



UNIVERSITÀ
DEGLI STUDI
DI PADOVA

Dept. of Civil, Environmental and
Architectural Engineering

Numerical Modeling of Fault Activation and Induced Seismicity in Gas Storage Reservoirs: The Netherlands Case

*Massimiliano Ferronato, Andrea Franceschini, Carlo Janna, Pietro
Teatini, Claudia Zoccarato*

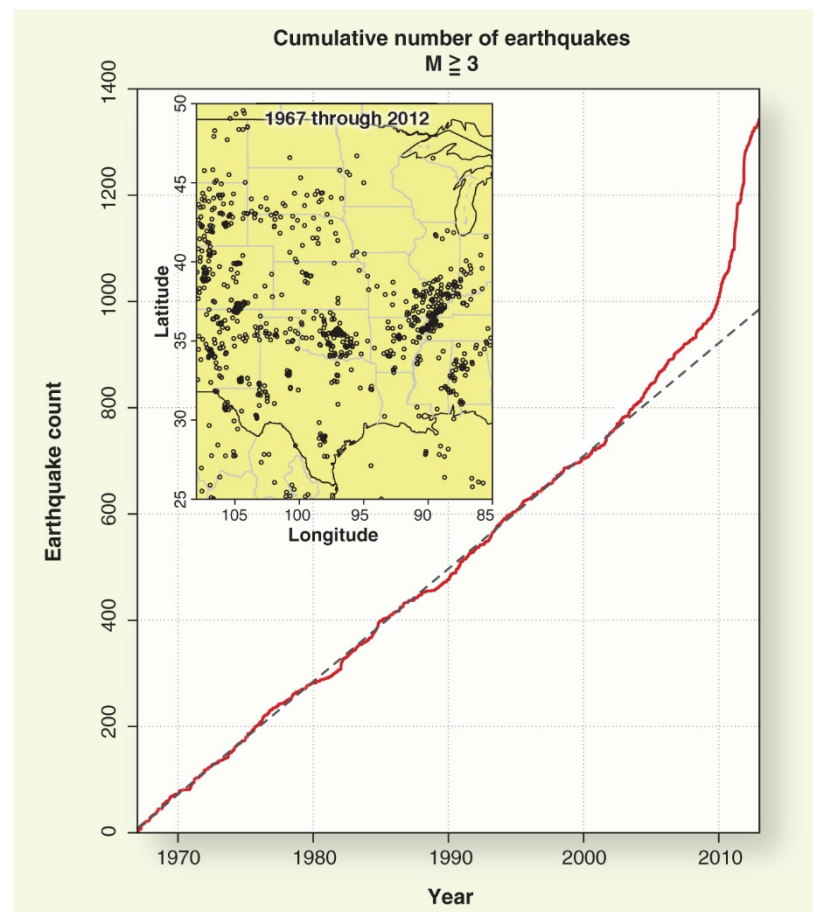
SIAM Conference on Mathematical and Computational Issues in the
Geosciences (GS19)

March 11-14, 2019 – Houston, TX



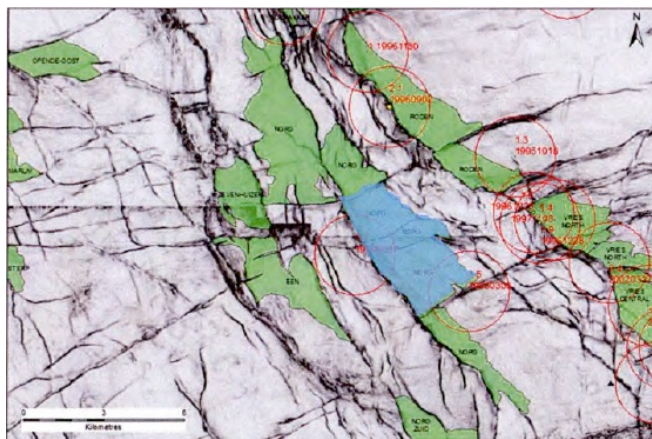
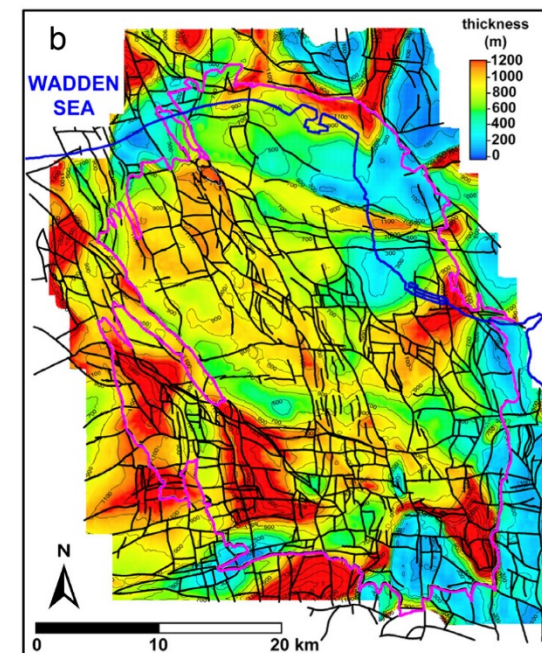
- ❑ Introduction: geohazards associated with Underground Gas Storage (UGS) activities
- ❑ The Netherlands case: conceptual model for the inception of fault movements
- ❑ Mathematical and numerical model: numerical discretization of fault mechanics by Lagrange multipliers
- ❑ Numerical results: factors that can enhance the probability of fault reactivation
- ❑ Conclusions: preliminary practical indications

