

Solving Eigenvalue and Linear System Problems for 3D Photonic Device Simulations

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A Celebration in Honor of Dianne P. O'Leary on the Occasion of her Retirement
SIAM Conference on Applied Linear Algebra, Atlanta, USA, 2015/10/29

Collaborators

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National Taiwan University

FAME

Fast Algorithms for Maxwell's Equations

Electromagnetic Waves



● Maxwell's equations

$$\nabla \times E = i\omega B$$

$$\nabla \times H = -i\omega D$$

$$\nabla \cdot B = 0$$

$$\nabla \cdot D = 0$$

E : the electric field

H : magnetic field

ω : frequency

B : magnetic flux density

D : electric flux density

● Constitutive relations

$$B = \mu H + \zeta E$$

$$D = \varepsilon E + \xi H$$

μ : magnetic permeability

ε : electric permittivity

ζ and ξ : magnetoelectric parameters

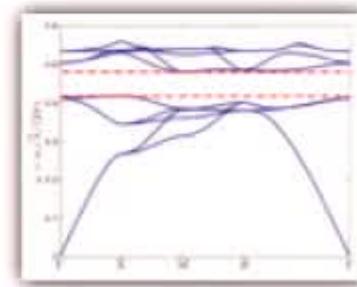
Dielectric Materials (Photonic Crystal)



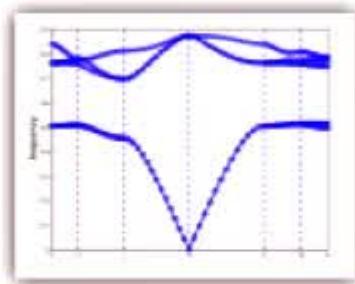
$$\nabla \times \nabla \times E = \mu_0 \omega^2 \epsilon(\mathbf{x}) E$$

$$\epsilon(\mathbf{x}) = \begin{cases} \epsilon_1 & \text{in material 1} \\ \epsilon_2 & \text{in material 2} \end{cases}$$

$$C^* C x = \lambda B x, \quad \lambda = \omega^2$$



Huang/Chang/Huang/Lin/Wang/W (JCP 2010)
Huang/Kuo/W (JSC 2013)



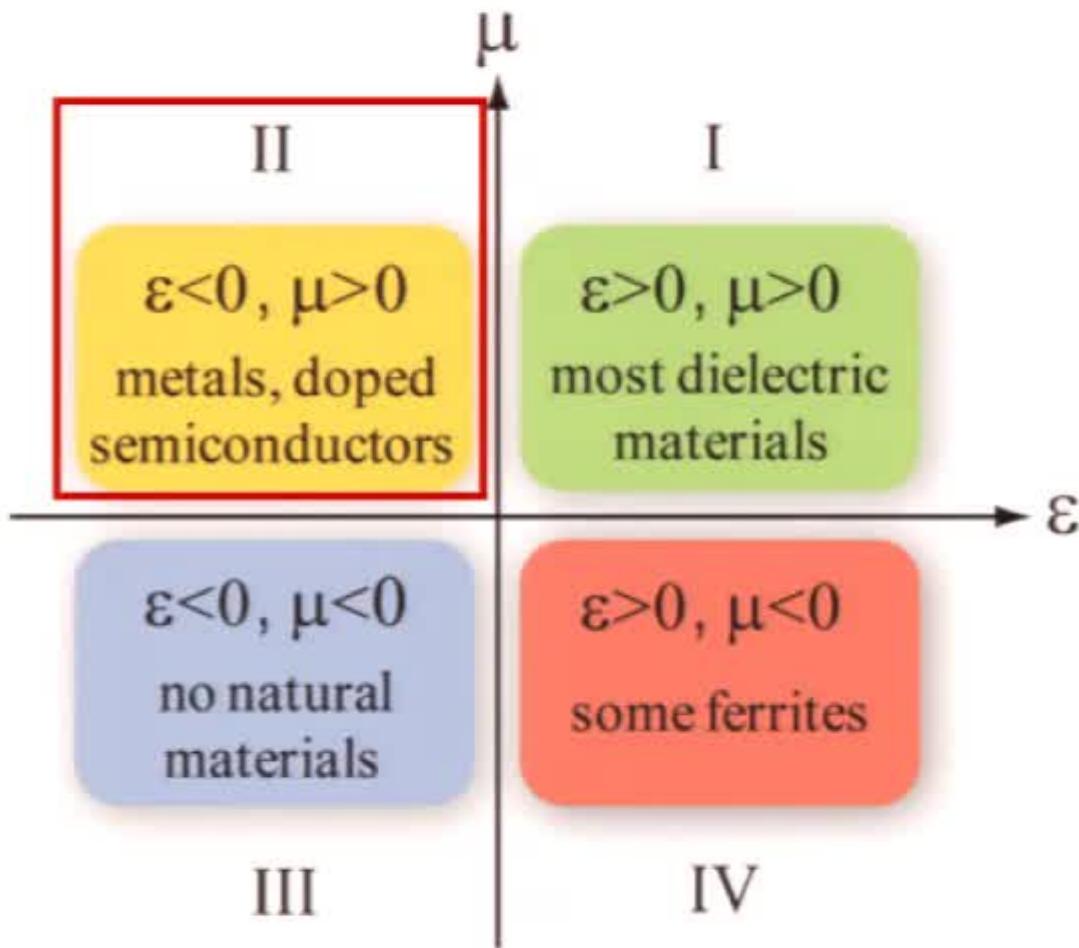
Eigendecompositoin of
discrete double-curl

Huang/Hsieh/Lin/W (SIMAX, 2013)

Dispersive Metallic Materials

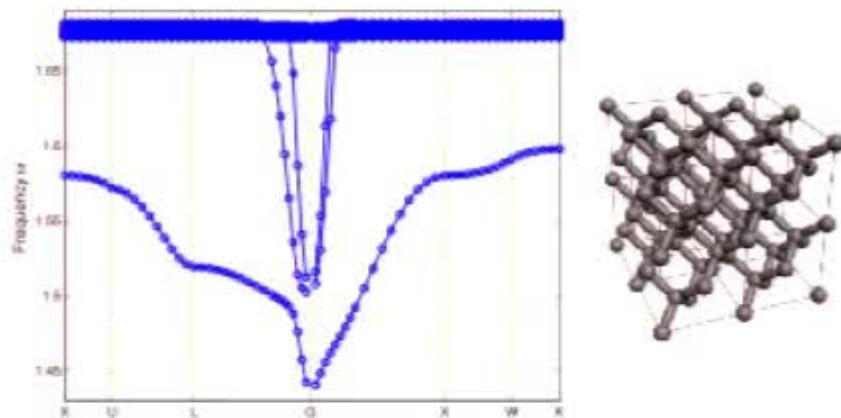


$$\nabla \times \nabla \times E = \mu_0 \omega^2 \epsilon(\mathbf{x}, \omega) E$$



$$\epsilon(\mathbf{x}, \omega) = \begin{cases} \epsilon_1(\omega) & \text{in material 1} \\ \epsilon_2 & \text{in material 2} \end{cases}$$

$$(C^* C)x = \omega^2 B(\omega)x$$

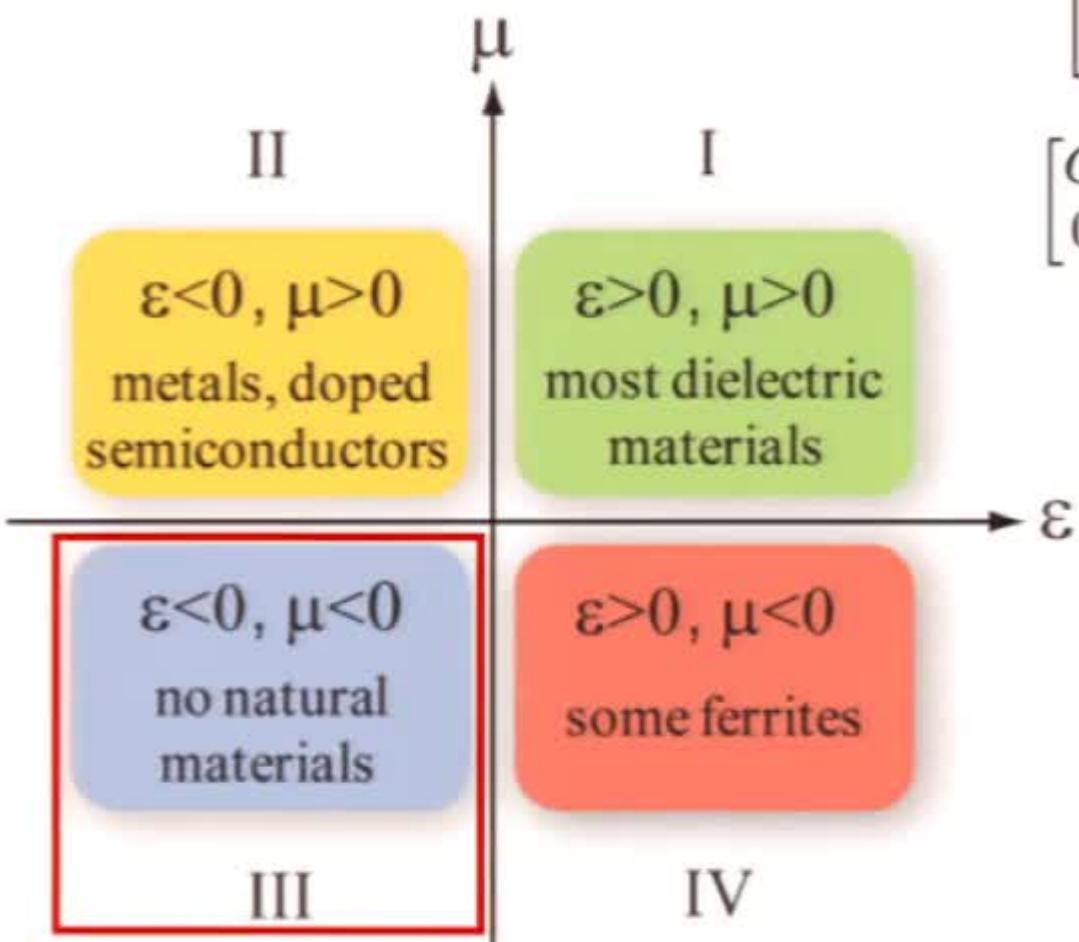


Eignesolver for
cluster eigenvalues

Huang/Lin/W (Preprint, 2015)

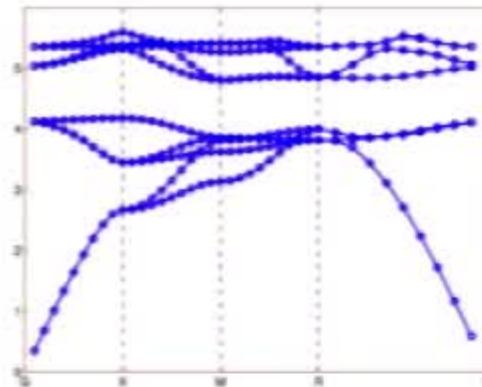


Artificial Complex Media



$$\begin{bmatrix} \nabla \times & 0 \\ 0 & \nabla \times \end{bmatrix} \begin{bmatrix} E \\ H \end{bmatrix} = i\omega \begin{bmatrix} \zeta & \mu \\ -\epsilon & -\xi \end{bmatrix} \begin{bmatrix} E \\ H \end{bmatrix}$$

$$\begin{bmatrix} C & 0 \\ 0 & C^* \end{bmatrix} \begin{bmatrix} E \\ H \end{bmatrix} = \omega \left(i \begin{bmatrix} \zeta_d & \mu_d \\ -\epsilon_d & -\xi_d \end{bmatrix} \right) \begin{bmatrix} E \\ H \end{bmatrix}$$

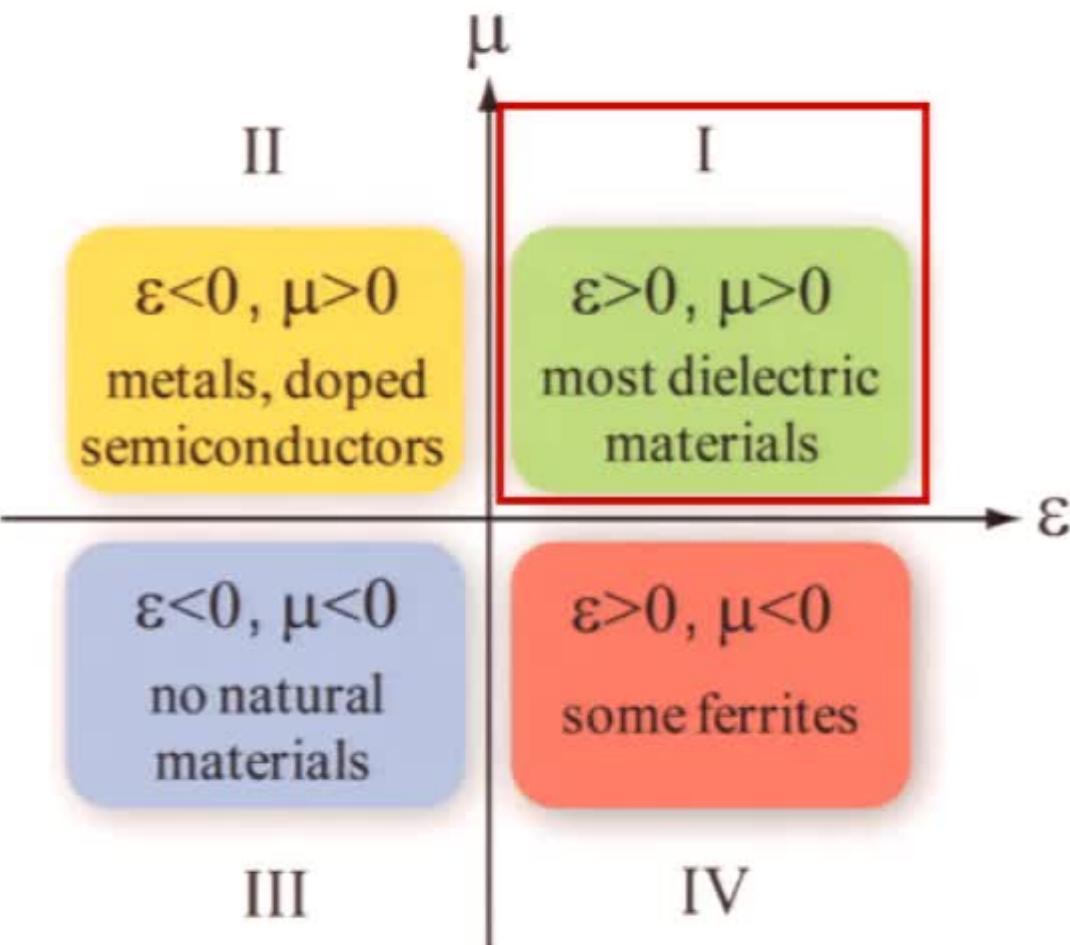


SVD of discrete single-curl

Chern/Hsieh/Huang/Lin/W (SIMAX, 2015)

Photonic Crystals

Dielectric Materials (Photonic Crystal)



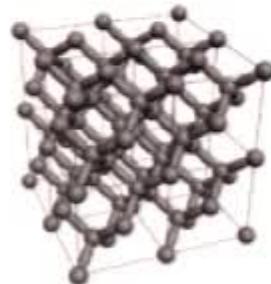
Photonics Crystal



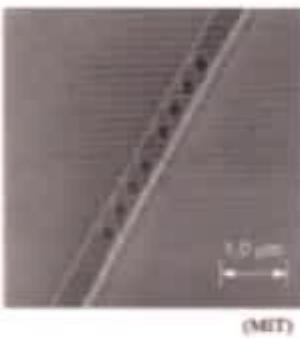
Periodic lattice composed of dielectric or metallic materials



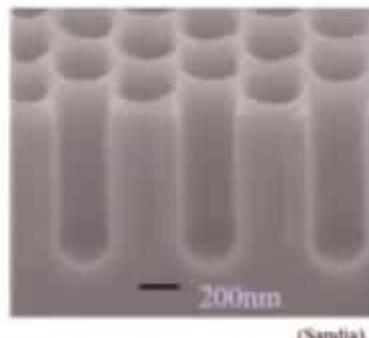
Simple cubic



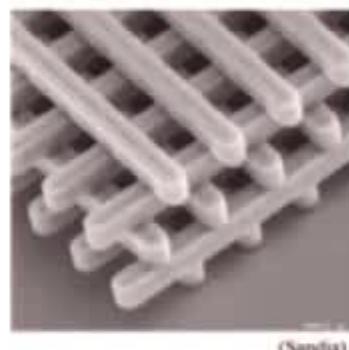
Face centered cubic



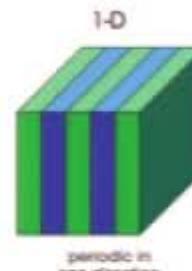
1D hole-array built on
A SOI substrate.



2D hole array built
from GaAs on Al-oxide.

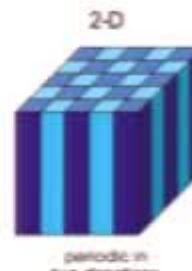


3D diamond lattice built
on a Si substrate.



