

The Multiscale Perturbation Method for Elliptic Equations

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Two Phase flow problem

Consider a two-phase oil-water flow problem:

$$-\lambda(s_w)\kappa\nabla p = \mathbf{u}$$

$$\nabla \cdot \mathbf{u} = q$$

$$\phi \frac{\partial s_w}{\partial t} + \nabla \cdot \left(\frac{\lambda_w(s_w)}{\lambda(s_w)} \mathbf{u} \right) = 0$$

- $\mathbf{u}(\mathbf{x}, t)$, $p(\mathbf{x}, t)$ – velocity and pressure fields, $s_w(\mathbf{x}, t)$ – water saturation, $q(\mathbf{x}, t)$ – source term, $\phi(\mathbf{x})$, $\kappa(\mathbf{x})$ – porosity, absolute permeability, $\lambda(s_w) = \lambda_w(s_w) + \lambda_o(s_w)$ – phase mobilities, no capillary pressure or gravity effects. Boundary and initial conditions have to be provided.

Operator Splitting

- One of the popular techniques to solve this problem is **Operator Splitting** ($s_w = s$):

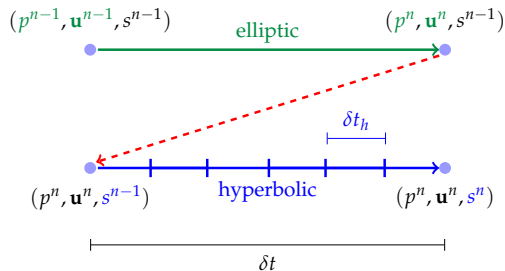


Figure: The operator splitting technique.

Operator Splitting: References

- **Two-phase:** J. Douglas, Jr., F. Furtado, F. P. (1997), On the numerical simulation of waterflooding of heterogeneous petroleum reservoirs, Computational Geosciences, v. 1, # 2, pp. 155-190.
- **Three-phase:** E. Abreu, J. Douglas Jr., F. Furtado and F. P. (2008), Operator Splitting Based on Physics for Flow in Porous Media, International Journal of Computational Science, v. 2, 3, pp. 315-335.
- **Compositional:** M. Akbarabadi, M. Borges, A. Jan, F. P. and M. Piri (2015), A Bayesian Framework for the Validation of Models for Subsurface Flows: Synthetic Experiments, Computational Geosciences, Volume 19, Issue 6, pp. 1231-1250.

Two Phase flow contd.

- Computational efficiency: elliptic time step $>$ saturation time step
- **Key point:** total velocity $\mathbf{u}^{n-1} \approx \mathbf{u}^n$ because

$$\lambda(s_w^{n-1})\kappa(\mathbf{x}) \approx \lambda(s_w^n)\kappa(\mathbf{x})$$

- $\frac{\kappa(\mathbf{x})_{MAX}}{\kappa(\mathbf{x})_{MIN}} \approx 10^\alpha, \quad \alpha = 5, 6, 7 \dots$
- $\frac{\lambda(s_w)_{MAX}}{\lambda(s_w)_{MIN}} \approx 4$

The velocity-pressure problem

Conclusion: The two-phase flow problem discussed above requires the solution of **several** problems of the following form

Find the velocity \mathbf{u} and pressure p fields satisfying

$$\begin{aligned}
 \mathbf{u} &= -K\nabla p && \text{in } \Omega \\
 \nabla \cdot \mathbf{u} &= q && \text{in } \Omega \\
 p &= g^p && \mathbf{x} \in \partial\Omega_p \\
 \mathbf{u} \cdot \mathbf{n} &= g^u && \mathbf{x} \in \partial\Omega_u
 \end{aligned} \tag{1}$$

where $K(\mathbf{x}) = \lambda(s_w(\mathbf{x}))\kappa(\mathbf{x})$ (a scalar in our discussion), q is the source term, g^p is the Dirichlet boundary data and g^u is the Neumann boundary data.

