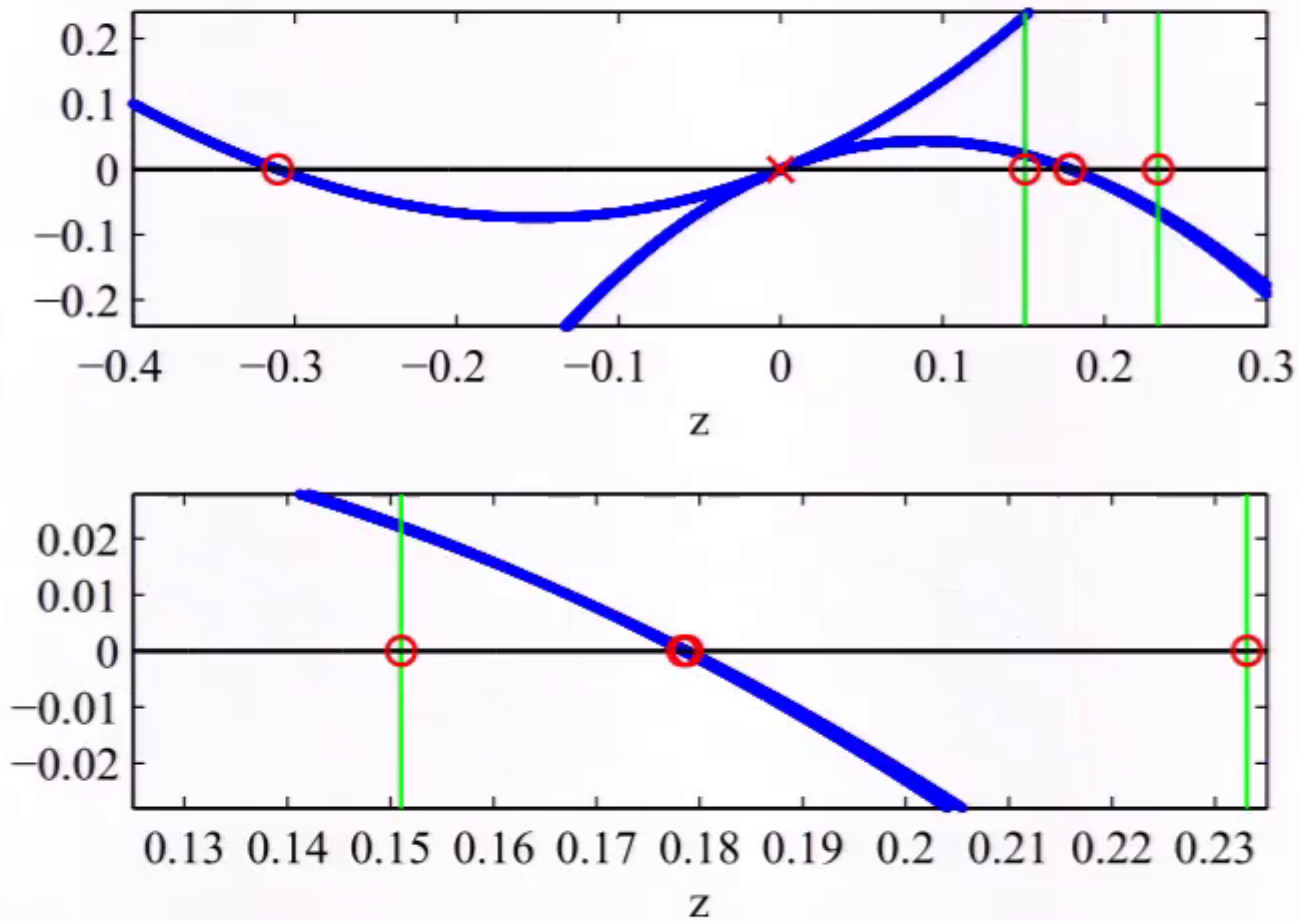
- $r_j'(z) > 0$  near poles (eigenvalues of  $P_{S^\perp} \mathcal{P}_2(iz) P_{S^\perp}$ )
- each pole of the Krein matrix is a removable singularity for all but one of the Krein eigenvalues.

Krein eigenvalue for small waves ( $\sim 0.05$ ) to KG:



$$\mu = 0.2435$$

Krein eigenvalue for unperturbed KG:



$$\mu = 0.2725$$