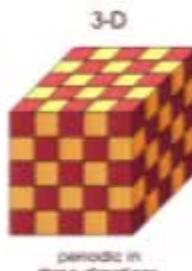
1-D

periodic in
one direction



2-D

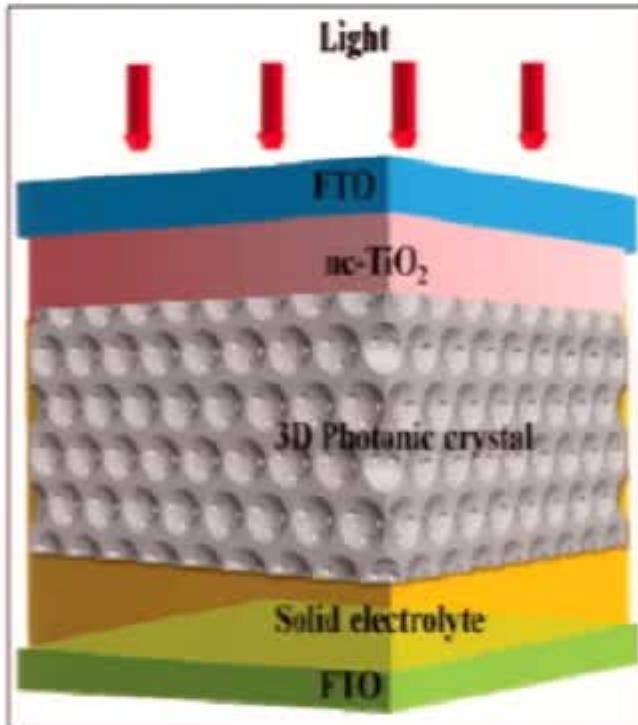
periodic in
two directions



3-D

periodic in
three directions

Control of Lights

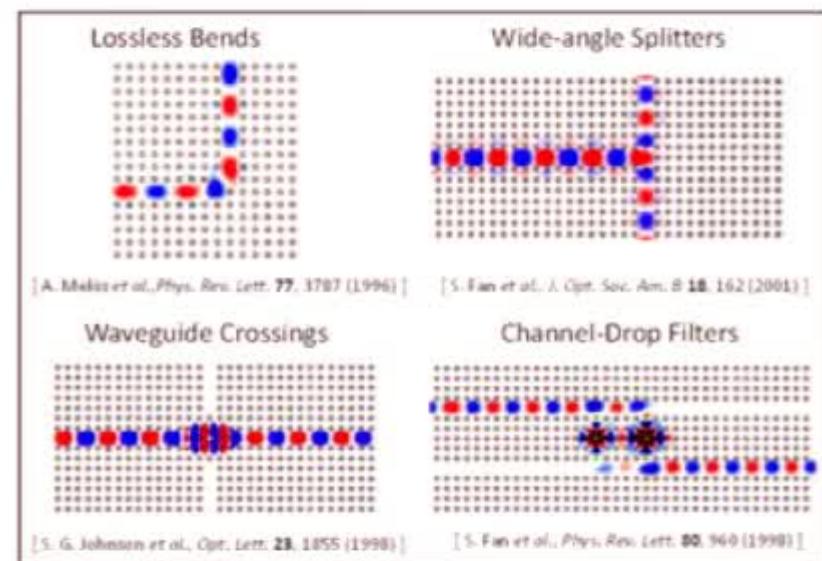
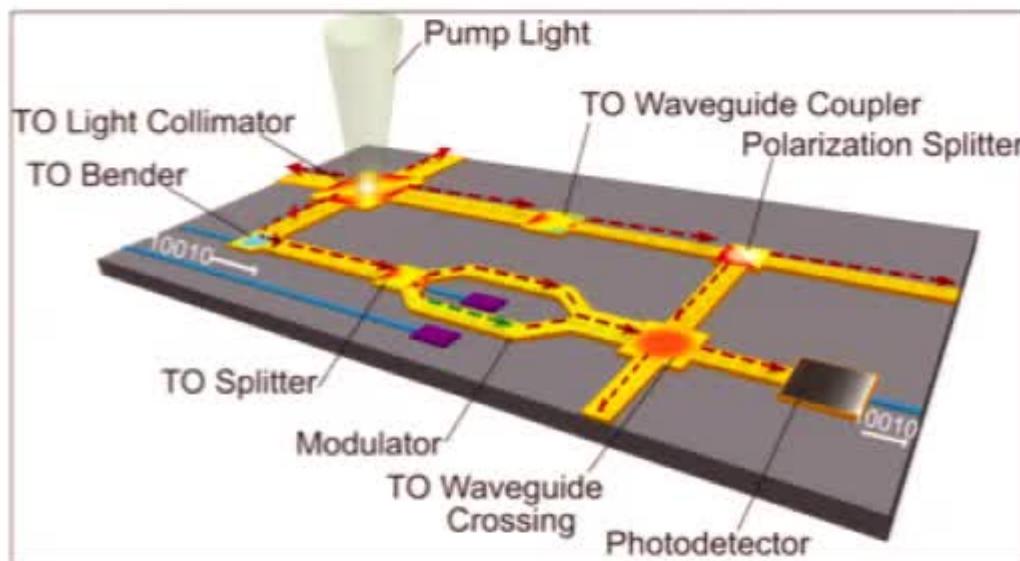


Solar Cells, Hwang et al,
RSC Adv. (2013)



Photonic crystal **light twisters** at Sandia.
<http://goo.gl/b3xtHk>

Control of Lights



A schematic **photonic integrated circuit** for imaging, communications, computing, and sensing. (2012) <http://goo.gl/bUeR0F>

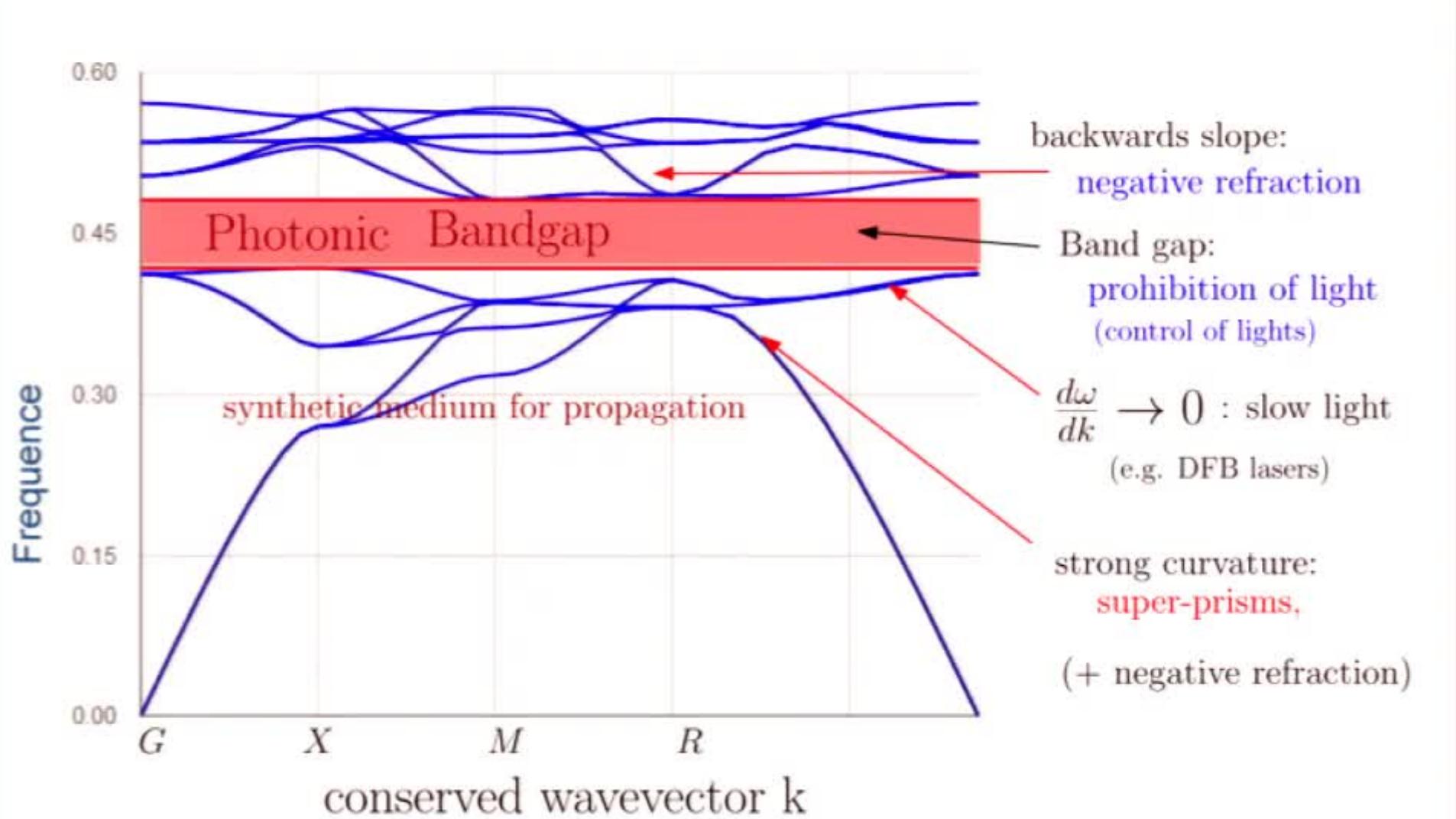
Bandgap



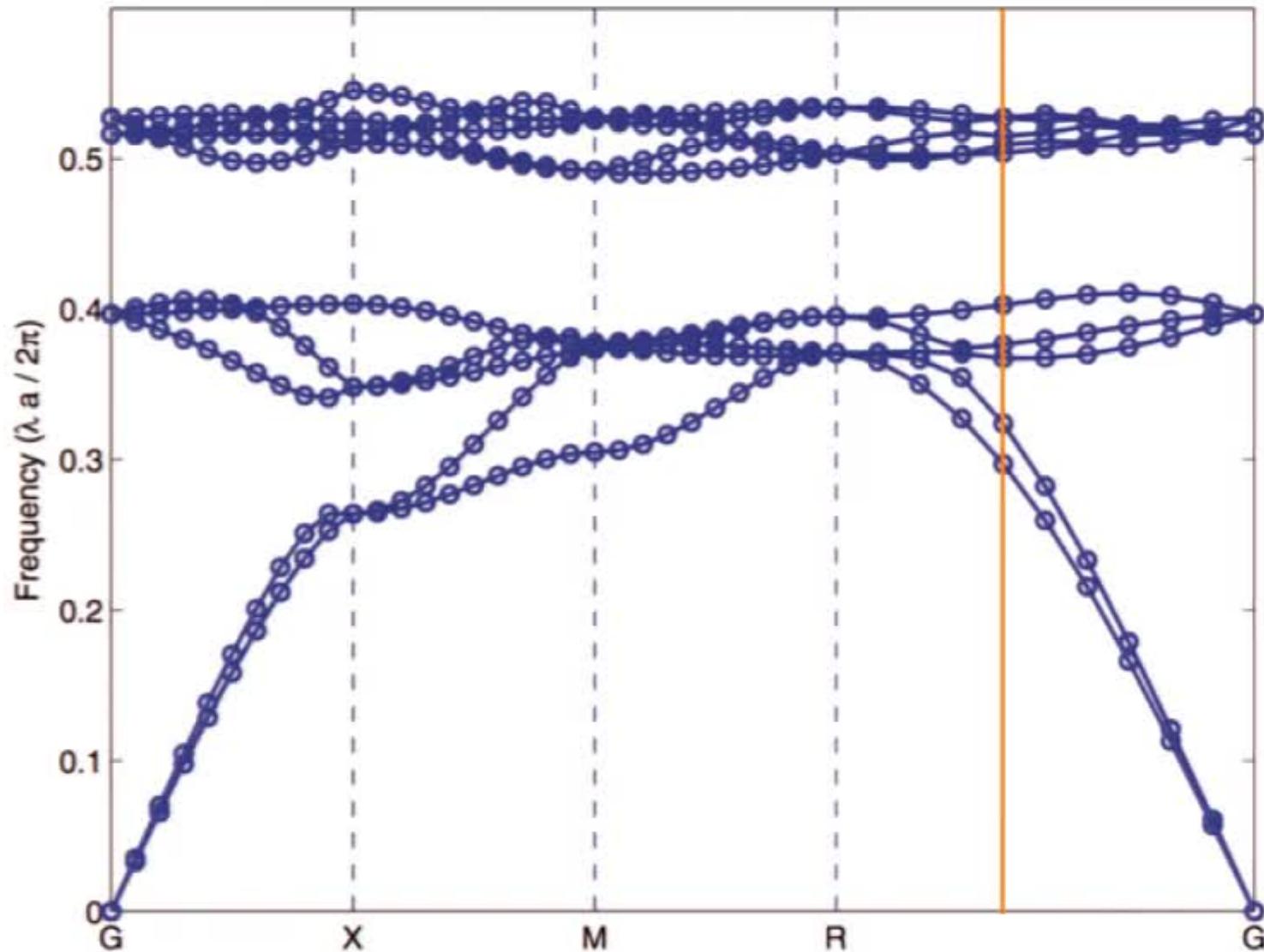
“If only it were possible to make dielectric materials in which electromagnetic waves cannot propagate at certain frequencies all kinds of almost-magical things would be possible.”

John Maddox, Nature, 348 (1990)

Band Structure Diagram



Many Large-scale Interior EVPs



Numerical Challenges



- Yee's scheme discretizes the equation

$$\nabla \times \nabla \times \tilde{E}(\mathbf{x}) = \mu_0 \omega^2 \varepsilon(\mathbf{x}) \tilde{E}(\mathbf{x})$$

to get the generalized eigenvalue problem (GEVP)

$$A\mathbf{x} = \lambda B\mathbf{x}.$$

Huang/Hsieh/Lin/W (Math Comp Model 2013)

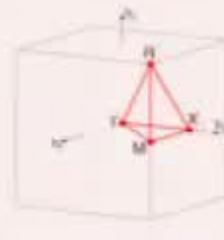
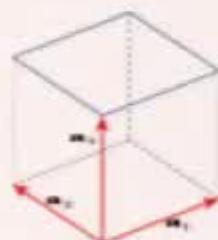
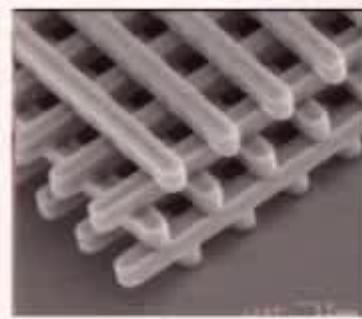
- A : complex Hermitian positive semi-definite
- B : positive diagonal (containing magnetic constant, frequency, material dependent permittivity)
- Dimension: $n=3n_1n_2n_3$
- Need a few of smallest interior positive eigenvalues
- A has big ($n/3$) null space

Quasi-Periodic Conditions

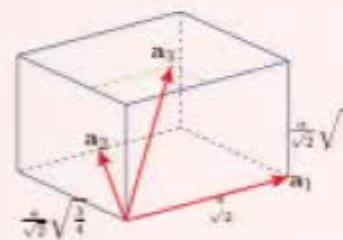


For a Bloch wave vector $2\pi\mathbf{k}$ in the *irreducible Brillouin zone*, the Bloch eigenfunctions E and H satisfy the quasi-periodic condition along the translation vectors \mathbf{a}_l .

$$E(\mathbf{x} + \mathbf{a}_\ell) = e^{i2\pi\mathbf{k}\cdot\mathbf{a}_\ell} E(\mathbf{x}), \quad H(\mathbf{x} + \mathbf{a}_\ell) = e^{i2\pi\mathbf{k}\cdot\mathbf{a}_\ell} H(\mathbf{x}).$$



CUB path: I-X-M-I-R-X-W

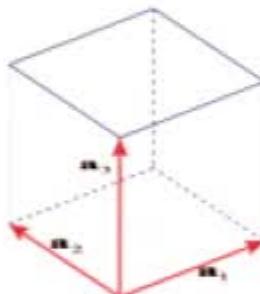


FCC path: I-X-W-X-C-L-U-C-R-S-U-X

Simple Cubic PC

FFT based Preconditioner

Simple Cubic



$$A\mathbf{x} = \lambda B\mathbf{x}$$

$$(A - \tau B)\mathbf{z} = \mathbf{b}$$

Is there a way to find a
preconditioner to cluster
the eigenvalues ?