The Multiscale Perturbation Method (MPM)

MPM combines:

- The Multiscale Mixed Method (MuMM)
- Classical Perturbation Theory
- HPC: Multi-core devices

References for the MuMM

- **Original DDM:** J. Douglas, Jr., P.J. Paes Leme, J.E. Roberts, and Junping Wang (1993), A parallel iterative procedure applicable to the approximate solution of second order partial differential equations by mixed finite element methods, Numer. Math. 65, 95-108.
- **MuMM:** A. Francisco, V. Ginting, F. P., J. Rigelo (2014), Design and Implementation of a Multiscale Mixed Method for Porous Media Flows, Mathematics and Computers in Simulation, 99, pp. 125-138.

MuMM : Formulation

- Three spatial scales enter in the definition of the MuMM:
 $h \leq \bar{H} \leq H$

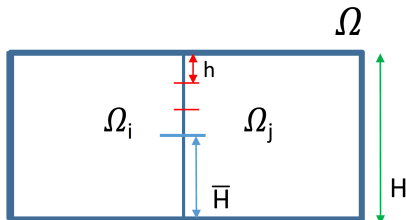


Figure: The three spatial scales for the definition of MuMM and MPM.

The MuMM (cont.)

- The solution of a local Robin BVP (with given Dirichlet and Neumann data on $\partial\Omega$) defined in Ω_i with Robin conditions (r_{ij}^ℓ) , $\ell = 1, 2, \dots, N$ is, obtained as

$$\{\mathbf{u}_h^i, p_h^i\} = \sum_{\Gamma_{ij} \subset \partial\Omega_i} \sum_{\ell=1}^N r_{ij}^\ell \psi_{ij}^\ell + \phi_i \quad (2)$$

- The global solution of MuMM is based on a red-black iteration over subdomains.

The MuMM (cont.)

- The divergence-free **Multiscale Basis Functions** (MSBF) $\Psi_{ij}^\ell = \{\psi_{ij}^\ell, \phi_{ij}^\ell\}$ are defined as the solution of the following Robin boundary value problem (BVP)

$$\Psi_{ij}^\ell : \begin{cases} \psi_{ij}^\ell = -K \nabla \phi_{ij}^\ell & \text{in } \Omega_i \\ \nabla \cdot \psi_{ij}^\ell = 0 & \text{in } \Omega_i \\ \phi_{ij}^\ell = 0 & \mathbf{x} \in \partial\Omega_p \\ \psi_{ij}^\ell \cdot \mathbf{n}^i = 0 & \mathbf{x} \in \partial\Omega_u \\ -\beta \psi_{ij}^\ell \cdot \mathbf{n}^i + \phi_{ij}^\ell = \mathbf{g}_{ij}^\ell & \mathbf{x} \in \Gamma_{ij} \\ \text{such that } \mathbf{g}_{ij}^\ell = 1 \text{ in } \Gamma_{ij}^\ell \text{ and } 0 \text{ in } \partial\Omega_i \setminus \Gamma_{ij}^\ell \end{cases}$$

The MuMM (cont.)

- Non-trivial Dirichlet/Neumann boundary conditions and source terms are incorporated in the local BVP solutions through an **auxiliary basis function** (that takes trivial Robin conditions) $\Phi_i = \{\chi_i, \varphi_i\}$ defined as

$$\Phi_i : \begin{cases} \chi_i = -K \nabla \varphi_i & \text{in } \Omega_i \\ \nabla \cdot \chi_i = q & \text{in } \Omega_i \\ \varphi_i = g^p & \mathbf{x} \in \partial\Omega_p \\ \chi_i \cdot \mathbf{n}^i = g^u & \mathbf{x} \in \partial\Omega_u \\ -\beta \chi_i \cdot \mathbf{n}^i + \varphi_i = 0 & \mathbf{x} \in \Gamma_{ij} \end{cases}$$

The Multiscale Perturbation Method (MPM)