- ❑ The interest in developing UGS projects is continuously increasing worldwide
 - May 2015: over 270 plants in Europe, 400 in the US
- ❑ Geohazards associated with UGS activities:
 - Formation integrity
 - Leakage from the reservoir
 - Land motion
 - Induced and/or triggered seismic events
- ❑ Coping with such issues is necessary for health and safety as related to public perception, economic risk and environmental impact



from Ellsworth (Science, 2013)

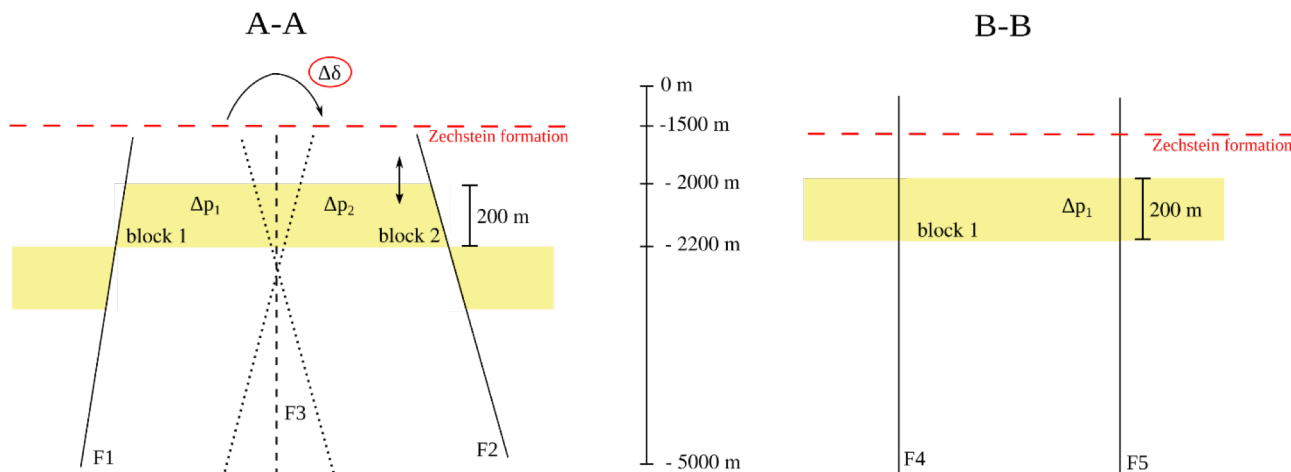
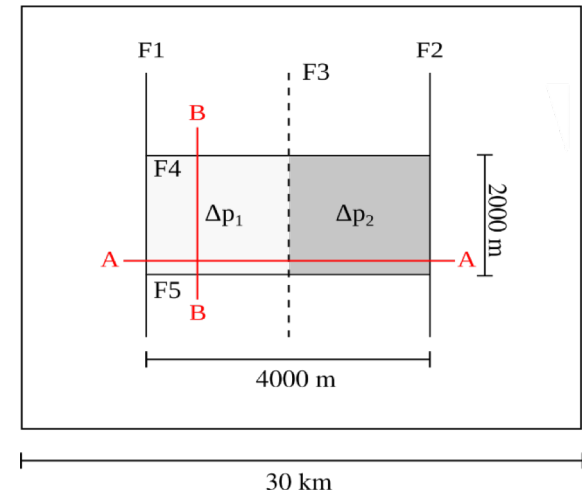
- ❑ Cases of seismic activity have been recently recorded in a few UGS plants in the Netherlands
 - Primary Production
 - Cushion Gas injection
 - Storage activities
- ❑ Highly compartmentalized and fractured reservoirs in stiff rocks overlain by salt deposits