Simple Cubic: Preconditioner



$$(A - \tau B)\mathbf{z} = \mathbf{b}$$

$$(A - \tau \varepsilon_0 I)\mathbf{z} = \mathbf{b} \quad \text{by averaging } \text{diag}(\mathbf{B})$$

$$\Leftrightarrow (e^{-ik \cdot x} \nabla \times \nabla \times e^{ik \cdot x} - \varepsilon_0 \tau) \mathbf{z} = \mathbf{b}$$

$$\nabla^2 \mathbf{F} = \nabla(\nabla \cdot \mathbf{F}) - \nabla \times (\nabla \times \mathbf{F})$$

$$\nabla \cdot (\nabla \times \mathbf{v}) = 0$$

$$\Leftrightarrow (-\nabla^2 - \varepsilon_0 \tau) e^{ik \cdot x} \mathbf{z} = e^{ik \cdot x} \mathbf{b} + \frac{1}{\varepsilon_0 \tau} \nabla(\nabla \cdot e^{ik \cdot x} \mathbf{b})$$

$$\Leftrightarrow \boxed{(\tilde{A} - \tau \varepsilon_0 I)} \tilde{\mathbf{z}} = \tilde{\mathbf{b}}$$

FFT based preconditioned

Simple Cubic: Preconditioner

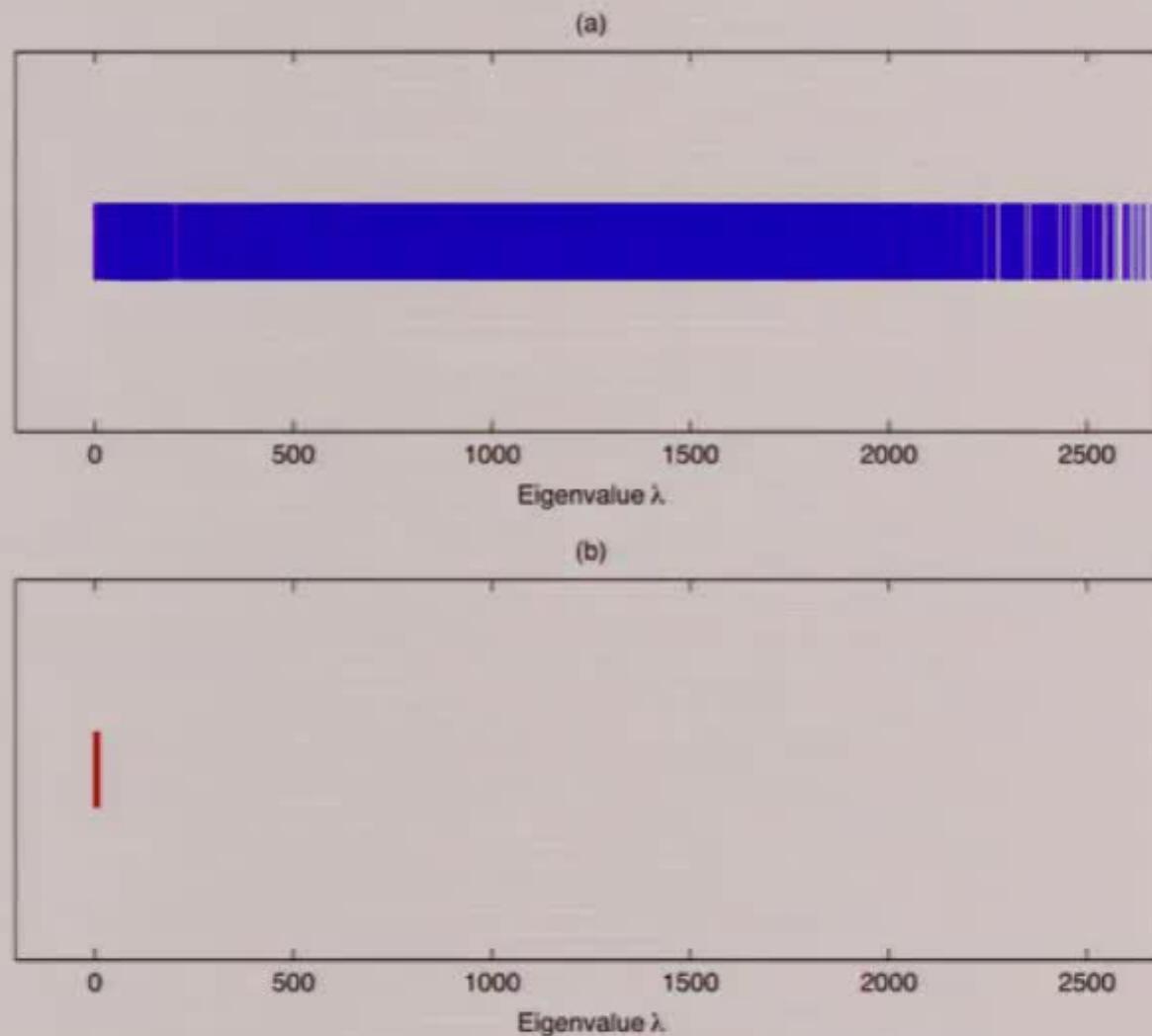
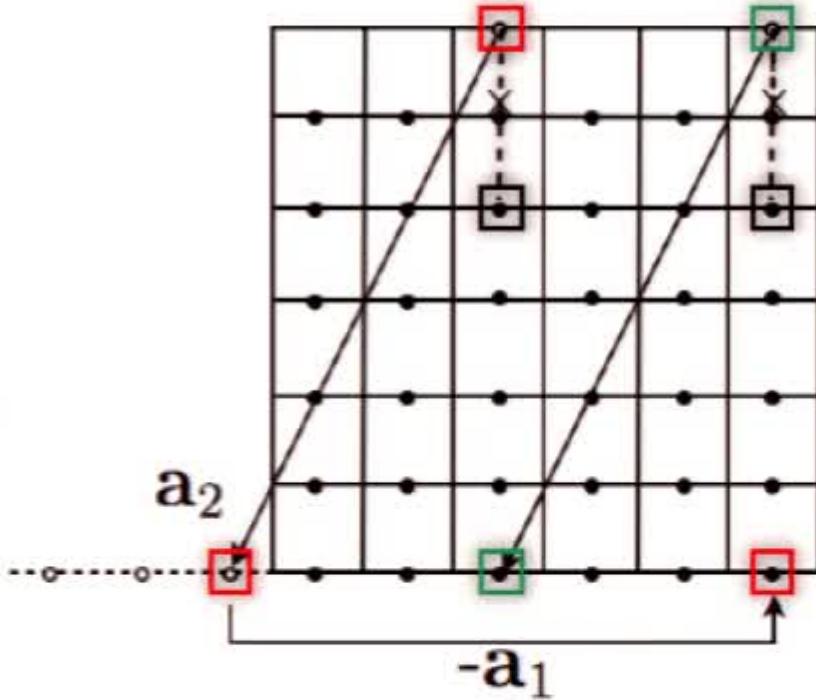


Fig. 4. Spectrum of (a) the matrix $(A - \tau B)$ and (b) the matrix $(A - \tau c_0)^{-1}(A - \tau B)$. The matrices sizes are equal to 10,125 (3×15^3), $\mathbf{k} = (\pi, \pi, 0)$, and $\tau = 0.01$.

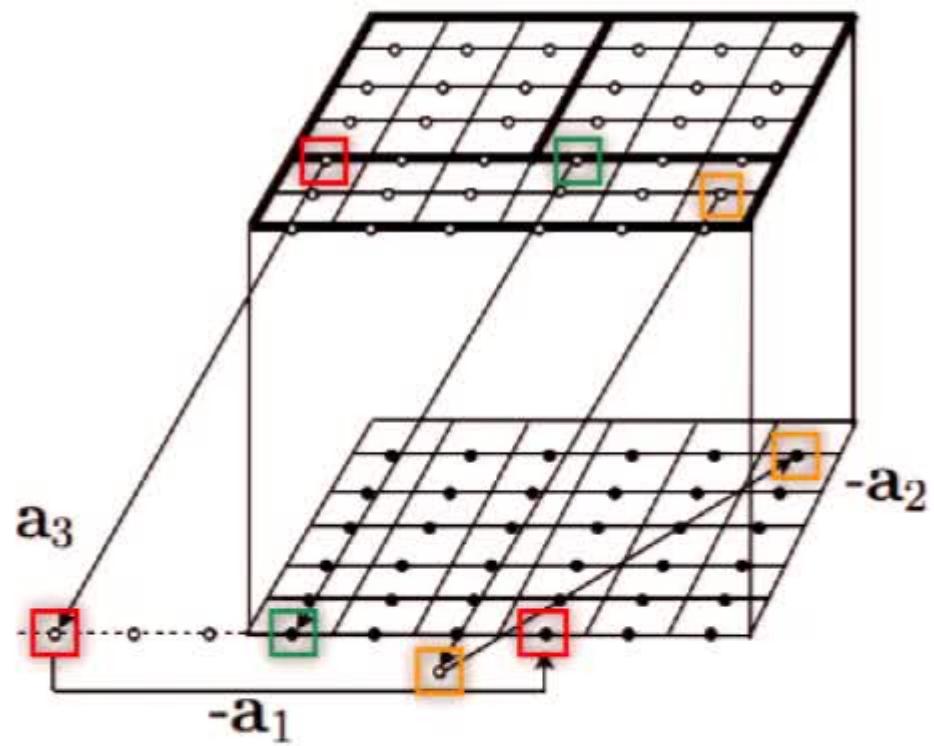
Face-Centered Cubic PC

Eigendecomposition of Double Curl

Preconditioner for 3D Coupling A?



$$\frac{\partial E_1}{\partial y}$$



$$\frac{\partial E_1}{\partial z}$$

Main Result: Eigendecomposition of A



$$A = Q_r \Lambda_r Q_r^*.$$

$$Q_r = \begin{bmatrix} Q_1 (3\Lambda_q^2 - \Lambda_q \Lambda_p)^{-\frac{1}{2}} & Q_2 (3\Lambda_q - \Lambda_p)^{-\frac{1}{2}} \end{bmatrix} \equiv (I_3 \otimes T) \Lambda.$$

Null space

Range space

Define

$$Q = \begin{bmatrix} Q_0 \\ Q_1 & Q_2 \end{bmatrix} \text{diag} \left(\Lambda_q^{-\frac{1}{2}}, (3\Lambda_q^2 - \Lambda_q \Lambda_p)^{-\frac{1}{2}}, (3\Lambda_q - \Lambda_p)^{-\frac{1}{2}} \right).$$

Then Q is unitary. Furthermore,

$$Q^* A Q = \text{diag} (0, \Lambda_q, \Lambda_q)$$

The Standard Eigenvalue Problem



$$\text{span} \left\{ B^{-1} Q_r \Lambda^{1/2} \right\} = \{ \mathbf{x} | A\mathbf{x} = \lambda B\mathbf{x}, \lambda > 0 \}$$



$$A\mathbf{x} = \lambda B\mathbf{x}$$

GEVP (3n-by-3n)

$$\left(\Lambda_r^{-\frac{1}{2}} Q_r^* \right) A \left(B^{-1} Q_r \Lambda_r^{\frac{1}{2}} \mathbf{y} \right) = \lambda \left(\Lambda_r^{-\frac{1}{2}} Q_r^* \right) B \left(B^{-1} Q_r \Lambda_r^{\frac{1}{2}} \mathbf{y} \right)$$

$$\left(\Lambda_r^{\frac{1}{2}} Q_r^* B^{-1} Q_r \Lambda_r^{\frac{1}{2}} \right) \mathbf{y} = \lambda \mathbf{y}$$

$$A_r \mathbf{y} = \lambda \mathbf{y}$$

SEVP (2n-by-2n)

Since $A = Q_r \Lambda_r Q_r^*$,

$$\Rightarrow \Lambda_r^{-\frac{1}{2}} (Q_r^* A) = \Lambda_r^{-\frac{1}{2}} (\Lambda_r Q_r^*) = \Lambda_r^{\frac{1}{2}} Q_r^*.$$

To Compute T^*p for Q_r^*



$$A_r \mathbf{y} = \lambda \mathbf{y}$$

$$A_r = \Lambda_r^{\frac{1}{2}} Q_r^* B^{-1} Q_r \Lambda_r^{\frac{1}{2}}$$

$$Q_r = \begin{bmatrix} Q_1 (3\Lambda_q^2 - \Lambda_q \Lambda_p)^{-\frac{1}{2}} & Q_2 (3\Lambda_q - \Lambda_p)^{-\frac{1}{2}} \end{bmatrix} \equiv (I_3 \otimes T) \Lambda.$$

- Rewrite T as **periodic** part (by FFT) and **diagonal** part

$$T_{i,j}^* \mathbf{p} = U_{\mathbf{z}}^* E_{\mathbf{z},i+j}^* P^\top (\overline{\mathbf{y}_{i,j}} \otimes \mathbf{x}_i)$$

via FFT

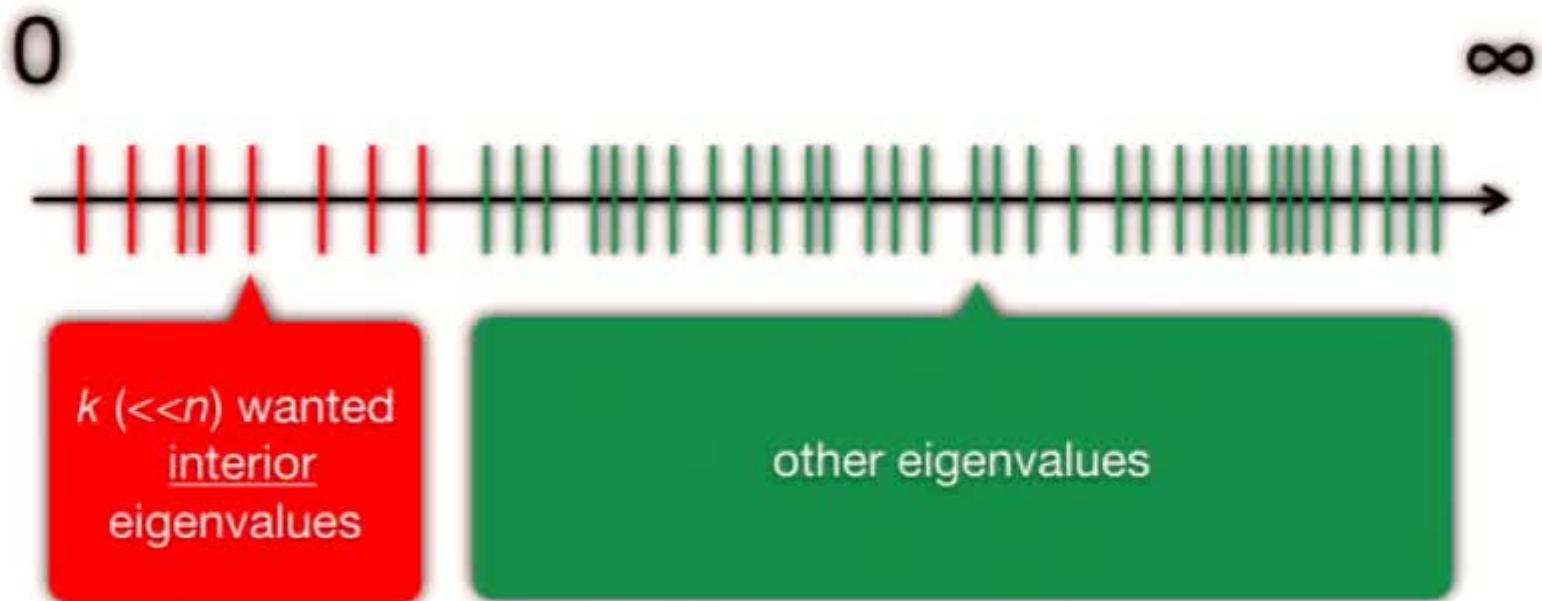
diagonal

matrix-vector multp.

Null-space Free Lanczos Method (NFLM)



- Dim. of GEVP is $3d$ and SVEP is $2n$
- GEVP and SEVP has same $2n$ positive eigenvalues
- SEVP has NO zero eigenvalues



Null-space Free Lanczos Method (NFLM)



- Hermitian positive definite linear system in NFLM

$$A_r \mathbf{y} = \lambda \mathbf{y}; \quad A_r = \Lambda_r^{\frac{1}{2}} (Q_r^* B_\varepsilon^{-1} Q_r) \Lambda_r^{\frac{1}{2}}.$$
$$(Q_r^* B_\varepsilon^{-1} Q_r) \mathbf{u} = \mathbf{c}$$

- Well-conditioned

$$\kappa(Q_r^* B_\varepsilon^{-1} Q_r) \leq \boxed{\kappa(B_\varepsilon^{-1})}$$

$B_\varepsilon = \text{diag}(\dots, 1, \dots, 13, \dots)$
electric permittivity of Si=13, Air=1

- CG convergence ratio

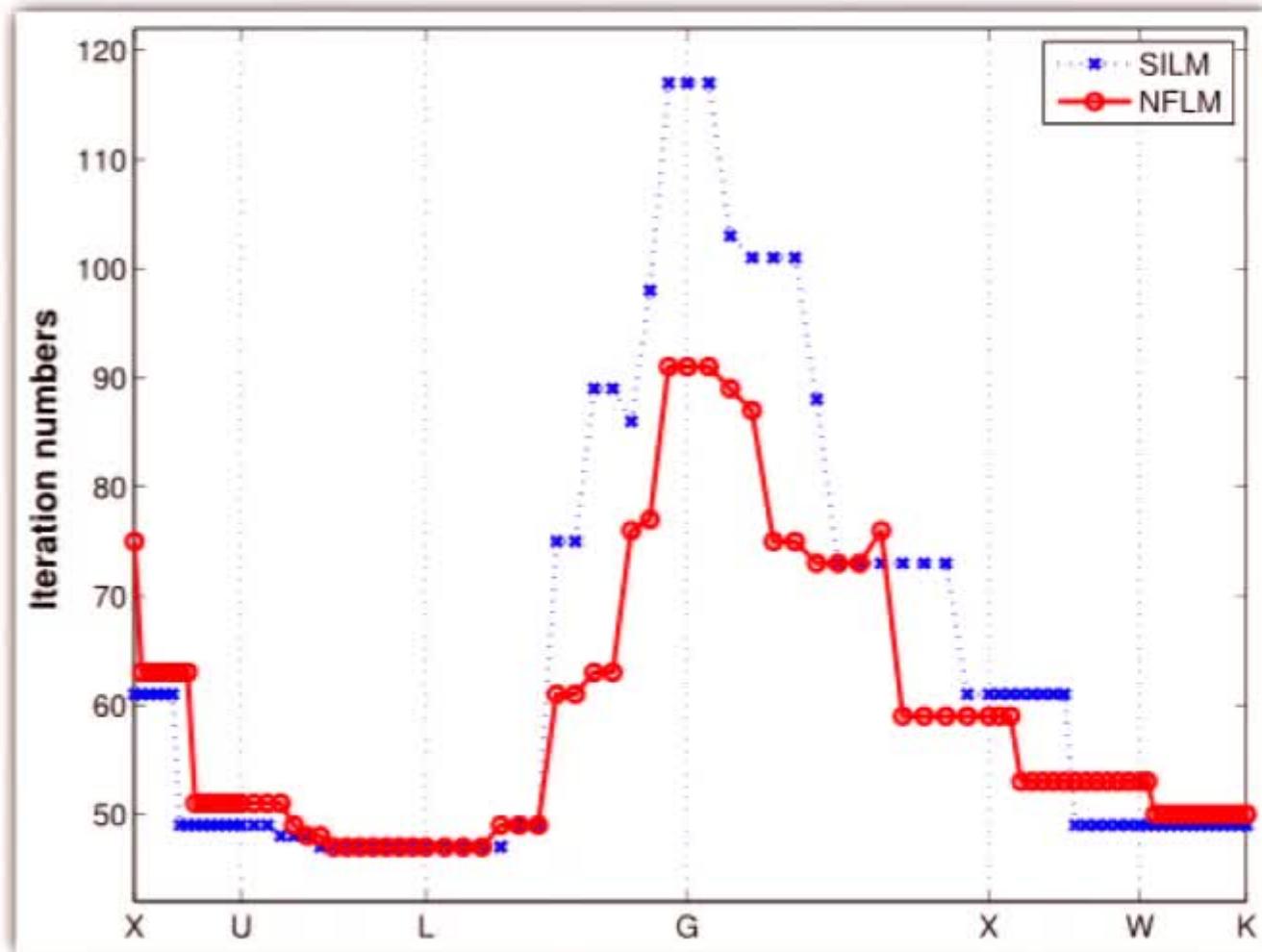
$$\gamma = \frac{\sqrt{\kappa(Q_r^* B_\varepsilon^{-1} Q_r)} - 1}{\sqrt{\kappa(Q_r^* B_\varepsilon^{-1} Q_r)} + 1} \leq \gamma_B = \frac{\sqrt{13} - 1}{\sqrt{13} + 1} \approx 0.5657$$