Consider the global problem (1) along with two fields,
 $K(\mathbf{x}) = K_0(\mathbf{x})$ and $K(\mathbf{x}) = K_t(\mathbf{x})$

- $K_0(\mathbf{x})$ corresponds to $t_0 = 0$
- $K_t(\mathbf{x})$ corresponds to $t_1 > 0$

To identify a **small parameter** that reflects the change in $K(\mathbf{x})$ from t_0 to t_1 , write

$$K_t(\mathbf{x}) = K_0(\mathbf{x}) + \epsilon K_p(\mathbf{x}), \quad \mathbf{x} \in \Omega \quad (3)$$

where $\epsilon = \frac{\|K_t(\mathbf{x}) - K_0(\mathbf{x})\|}{\|K_t(\mathbf{x}) - K_0(\mathbf{x})\|}$ and $K_p(\mathbf{x}) = \frac{K_t(\mathbf{x}) - K_0(\mathbf{x})}{\|K_t(\mathbf{x}) - K_0(\mathbf{x})\|}$;

Here ϵ is a dimensionless parameter with its value < 1 .

The MPM scheme (cont.)

The target global BVP can be written as

$$P_t : \begin{cases} \mathbf{u} = -(K_0 + \epsilon K_p) \nabla p & \text{in } \Omega \\ \nabla \cdot \mathbf{u} = q & \text{in } \Omega \\ p = g^p & \mathbf{x} \in \partial\Omega_p \\ \mathbf{u} \cdot \mathbf{n} = g^u & \mathbf{x} \in \partial\Omega_u \end{cases}$$

- Let $\mathbf{u}_{ref}, p_{ref}$ be a guess to the solution of the target problem, such that,

$$\begin{aligned} \mathbf{u} &= \mathbf{u}_{ref} + \delta \mathbf{u} \\ p &= p_{ref} + \delta p \end{aligned}$$

The MPM scheme contd.

- Considering the expansion,

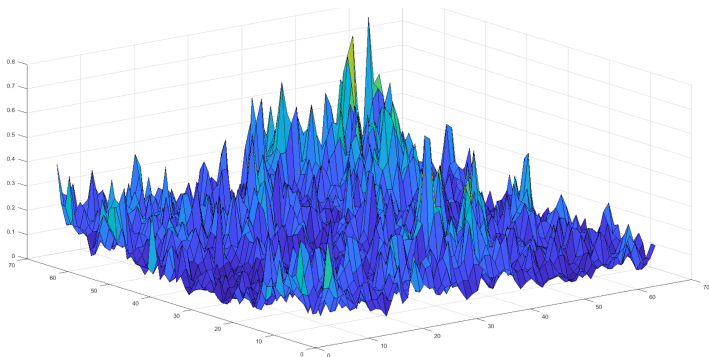
$$\delta \mathbf{u} = \delta \mathbf{u} + \epsilon \delta \mathbf{u} + \epsilon^2 \delta \mathbf{u} + \cdots + \epsilon^r \delta \mathbf{u}$$

$$\delta p = \delta p_0 + \epsilon \delta p_1 + \epsilon^2 \delta p_2 + \cdots + \epsilon^r \delta p_r$$

we derive a hierarchy of second order elliptic equations that refer to K_0 and not K_t .

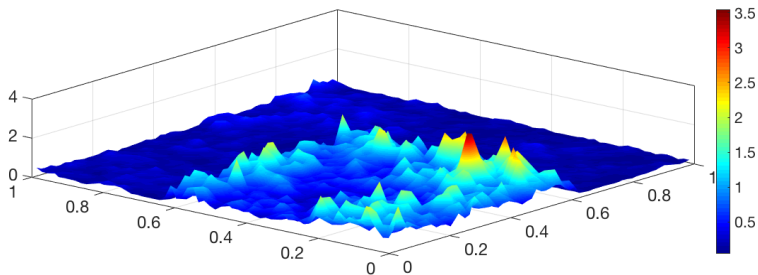
- These BVPs can be efficiently solved using the **MuMM** in multi-core computers.
- The multiscale basis functions computed at time $t = 0$ can be **reused** for all BVPs, once they are defined in terms of $K_0(\mathbf{x})$.