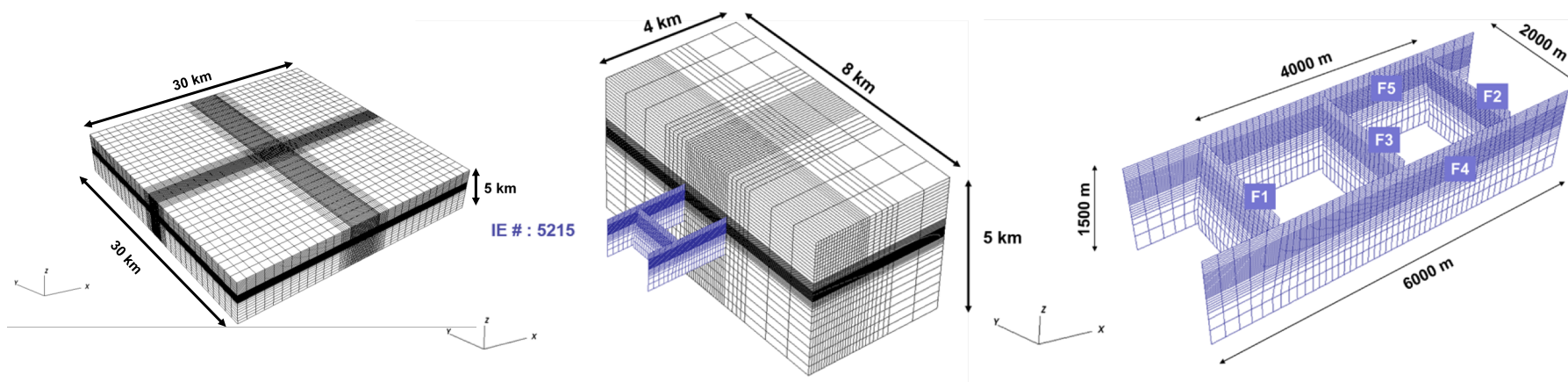
Develop numerical models to simulate the possible inception of fault motion also in «unexpected» configurations

□ Schematic geometry representative of the typical configuration of Dutch UGS fields:

- Independent blocks with different pressure variations
- Bounding vertical and sub-vertical faults
- Viscous salt formations on top of the reservoir

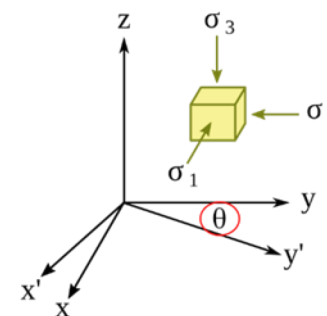


- Geomechanical parameters typical of Dutch UGS formations
- Almost isotropic initial stress state ($M_1=0.74$, $M_2=0.83$) with principal directions oriented like the bounding faults



IE # : 5215

LAYER	DENSITY (kg/m ³)	YOUNG MODULUS (GPa)	POISSON RATIO
Overburden	2200	10.0	0.25
Zechstein Salt	2100	40.0	0.3
Reservoir (Upper Rotliegend)	2400	11.0	0.15
Underburden	2600	30	0.2



Initial stress regime:

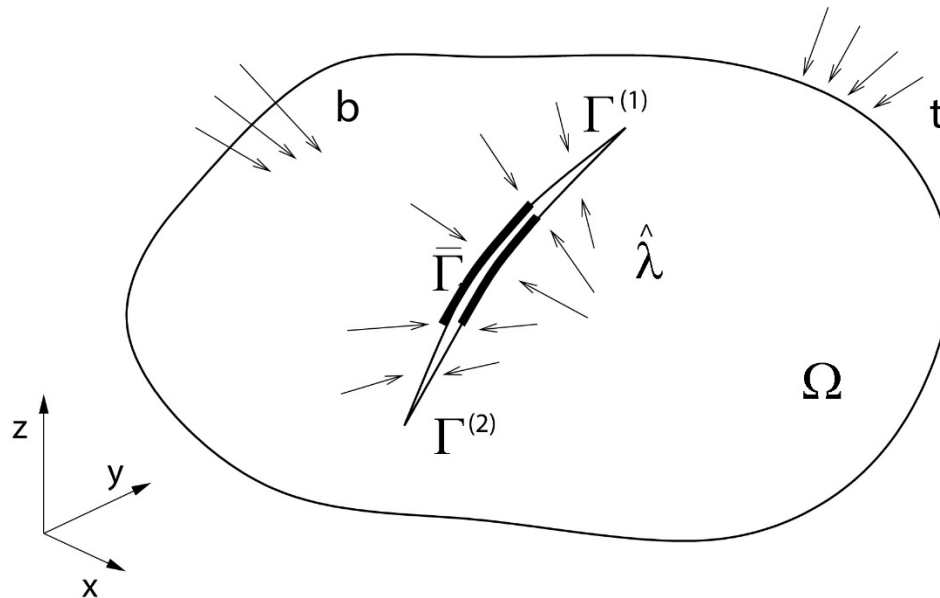
$$\sigma_h = \sigma_1 = M_1 * \sigma_3$$

$$\sigma_H = \sigma_2 = M_2 * \sigma_3$$

$$\sigma_v = \sigma_3$$



- ❑ Modeling fault/fracture mechanics involves a number of numerical open issues
- ❑ Available numerical approaches:
 1. Continuous Finite Elements with a different rheology, e.g. [Rutqvist et al. 2008]
 - Ease of implementation
 - Inability to describe slippage and opening
 2. Interface frictional elements by penalties, e.g. [Beer, 1985; Cescotto & Charlier, 1993; Juanes et al., 2002; Ferronato et al., 2008]
 - Definiteness preservation, controlled number of DoFs
 - Ill-conditioning, instability, non-linear convergence difficulties
 3. Lagrange multipliers, e.g. [Aagaard et al., 2013; Jha & Juanes, 2014; Franceschini et al. 2016]
 - Mathematically robust prescription of constraints
 - Increase of DoFs, saddle-point problem



- ❑ A fracture is a discontinuity within a 3D porous body made of a pair of friction surfaces in contact each other
- ❑ The surfaces can't penetrate and continuity is preserved if:

$$\tau_s \leq \tau_L = c - \sigma_n \tan \varphi, \quad \sigma_n < 0$$

- ❑ The Mohr-Coulomb criterion defines τ_L , but gives no indication as to the direction of the limiting shear vector \mathbf{t}_L
- ❑ According to the Principle of Maximum Plastic Dissipation, \mathbf{t}_L is such that the friction work W_f is maximum, i.e. is parallel to the slip vector \mathbf{u}_r

$$\mathbf{t}_L = \tau_L \frac{\mathbf{u}_r}{\|\mathbf{u}_r\|_2}$$



- From a mathematical standpoint, the problem requires the solution of a set of governing PDEs:

$\nabla \cdot \boldsymbol{\sigma} - \mathbf{b} = 0$	$\forall \mathbf{x} \in \Omega \setminus \Gamma_f$	Momentum balance
$\mathbf{u} = \bar{\mathbf{u}}$	$\forall \mathbf{x} \in \Gamma_u$	Boundary displacement
$\boldsymbol{\sigma} \cdot \mathbf{n} = \bar{\mathbf{t}}$	$\forall \mathbf{x} \in \Gamma_\sigma$	Boundary traction
$\boldsymbol{\sigma} \cdot \mathbf{n}_f^- = \boldsymbol{\sigma} \cdot \mathbf{n}_f^+ = \mathbf{t}$	$\forall \mathbf{x} \in \Gamma_f$	Traction continuity

subject to normal contact conditions along Γ_f :

$$t_N = \mathbf{t} \cdot \mathbf{n}_f \leq 0, \quad g_N = \llbracket \mathbf{u} \rrbracket \cdot \mathbf{n}_f \geq 0, \quad t_N g_N = 0$$

and Coulomb frictional contact conditions along Γ_f :

$$\Phi = \|\mathbf{t}_T\|_2 - (c - t_N \tan \varphi) \leq 0, \quad \mathbf{g}_T - \alpha \frac{\mathbf{t}_T}{\|\mathbf{t}_T\|_2} = 0,$$
$$\alpha \geq 0, \quad \Phi \alpha = 0$$