$$(\gamma_B)^{40} \approx 1.27 \times 10^{-10}$$

Small Iteration Numbers for SILM & NFLM



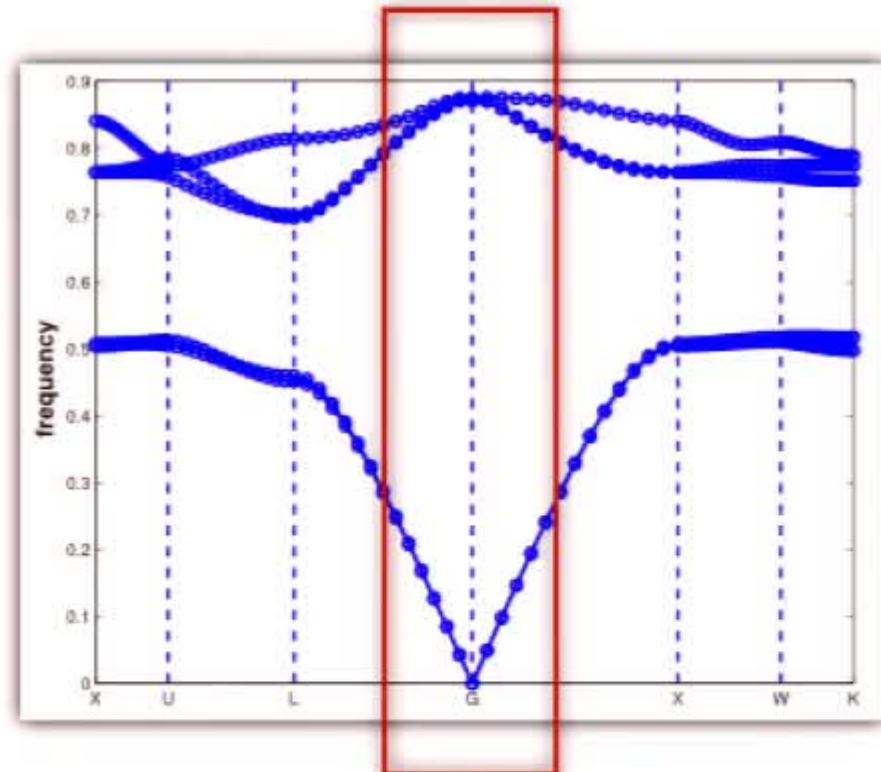
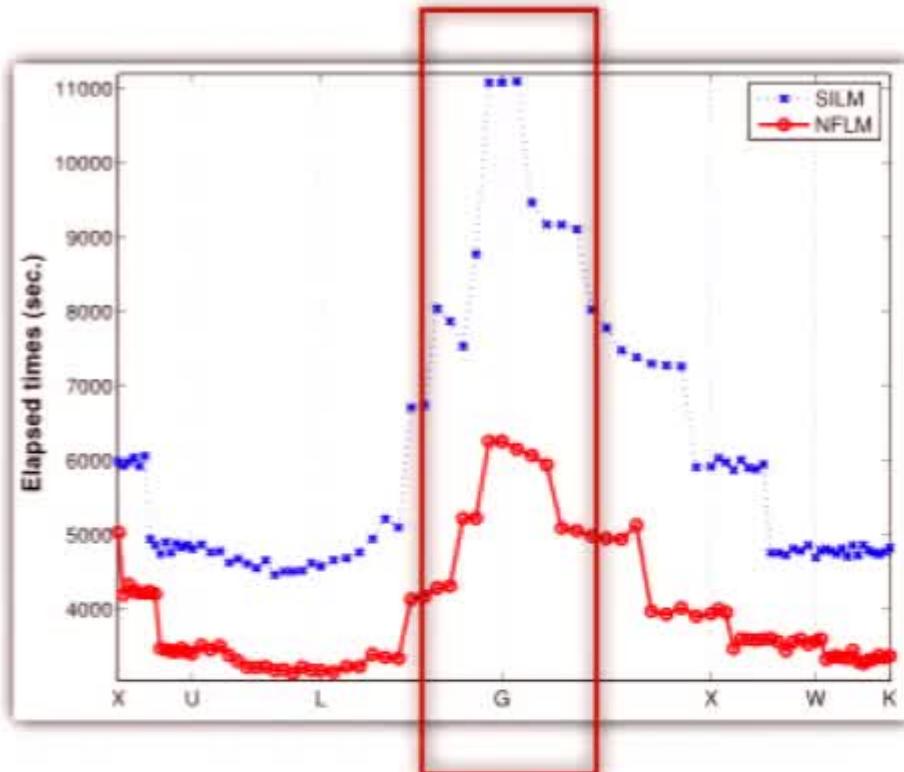
- SILM: 47-115 ite.; D= 5,184,000;
- NFLM: 47-91 iter; D= 3,456,000;



Time for SILM and NFLM



- SILM: 47-115 ite.; D=5,184,000; 2 MVmult in CGS
- NFLM: 47-91 iter; D= 3,456,000; 1 MVmult in CG



Maxwell's Equation for 3D Photonic Crystal

Eigendecomposition of A

GEVP $Ax = \lambda Bx$
3n x 3n (n zero e.v.)

SEVP $Ax = \lambda x$
2n x 2n (no zero e.v.)

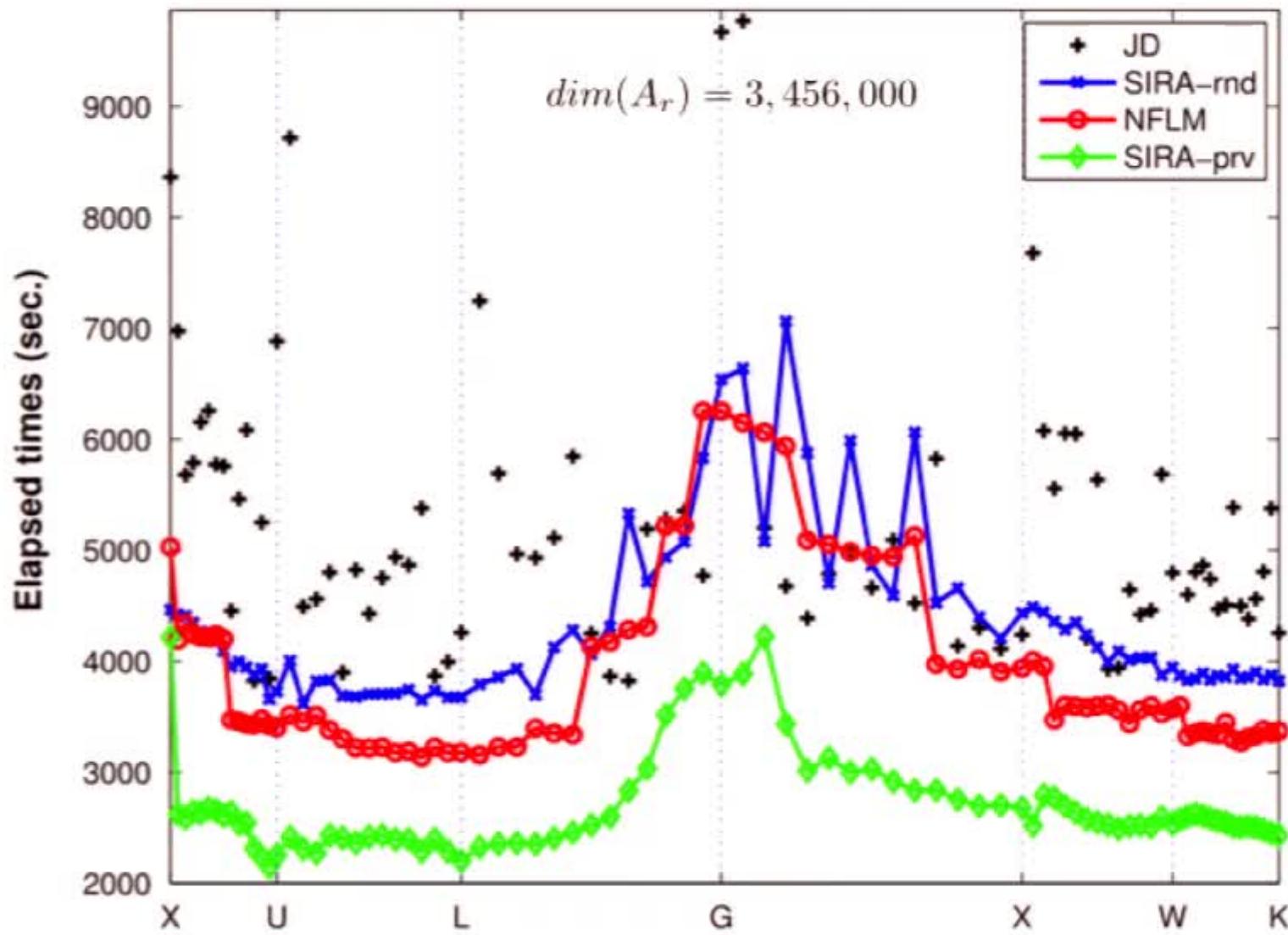
Shift-and-Invert
Lanczos Method

Inv. Lanczos, Jacobi-
Davidson, SI Res Arnoldi

Iterative Linear System Solver and Preconditioning

FFT-based Matrix Vector Multiplication

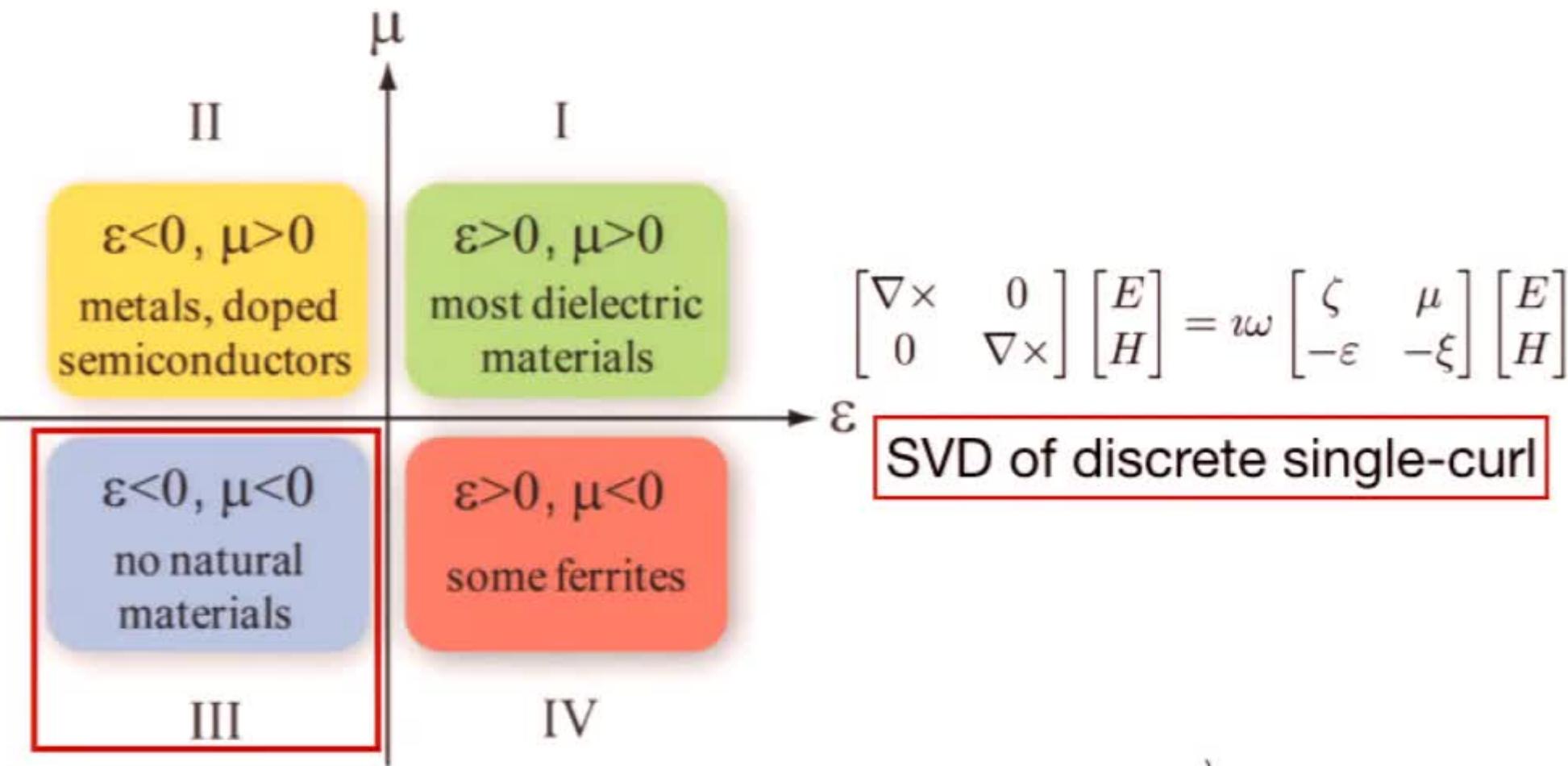
Comparison of Eigensolvers



Complex Media

SVD of Single Curl

Artificial Complex Media



Complex Media



Singular Value Decomposition of Single Curl

$$C = P \operatorname{diag} \left(\Lambda_q^{1/2}, \Lambda_q^{1/2}, 0 \right) Q^*$$

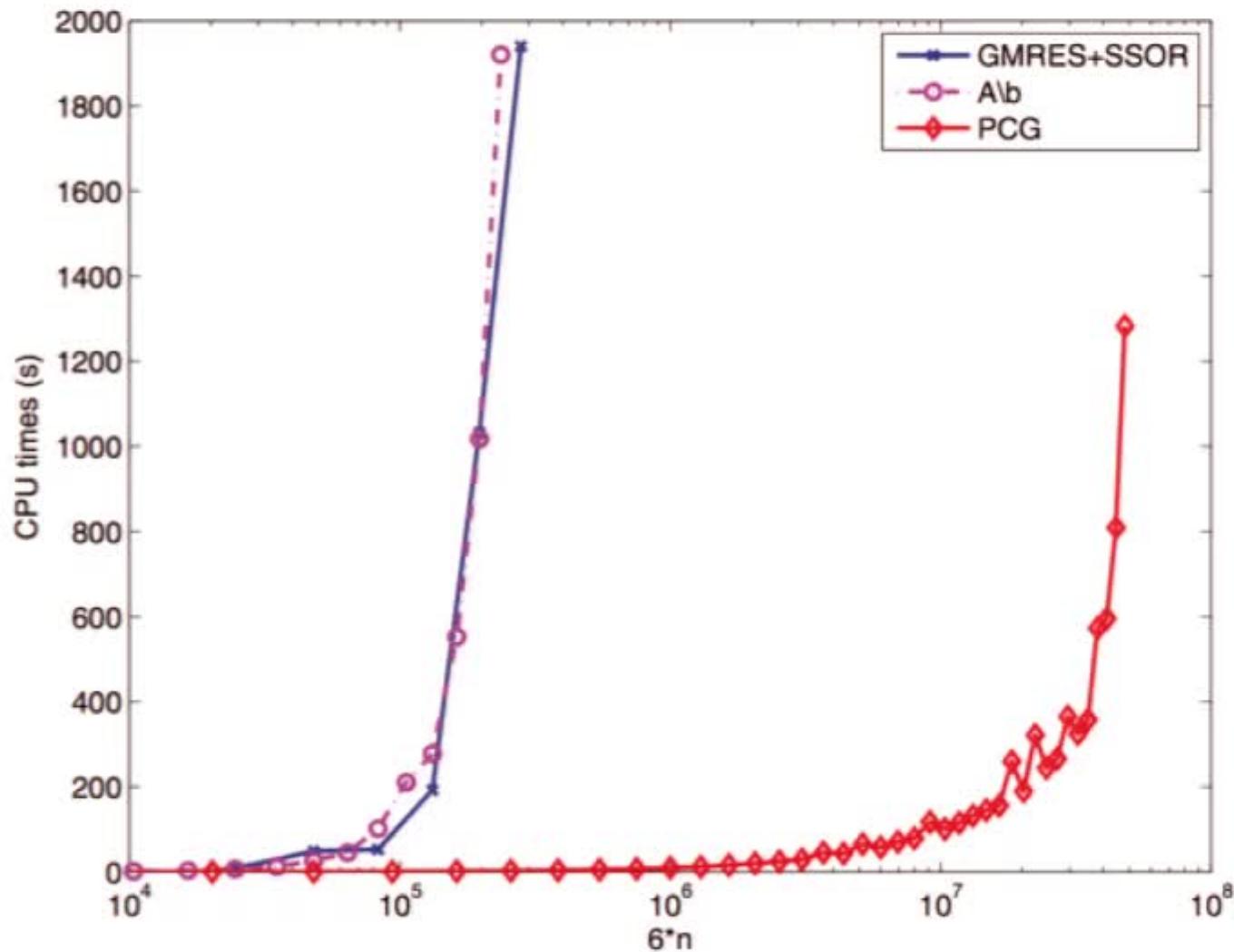
Eigenvalue problem	Generalized non-Hermitian	Generalized Hermitian & HPD
Eigenvalue	Complex	Real
Dimension	6n x 6n	4n x 4n
# of zero eigenvalue	2n	0 (null space free)
Eigensolver	S.I. Arnoldi	S.I. Lanczos
Shift	Hard to choose	0
Linear system solver	LU or GMRES (not efficient)	CG w/ FFT mtx-vec mult
Preconditioner	Hard to find	SVD+FFT (well-cond.)
Applications	Complex media	Chiral/Pseudochiral
Embedded linear system	$\left(\begin{bmatrix} C & 0 \\ 0 & C^* \end{bmatrix} - \sigma \begin{bmatrix} \zeta_d & I_{3n} \\ -\varepsilon_d & -\xi_d \end{bmatrix} \right) y = b$	See below

$$A_r u \equiv \begin{bmatrix} P_r^* & Q_r^* \\ I_{3n} & 0 \end{bmatrix} \begin{bmatrix} \zeta_d & -I_{3n} \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \Phi^{-1} & 0 \\ 0 & I_{3n} \end{bmatrix} \begin{bmatrix} \zeta_d^* & I_{3n} \\ -I_{3n} & 0 \end{bmatrix} \begin{bmatrix} P_r & Q_r \\ Q_r & 0 \end{bmatrix} u = b$$