The Model Problem



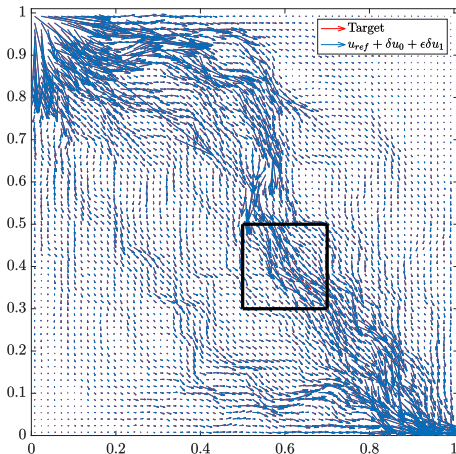
The absolute permeability field. $\frac{\kappa_{max}}{\kappa_{min}} \approx 10^3$

The Model Problem (cont.)



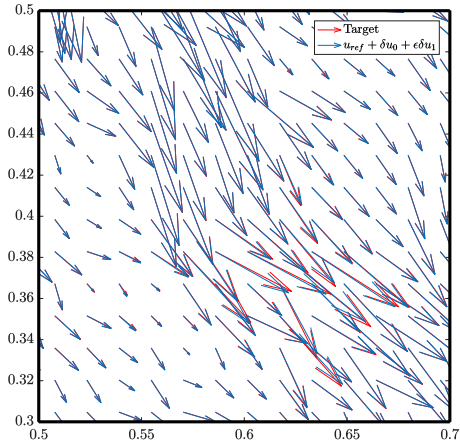
The target permeability field showing the area that has been perturbed from $t = 0$ to $t = t_1$. The region with larger variability mimics an invading water front.

Vector field: $\epsilon = 0.05$



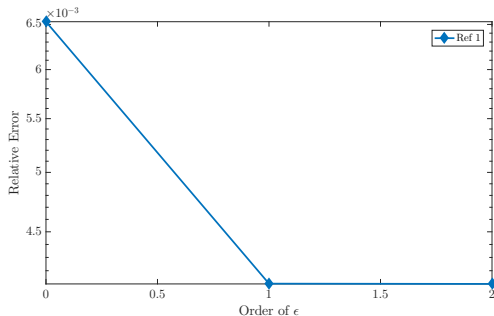
Comparing the fine scale solution with the first order ϵ

Vector field: $\epsilon = 0.05$



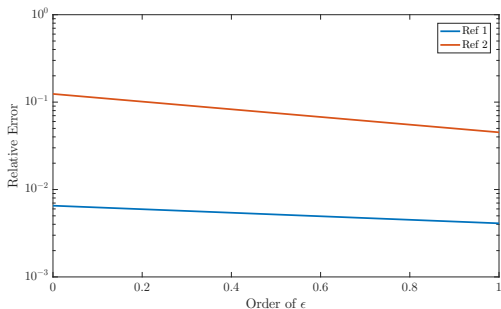
Comparing the fine scale solution with the first order ϵ

Vector field: $\epsilon = 0.05$



Relative Error for flux vs several orders of ϵ

Vector field: $\epsilon = 0.05$



Relative Error for flux vs several orders of ϵ

Conclusion/Future Work

- **MPM** has been introduced and tested on simple problems.
- The solution of a two phase flow problem obtained using techniques like Operator splitting and IMPES (IMplicit Pressure, EXplicit Saturation) ¹ can be obtained much more efficiently without recomputing the basis functions at every time step.

¹P. Jenny and S.H. Lee and H.A. Tchelepi (2006), Adaptive fully implicit multi-scale finite-volume method for multi-phase flow and transport in heterogeneous porous media, Journal of Computational Physics, V. 217, pp 627-641

Thank You.