Time for Embedded Linear Systems



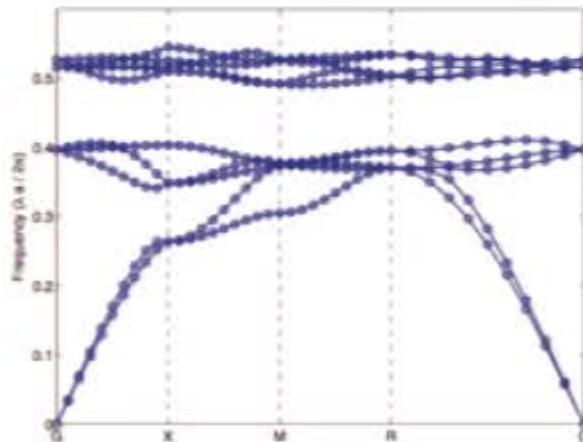
Quad-Core Xeon X5687 3.6GHz CPUs, 48GB, MATLAB 2014a



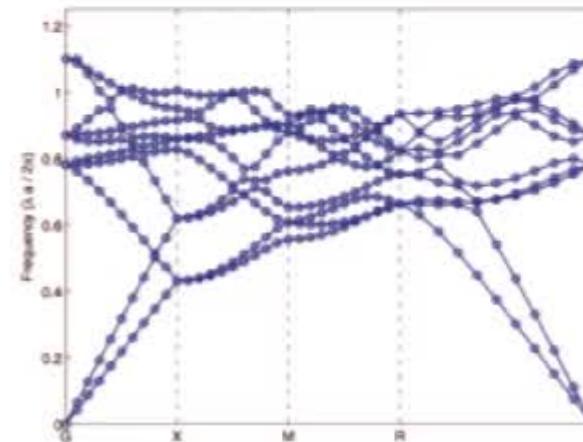
Iteration Number



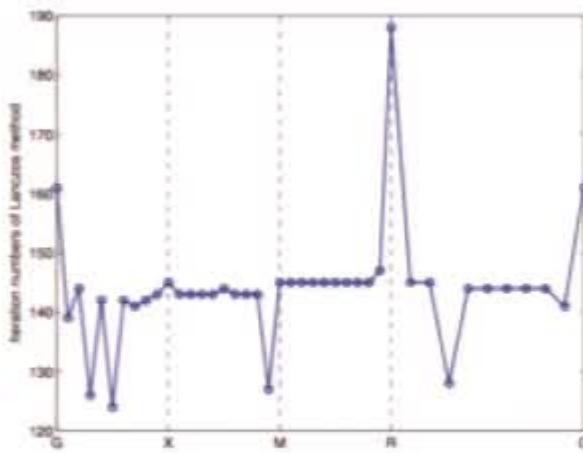
Dim: $4 \times 128 \times 128 \times 128 = 8,388,608$
Quad-Core Xeon X5687 3.6GHz CPUs, 48GB, MATLAB 2014a



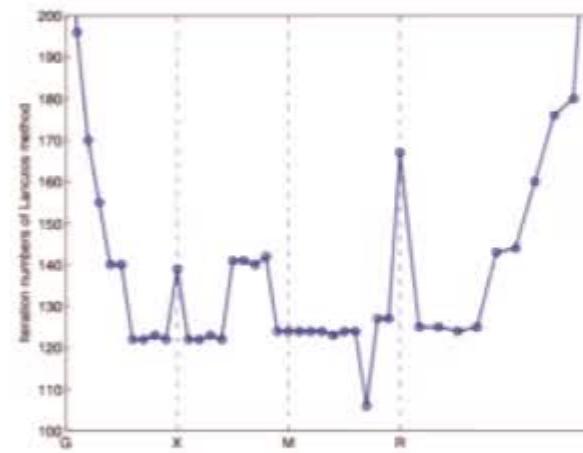
(a) Band structure for $(\varepsilon_i, \varepsilon_o, \gamma) = (13, 1, 0.5)$



(b) Band structure for $(\varepsilon_i, \varepsilon_o, \gamma) = (1, 1, 0.8)$



(c) Iteration numbers ranging from 120 to 190
with average 143 for $(\varepsilon_i, \varepsilon_o, \gamma) = (13, 1, 0.5)$

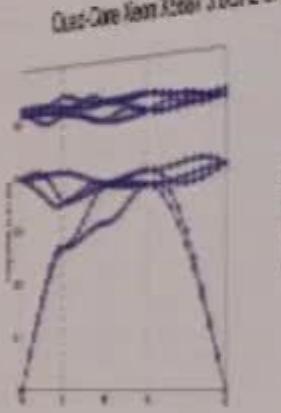


multiplicity of eigenvalues

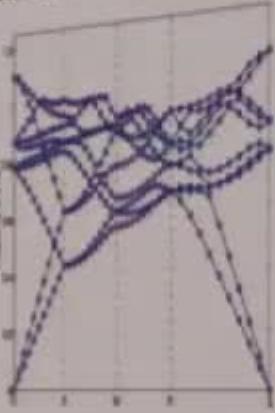
(d) Iteration numbers ranging from 100 to 200
with average 136 for $(\varepsilon_i, \varepsilon_o, \gamma) = (1, 1, 0.8)$

Iteration Number

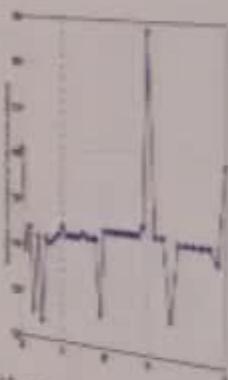
Dim $4 \times 28 \times 125 \times 125 = 8,388,800$
Quad-Core Xeon X5687 3.6GHz CPUs, 4GB, MATLAB 2014a



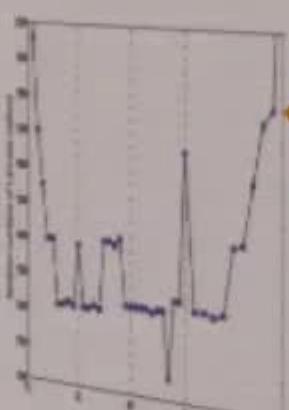
(a) Bad strategy for $\{r_1, t_0, \gamma\} = \{1, 1, 0.5\}$



(b) Good strategy for $\{r_1, t_0, \gamma\} = \{1, 1, 0.8\}$



(c) Bad strategy ranging from 22 to 36 with varying r_1 for $\{t_0, \gamma\} = \{1, 1, 0.5\}$



(d) Bad strategy ranging from 100 to 200 with varying r_1 for $\{t_0, \gamma\} = \{1, 1, 0.8\}$

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Software

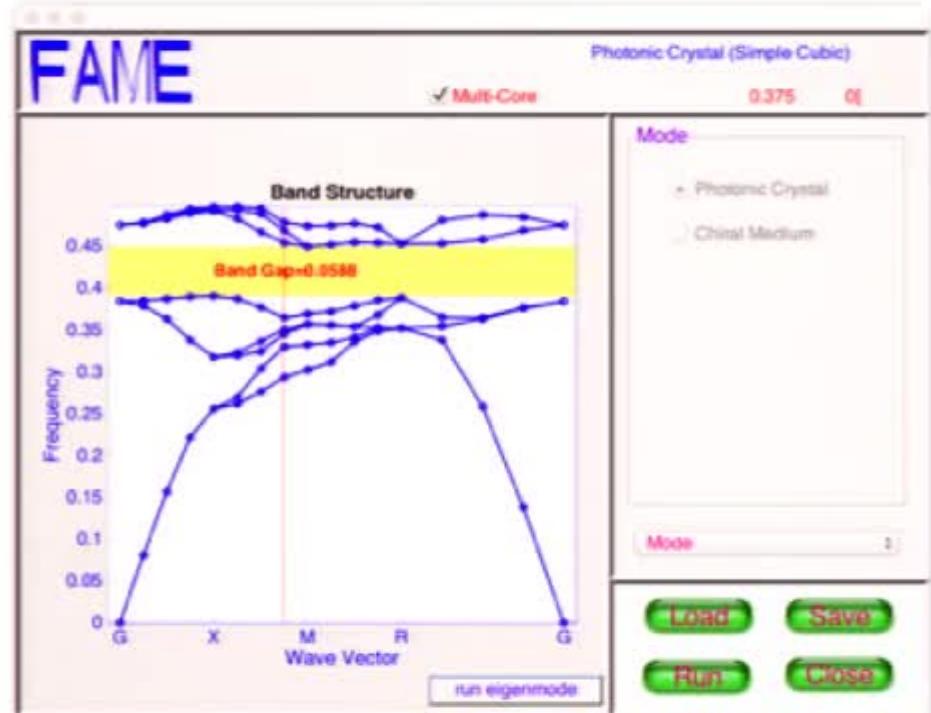
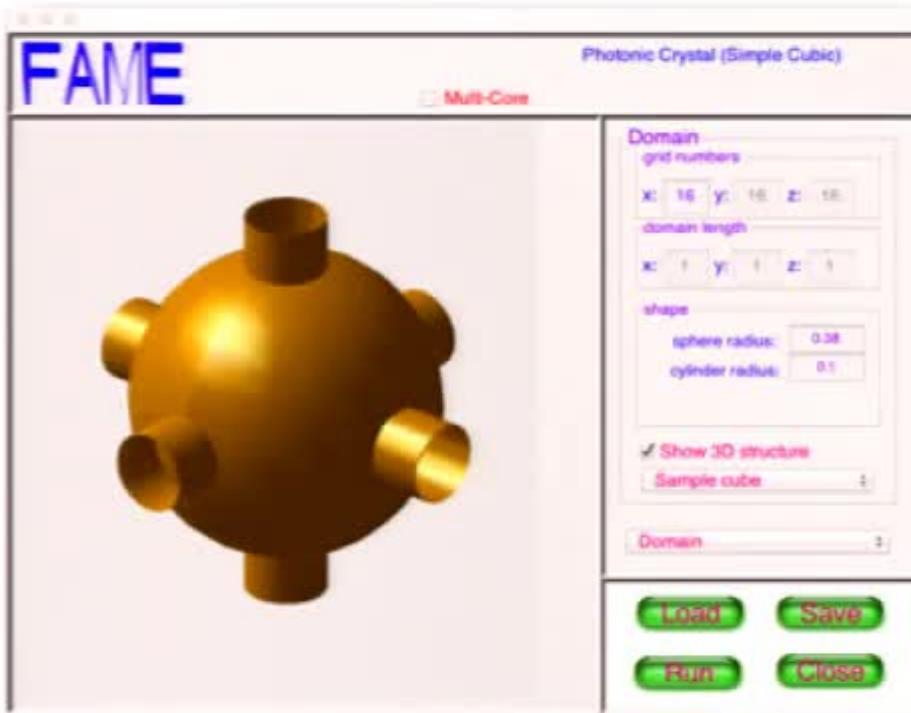
38

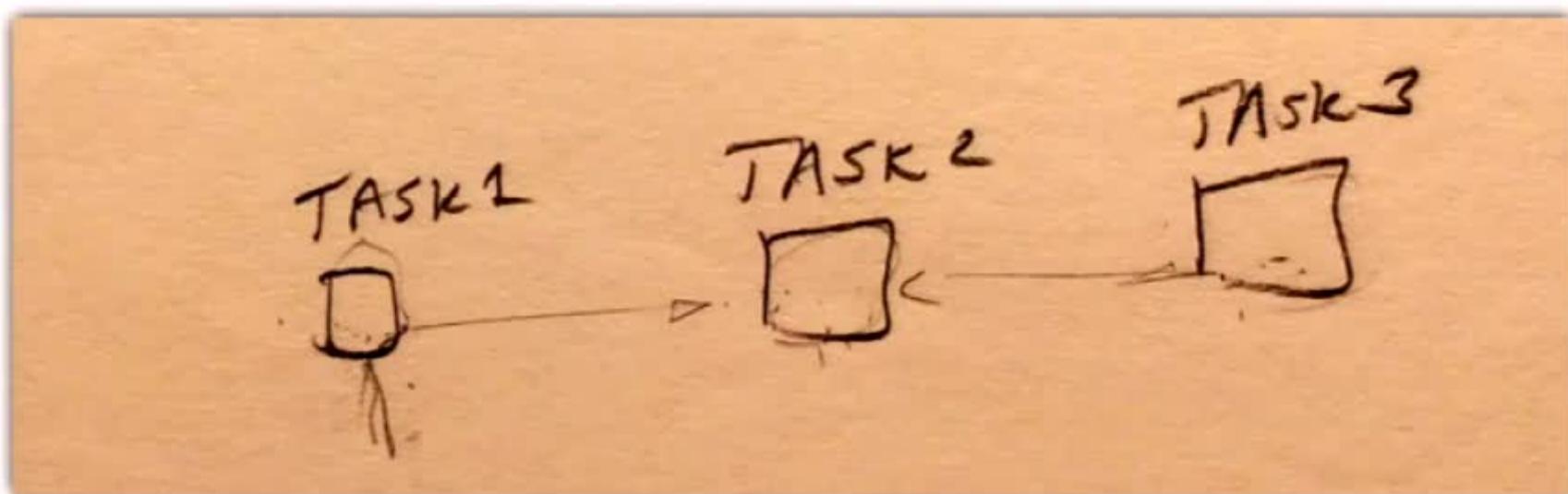
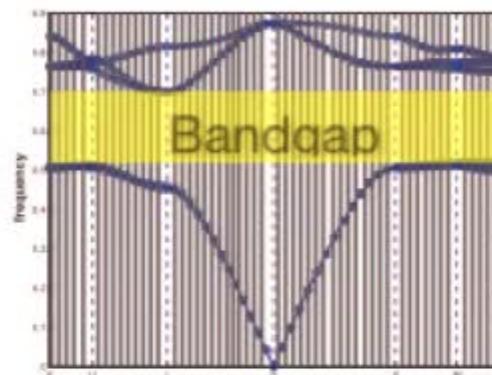
FAME.MATLAB



- Various applications
(e.g. photonic crystal, metallic materials, complex material)
- GUI for photonic structure design
- Choices of eigensolvers and linear system solvers
- Built-in functions for efficient algorithms
- Visualization of results
- Extendable to new algorithms and applications

FAME.MATLAB





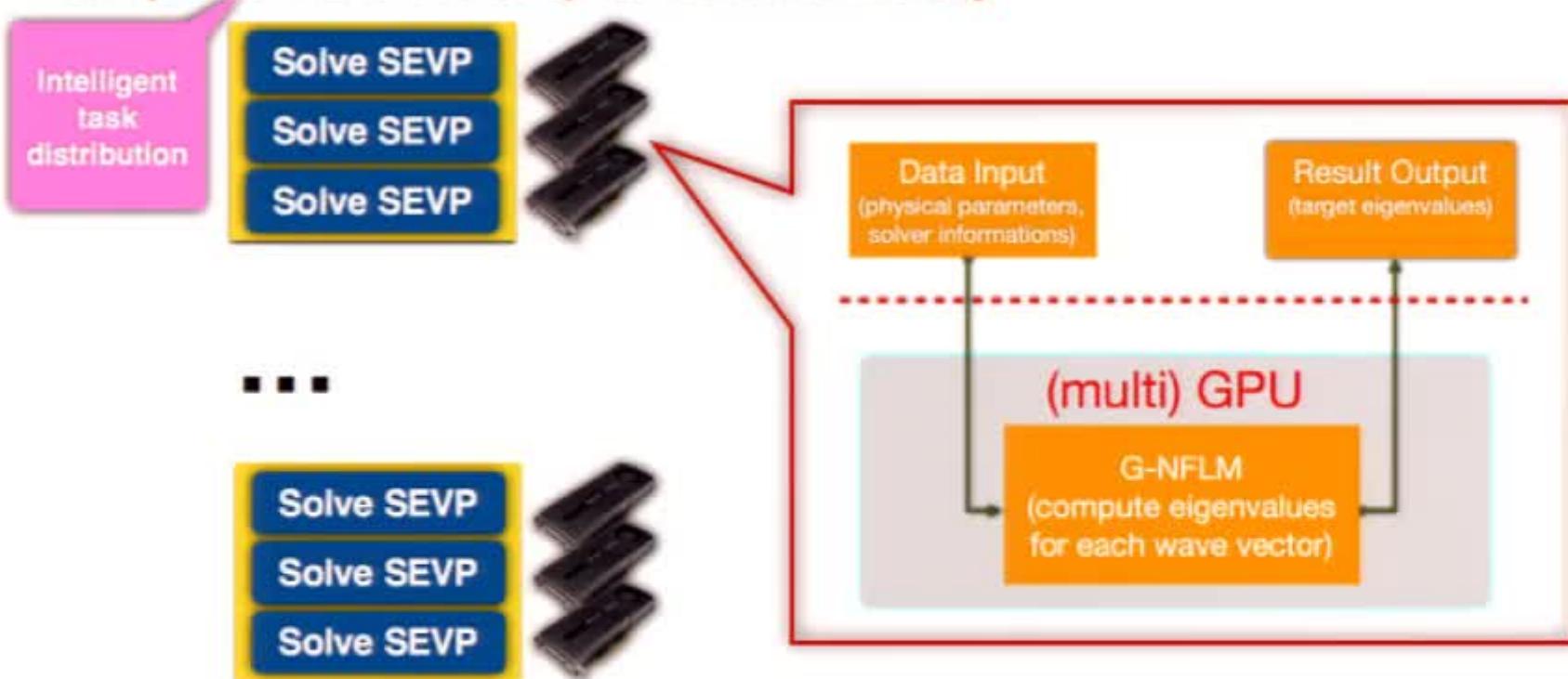
A hand-drawn diagram on a light orange background showing the letters "MPI" in a bold, black, sans-serif font. The letters are slightly irregular and have a hand-drawn appearance.

FAME.GPU



- Cluster with 16 MPI nodes (each has 3 NVIDIA Tesla M2070)
- Minimal communications between CPU and GPU

for (some wave vectors) do simultaneously

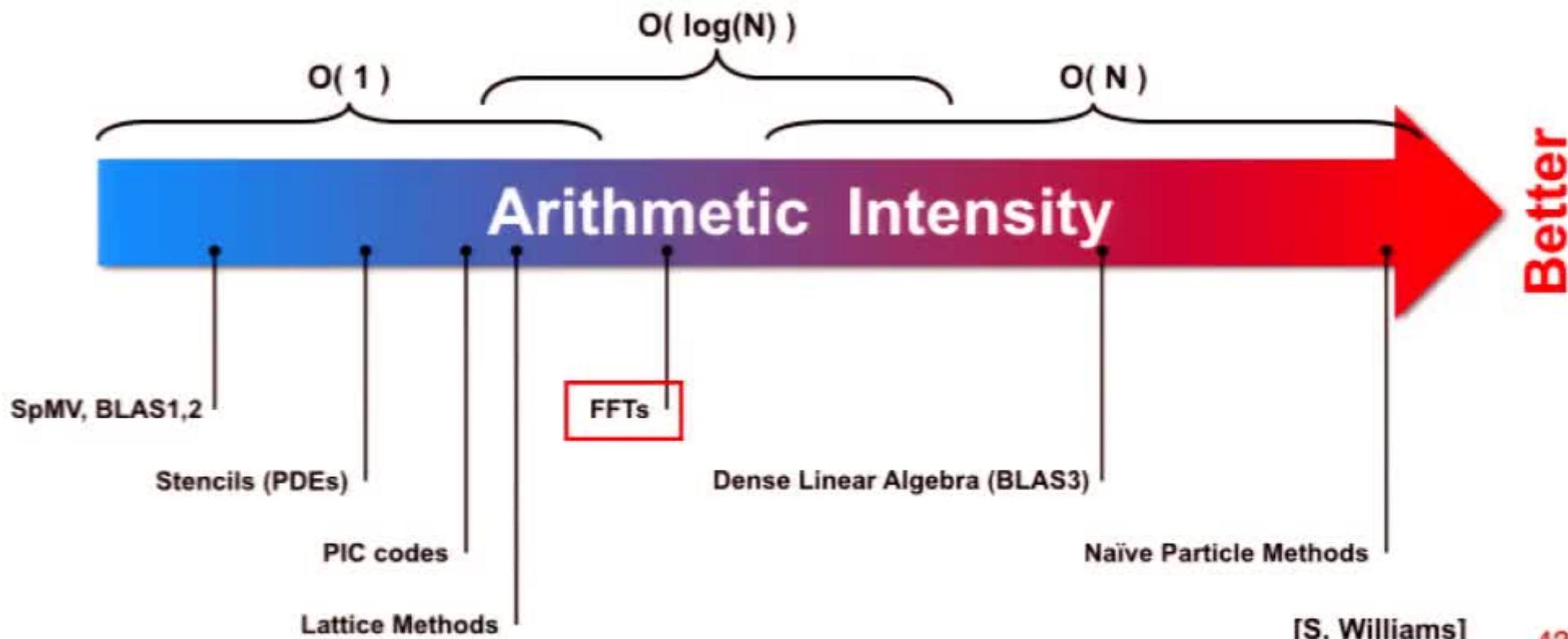


end for

FAME.GPU



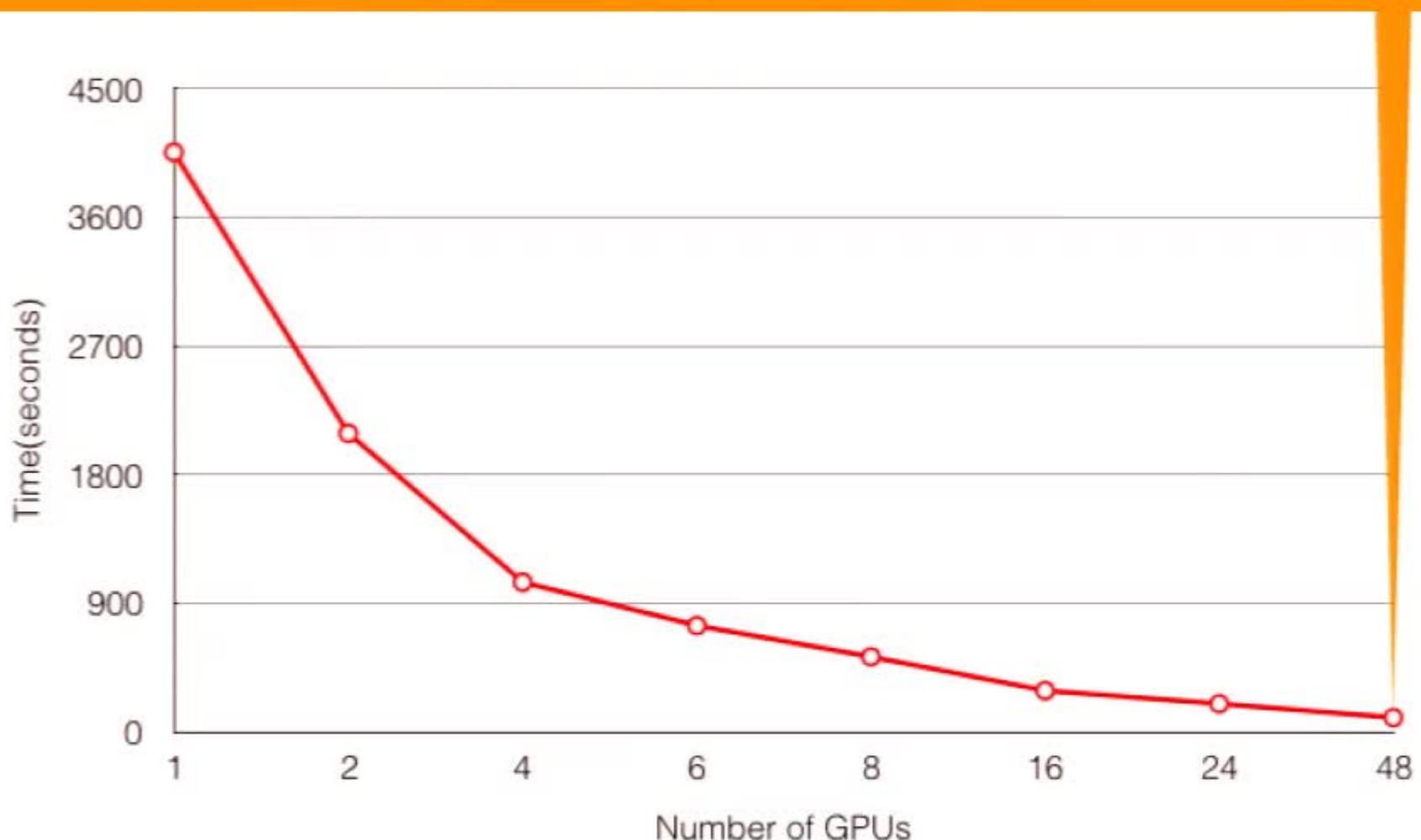
- Algorithm:
Inverse Lanczos (w/ restart) + Nullspace free + CG (well-cond.)
- Parallelism:
Embarrassing parallel in terms of wave vectors w/ intelligent task distribution



Scalability on Multiple GPU (6.29M)



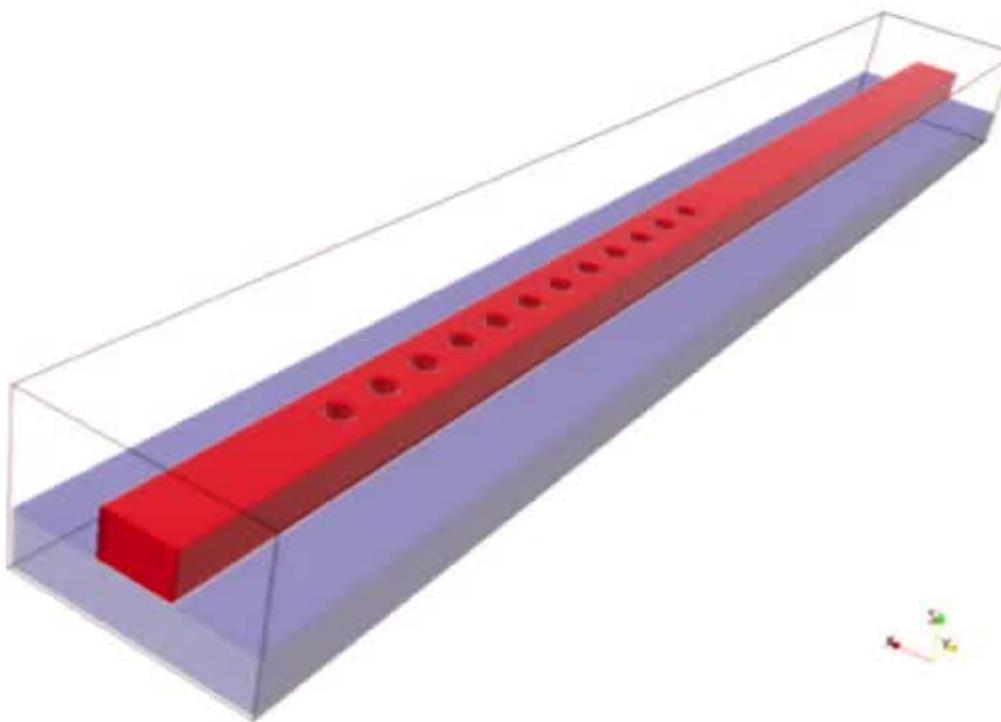
~1.5 minute for a complete band structure diagram



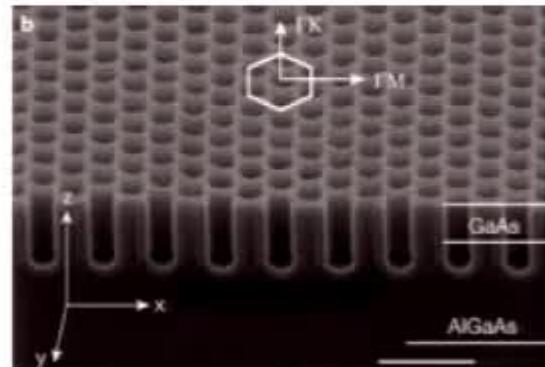
Compressed Hierarchical Schur

Memory Saving BLAS3 Direct Solver

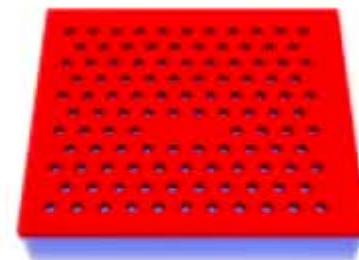
Photonic Devices



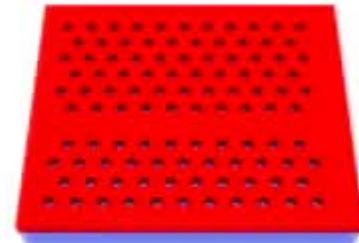
Periodic airhole array on a CdTe ridge waveguide



Ref: Chow et al., "Three-dimensional control of light in a two-dimensional **photonic crystal slab**," Nature, vol. 407, 983-986 (26 October 2000)



Point defect photonic crystal slab

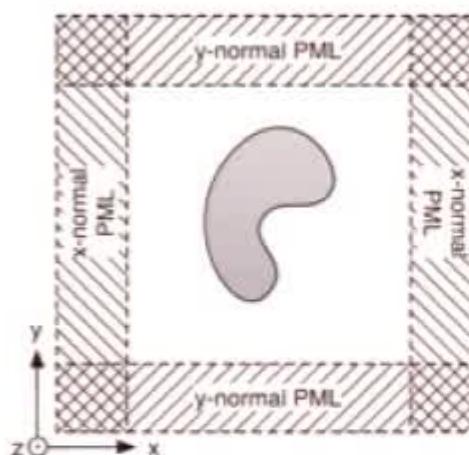
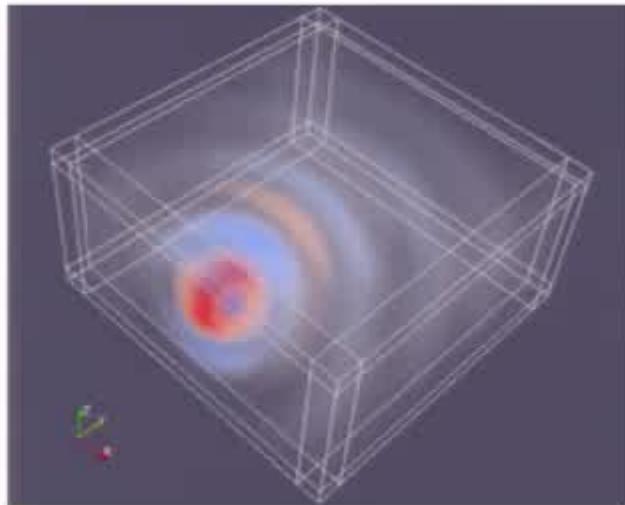


Linear defect photonic crystal slab

Perfectly Matched Layer (PML)



- $\nabla_s = \hat{x} \frac{1}{s_x} \frac{\partial}{\partial x} + \hat{y} \frac{1}{s_y} \frac{\partial}{\partial y} + \hat{z} \frac{1}{s_z} \frac{\partial}{\partial z}$ on complex stretching coordinate
- w -normal scale stretching factors: $s_w(l) = 1 + \frac{j\sigma_w(l)}{\omega\epsilon_0}$
($w = x, y, z$ and l : depth from PML interface)
- Artificial absorptive layer that absorb outward wave and reduce reflection from boundary



Ref: Shin and Fan, JCP, vol. 231, pp. 3406, 2012.

Linear System Problems



- Nonmagnetic vector wave equation ($\mu_r = 1$)

- $$-\nabla_s \times \nabla_s \times \vec{E} + k_0^2 \varepsilon_r \vec{E} = j\omega \mu_0 \vec{J}$$

\vec{E} : electric field
 \vec{H} : magnetic field
 ω : angular frequency
 μ : permeability
 μ_0 : vacuum permeability
 μ_r : relative permeability
 ε : permittivity
 ε_0 : vacuum permittivity
 ε_r : relative permittivity
 k_0 : vacuum wavenumber
 λ_0 : vacuum wavelength

- Linear problem:
For $k_0 = \omega \sqrt{\mu_0 \varepsilon_0} = \frac{2\pi}{\lambda_0}$, given \vec{J} , λ_0 , ε_r , solve \vec{E}
- Photonic band analysis equation without source

- $$\frac{1}{\varepsilon_r} \nabla_s \times \nabla_s \times \vec{E} = k_0^2 \vec{E}$$

- Eigenvalue problem:

Given ε_r , \vec{k} , solve k_0^2 and band field pattern \vec{E}

- Shift-and-invert method with given shift k_s^2

- Solve
$$-\frac{1}{\varepsilon_r} (-\nabla_s \times \nabla_s \times \vec{\psi}_{i+1} + k_s^2 \varepsilon_r \vec{\psi}_{i+1}) = \vec{\psi}_i$$

- Non-Hermitian complex ill-conditioned linear systems

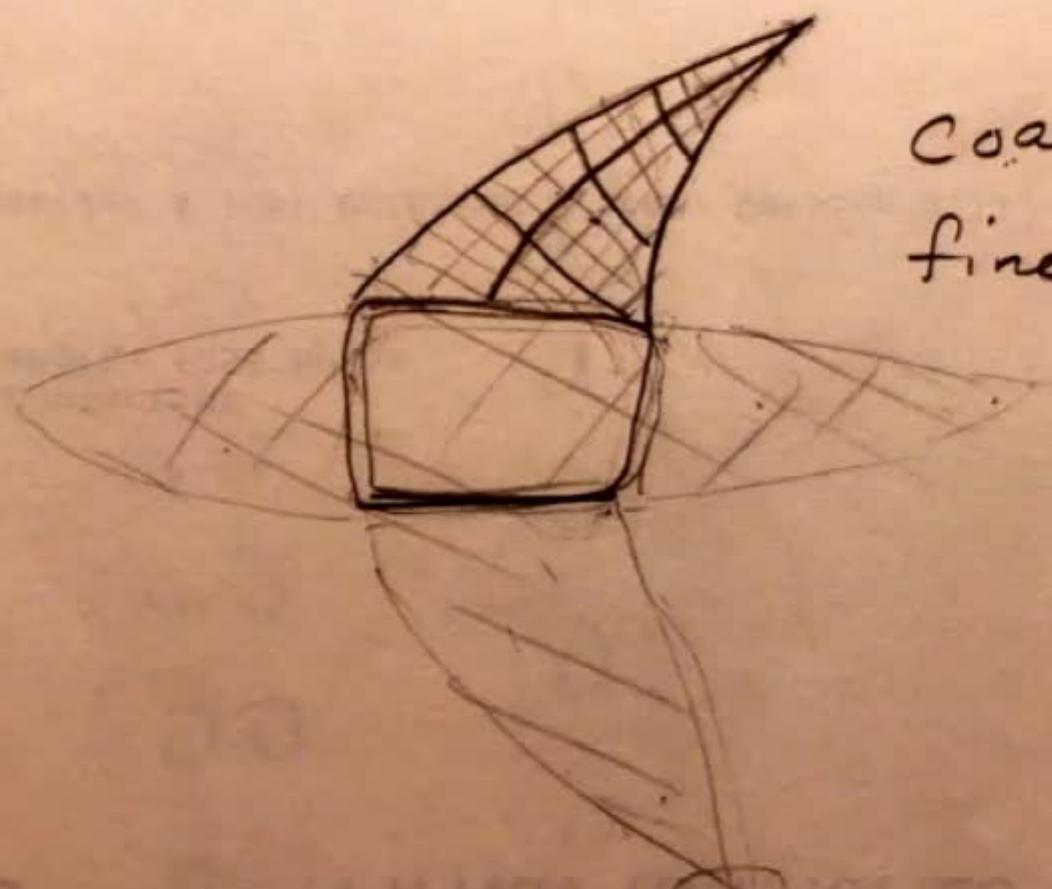
Challenges



- Non-Hermitian ill-conditioned linear systems
- Without eigendecomposition of the coefficient matrices
- Hard to find efficient preconditioner

DIRECT

Coarse
fine -

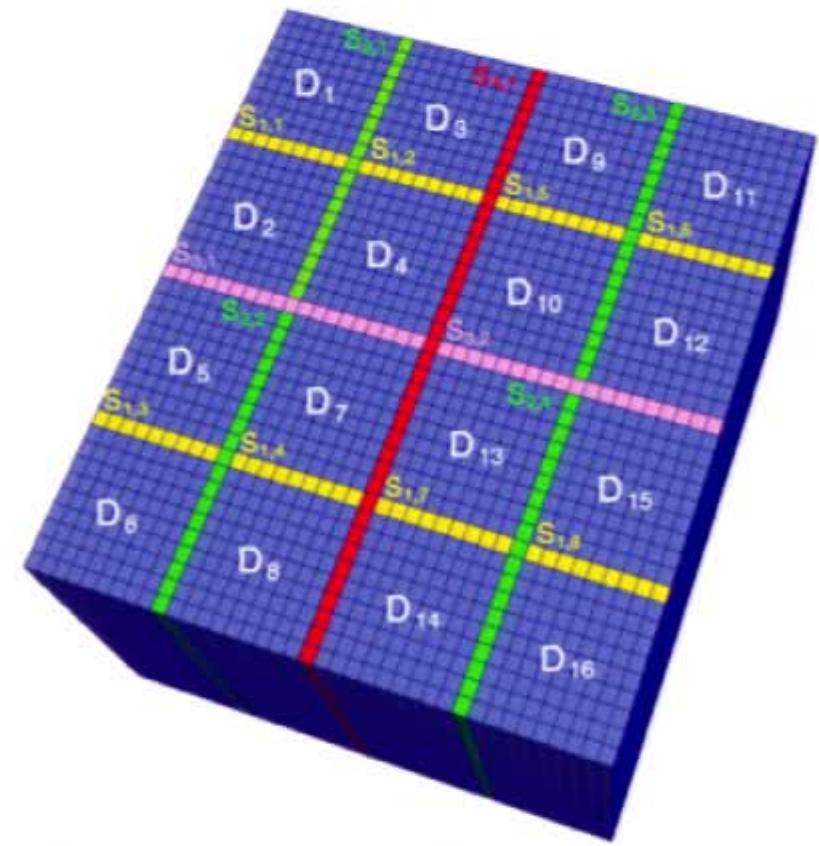
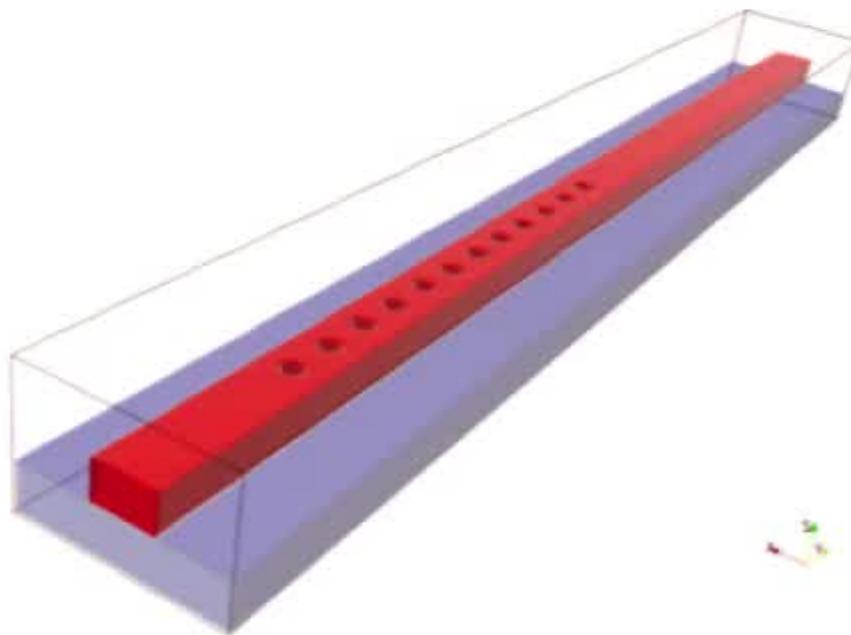




Approaches

- Hierarchical Schur direct solver
- Memory saving
- Multilevel parallelism
- BLAS3 operations

Physical and Computational Domains

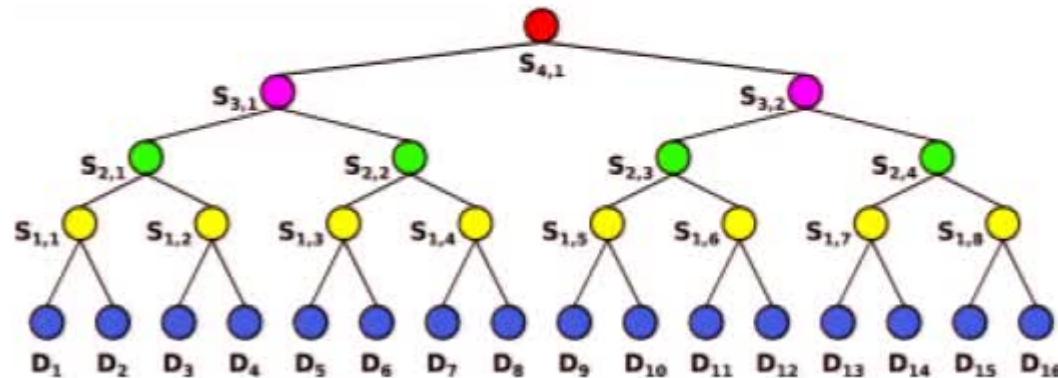
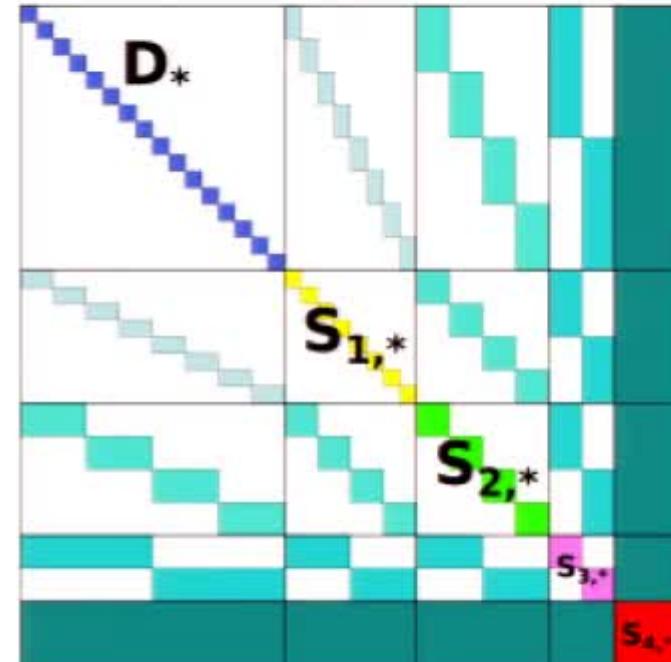
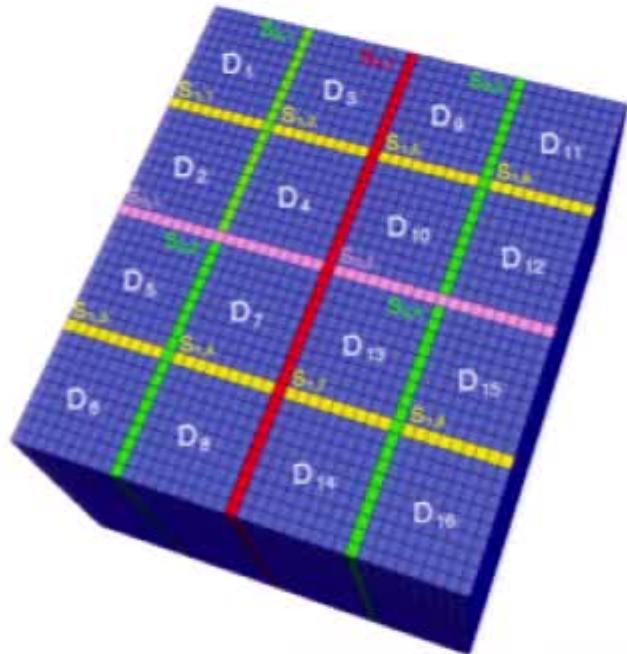


Subdomains, interfaces, separators
in domain decomposition

Elimination Tree and Matrix



- Customized nested dissection ordering



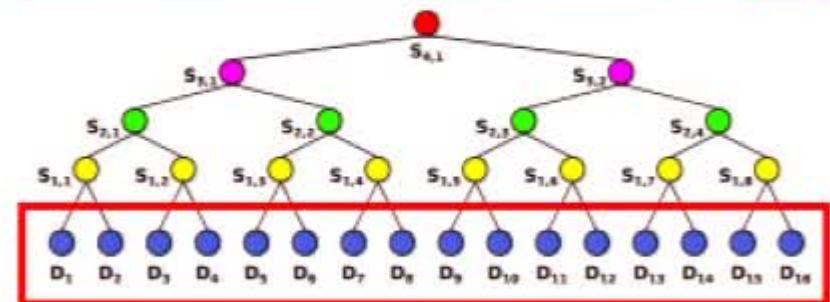
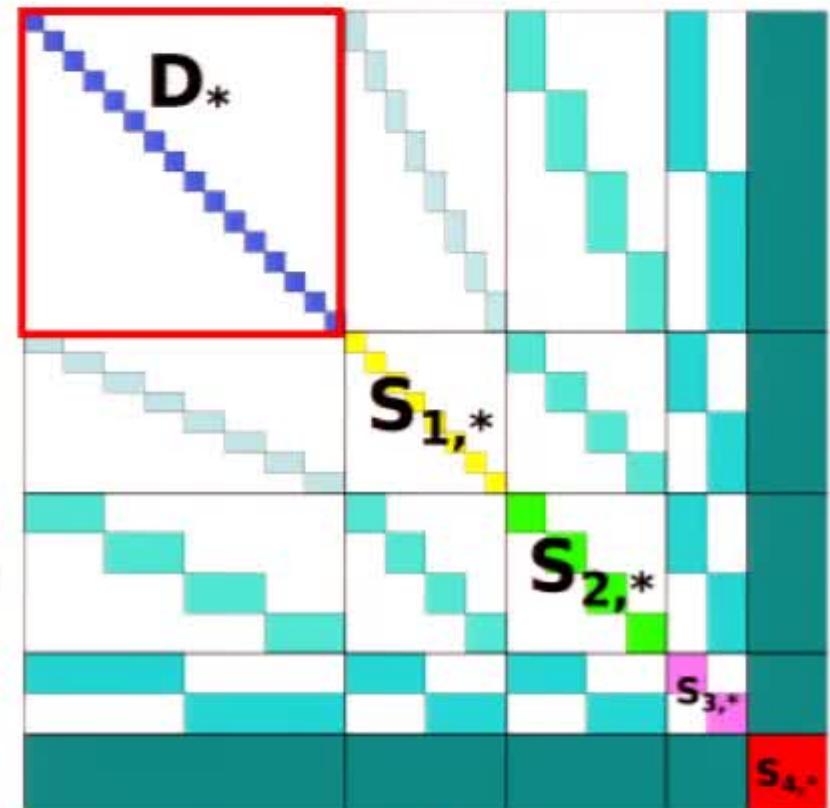
Subdomain Level



- Operations
 - Factor sparse diagonal blocks
 - Scale sparse interface blocks
 - Update Schur submatrices

$$S_{1,1} = A_{17,17} - A_{17,1}A_{1,1}^{-1}A_{1,17} - A_{18,2}A_{2,2}^{-1}A_{2,18}$$

- Parallelisms
 - One node for each D_i
 - Multi-threads for each D_i



Separator Levels

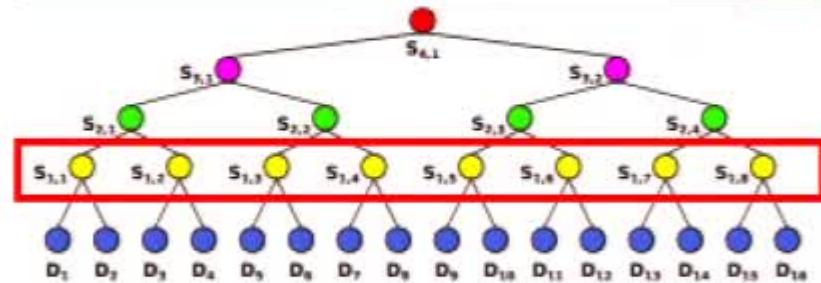
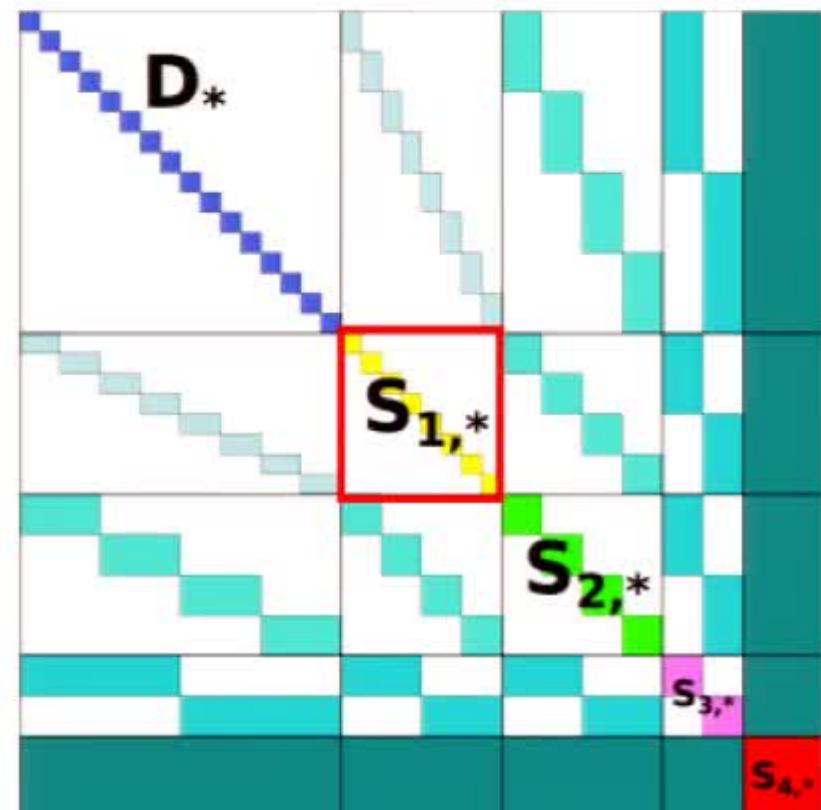


- Operations

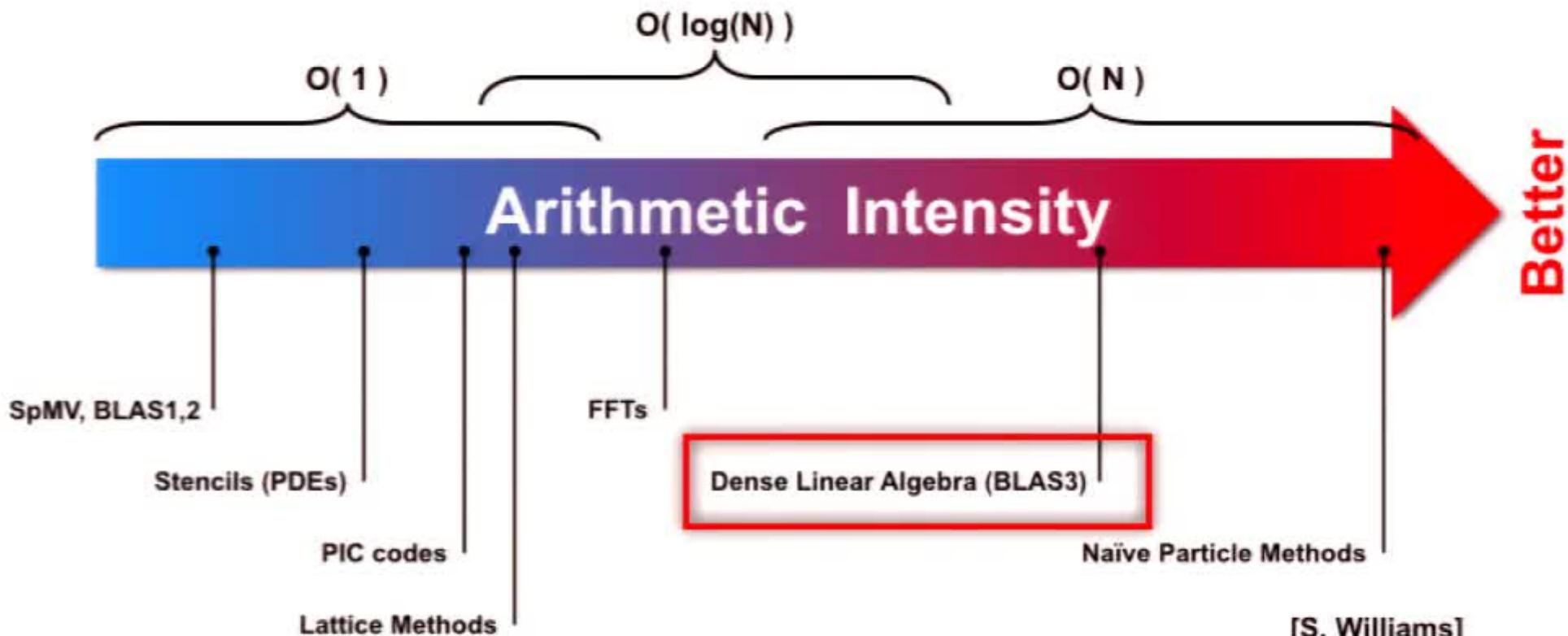
- Factor dense diagonal blocks (ZGETRF)
- Scale dense interface blocks (ZGETRS)
- Update dense Schur submatrices (ZGEMM)

- Parallelism

- Concurrent dense BLAS3 operations
- Fine-grained parallelism via GPU/Xeon Phi
- “Right sizes” dense matrices



Arithmetic Intensity

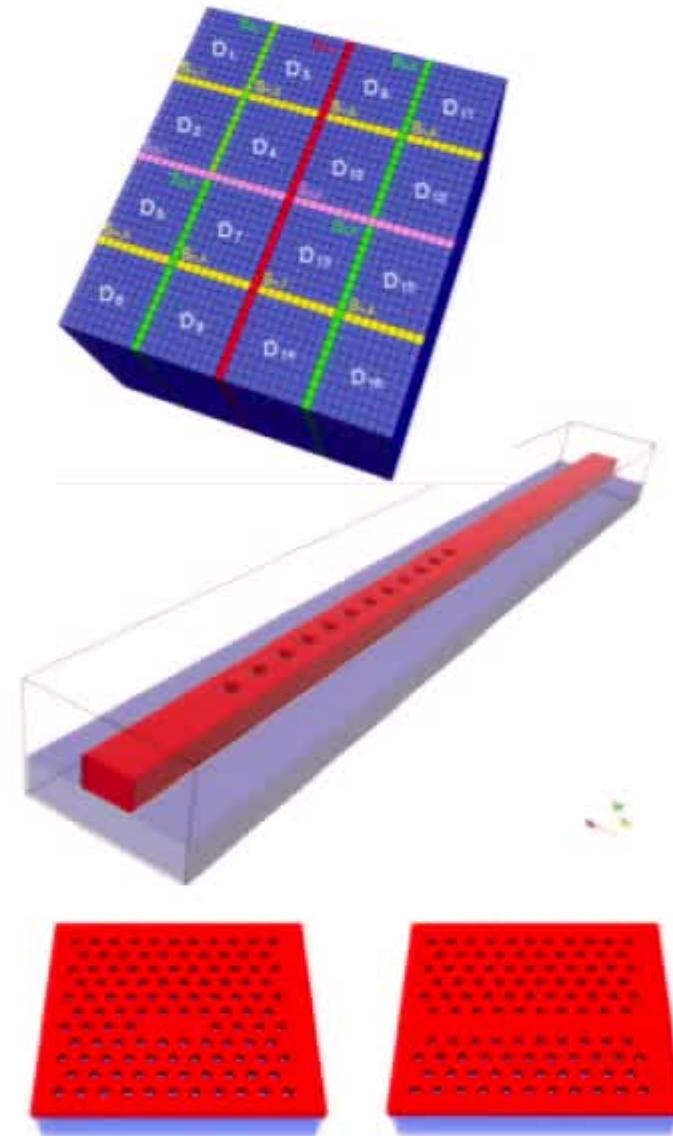


- Arithmetic Intensity $\sim (\text{Total Flops}) / (\text{Total DRAM Bytes})$
- Example: dense matrix-matrix multiplication: $(N^3 \text{ flops}) / (N^2 \text{ memory})$
- Higher arithmetic intensity \sim better locality \sim higher chance to achieve machine peak

Memory/Computation Saving



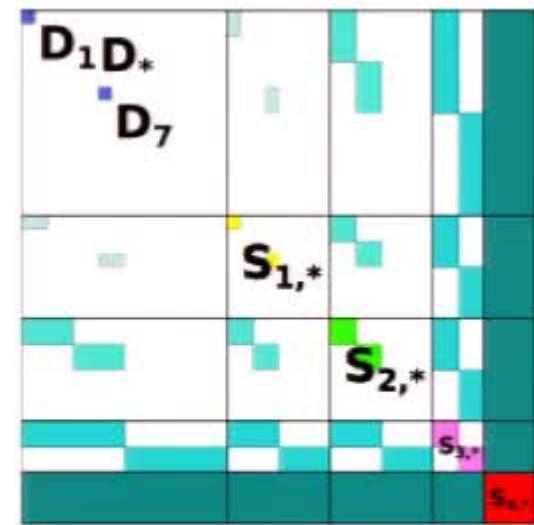
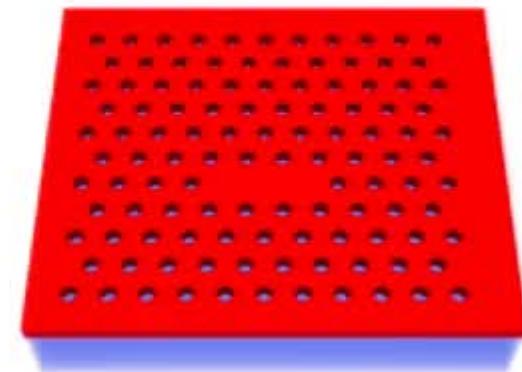
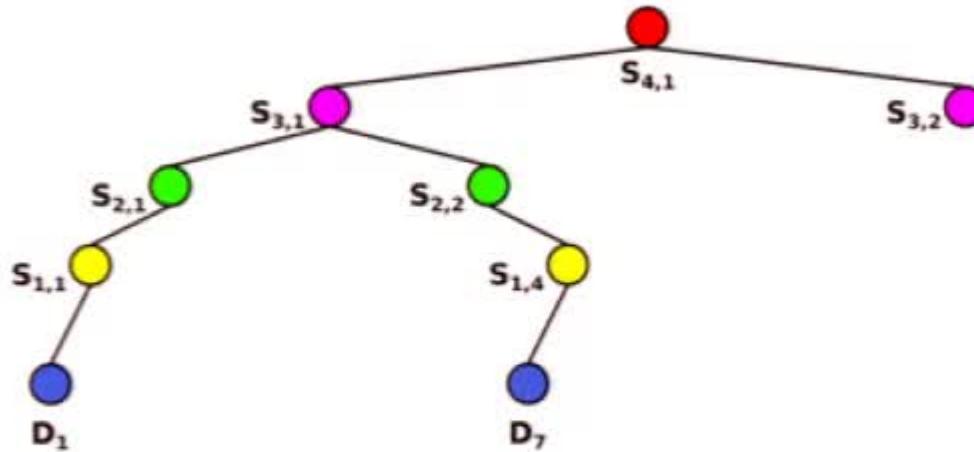
- Criteria of redundancy
 - Identical dimensions
 - Identical physical properties (e.g. periodic structures or homogeneous material)
 - Identical discretization (e.g. No PML and uniform grids)
 - Identical grid index ordering



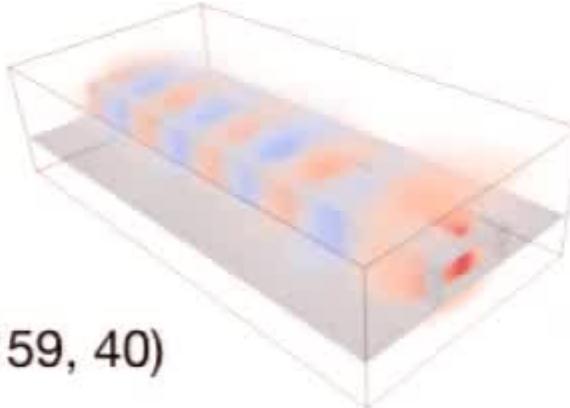
Memory/Computation Savings



- Identical redundant matrix blocks removal
- For defected parts
 - D_7 is different from other subdomains
 - No redundancy to its parent separators



Memory for Dielectric Waveguide



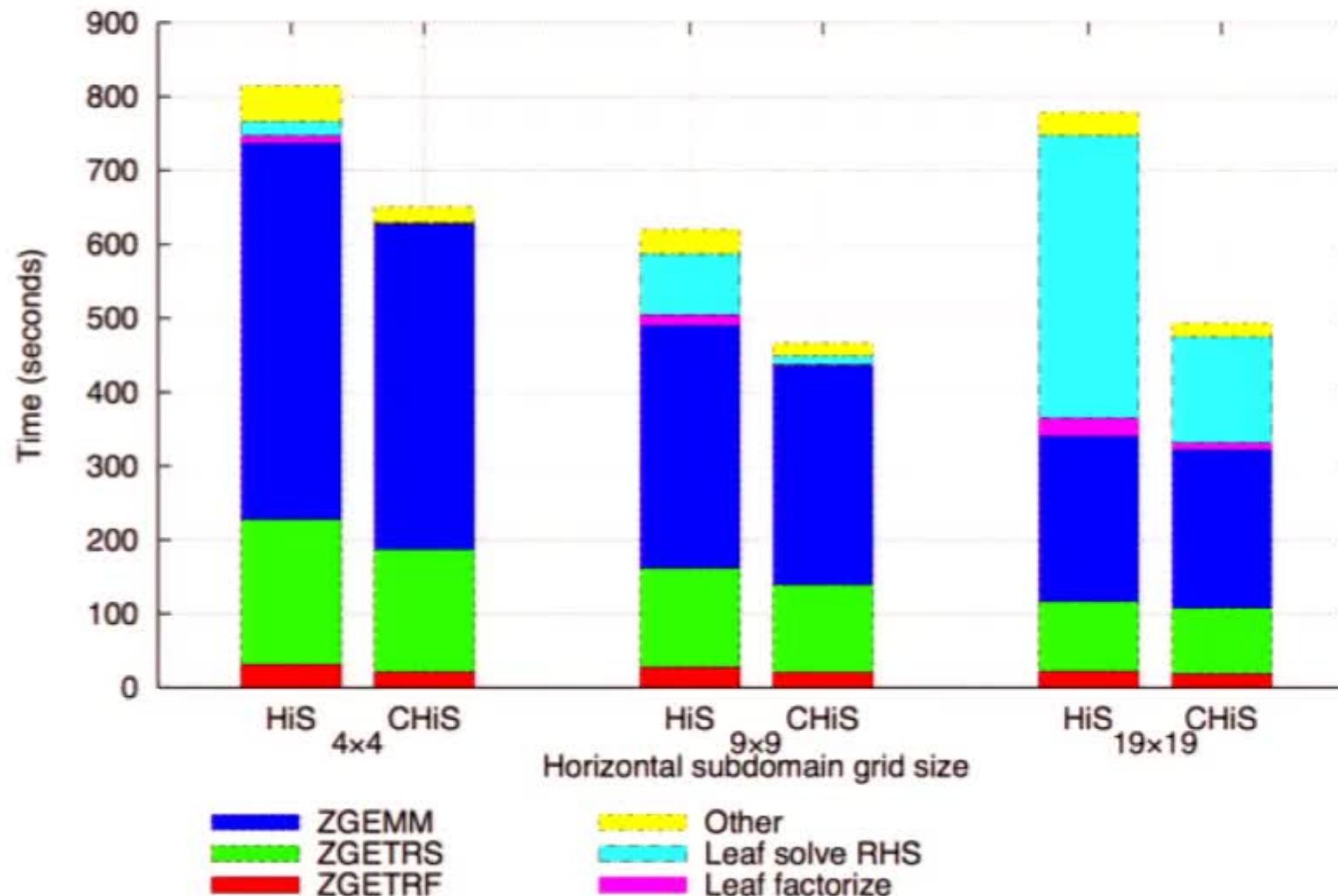
(79, 159, 40)

	Peak Memory	Number of non-redundant elements									
		D_*	$S_{1,*}$	$S_{2,*}$	$S_{3,*}$	$S_{4,*}$	$S_{5,*}$	$S_{6,*}$	$S_{7,*}$	$S_{8,*}$	$S_{9,*}$
HiS _{4×4}	115 GB	512	256	128	64	32	16	8	4	2	1
CHiS _{4×4}	75 GB	40	30	18	12	12	6	6	3	2	1
HiS _{9×9}	89 GB	128	64	32	16	8	4	2	1	0	0
CHiS _{9×9}	64 GB	18	12	12	6	6	3	2	1	0	0
HiS _{19×19}	68 GB	32	16	8	4	2	1	0	0	0	0
CHiS _{19×19}	52 GB	12	6	6	3	2	1	0	0	0	0
UMFPACK	245 GB										

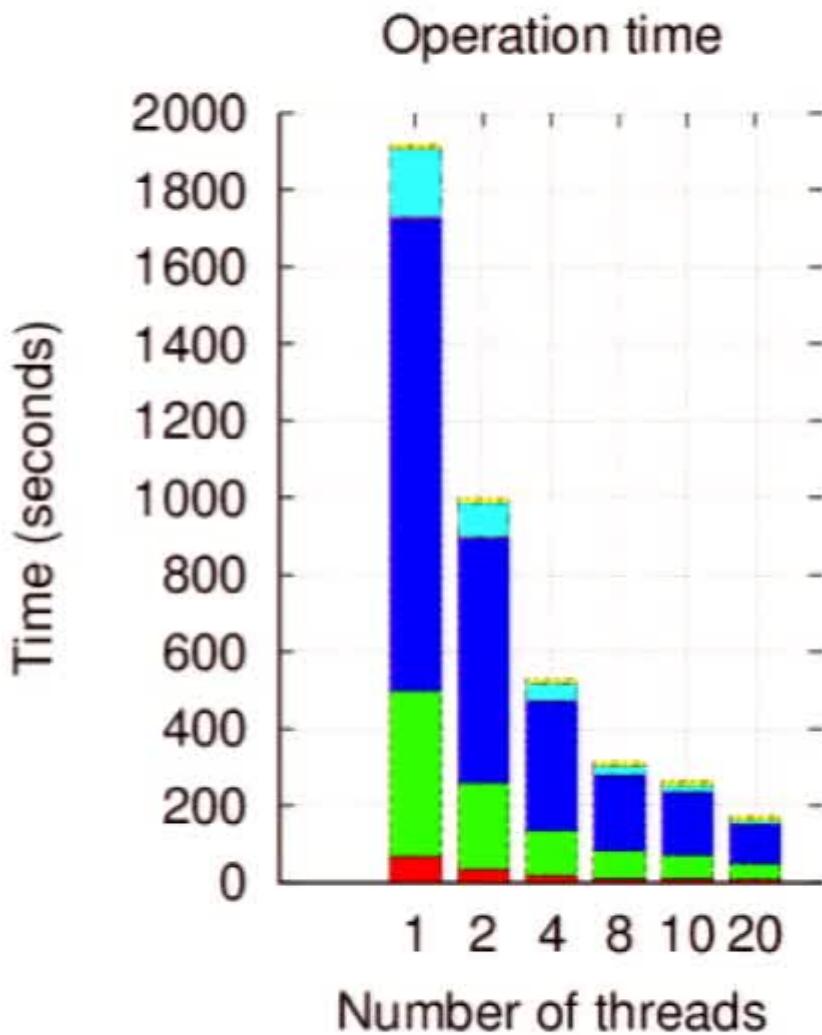
Timing Breakdown Analysis



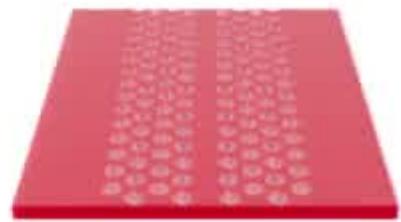
Dual Intel E5-2650 v3 CPUs + 256 GB memory



Factorization Time for PC Slab



- Other
- Leaf solve RHS
- Leaf factorize
- ZGEMM
- ZGETRS
- ZGETRF



- Matrix dimension: $95 \times 289 \times 20 = 1,362,300$
- Intel E5-2650 v3 (10 cores) x 2
- 256 GB DDR4-2133 main memory
- Intel Parallel Studio XE 2015 Cluster edition with ICC 15.0
- Flags: -O3 -openmp
- OS: CentOS 6.5
- Libraries: MKL 11.2, PARDISO (included in Intel MKL)

ZGEMM on CPU/GPU for PC Slab



- Projected GPU acceleration performance of dense linear algebra
 - Acceleration on ZGETRF, ZGETRS, and ZGEMM
 - Dual E5-2650 v3 \approx 490 GFLOPS
 - Dual NVIDIA K40 only: 1000 ~ 2400 GFLOPS
 - Expect more than 3x speedup with tuning and hybrid computing