- The discretization is obtained with a mixed finite element approach where displacement \mathbf{u} and traction \mathbf{t} on the fault are the main unknowns
- Find $\{\mathbf{u}, \mathbf{t}\} \in \mathcal{S}_u(\Omega) \times \mathcal{S}_t(\Gamma_f)$ such that for all $\{\mathbf{v}, \boldsymbol{\mu}\} \in \mathcal{V}_u(\Omega) \times \mathcal{V}_t(\Gamma_f)$

$$\int_{\Omega \setminus \Gamma_f} \nabla^S \mathbf{v} : \boldsymbol{\sigma} \, d\Omega - \int_{\Omega \setminus \Gamma_f} \mathbf{v} \cdot \mathbf{b} \, d\Omega - \int_{\Gamma_\sigma} \mathbf{v} \cdot \bar{\mathbf{t}} \, d\Gamma + \int_{\Gamma_f} \llbracket \mathbf{v} \rrbracket \cdot \mathbf{t} \, d\Gamma = \mathbf{0}$$

$$\int_{\Gamma_f} \boldsymbol{\mu} \cdot \llbracket \mathbf{u} \rrbracket \, d\Gamma = 0$$

- Introducing as usual the discrete finite element spaces we obtain:

$$\int_{\Omega \setminus \Gamma_f} \nabla^S \mathbf{N}_u : \boldsymbol{\sigma}^h \, d\Omega + \int_{\Gamma_f} \llbracket \mathbf{N}_u \rrbracket \cdot \mathbf{t}^h \, d\Gamma = \mathbf{f}$$

$$\int_{\Gamma_f} \mathbf{N}_t \cdot \llbracket \mathbf{u}^h \rrbracket \, d\Gamma = 0$$



- The integrals along the fractures Γ_f are computed as the sum of the contributions arising from the portions operating in the «stick» and «slip» modes

$$\int_{\Gamma_f^{st}} \llbracket \mathbf{N}_u \rrbracket \cdot \mathbf{t}^h d\Gamma + \int_{\Gamma_f^{sl}} \llbracket \mathbf{N}_u \rrbracket \cdot (\mathbf{n}_f \otimes \mathbf{n}_f) \cdot \mathbf{t}^h d\Gamma + \int_{\Gamma_f^{sl}} \llbracket \mathbf{N}_u \rrbracket \cdot \mathbf{t}_T^{*,h} d\Gamma$$

$$\int_{\Gamma_f^{st}} \mathbf{N}_t \cdot \llbracket \mathbf{u}^h \rrbracket d\Gamma + \int_{\Gamma_f^{sl}} \mathbf{N}_t \cdot (\mathbf{n}_f \otimes \mathbf{n}_f) \cdot \llbracket \mathbf{u}^h \rrbracket d\Gamma$$

- The problem is highly non-linear because also the «stick» and «slip» portions of the fracture are unknown
- The non-linear problem is finally addressed by a Newton-Raphson scheme



- A weakening model is implemented to account for the variation of the friction angle from static to dynamic conditions at the fault slipping:

$$\varphi = \varphi_s + \frac{\varphi_d - \varphi_s}{d_c} \|\mathbf{g}_T\|_2, \quad \|\mathbf{g}_T\|_2 \leq d_c$$

$$\varphi = \varphi_d, \quad \|\mathbf{g}_T\|_2 > d_c$$

- A Maxwell model is used to simulate the viscous behaviour of the top salt Zechstein formation:

$$\dot{\boldsymbol{\varepsilon}}_v = V(\mu)\boldsymbol{\sigma}, \quad \mu = 10^{17} \text{ Pa} \cdot \text{s}$$

- The faults are partially sealing and are characterized by an inner pressure variation p_f



- The discrete Jacobian is a large, sparse matrix with a generalized saddle-point structure:

$$J(\mathbf{u}, \mathbf{t}) = \begin{bmatrix} K(\mathbf{u}) + E(\mathbf{u}, \mathbf{t}) & C - F(\mathbf{u}, \mathbf{t}) \\ C^T & 0 \end{bmatrix}$$

Classic FE stiffness
+
Maximum plastic
dissipation

Fault
equilibrium
+
Maximum
plastic
dissipation

Fault
congruence

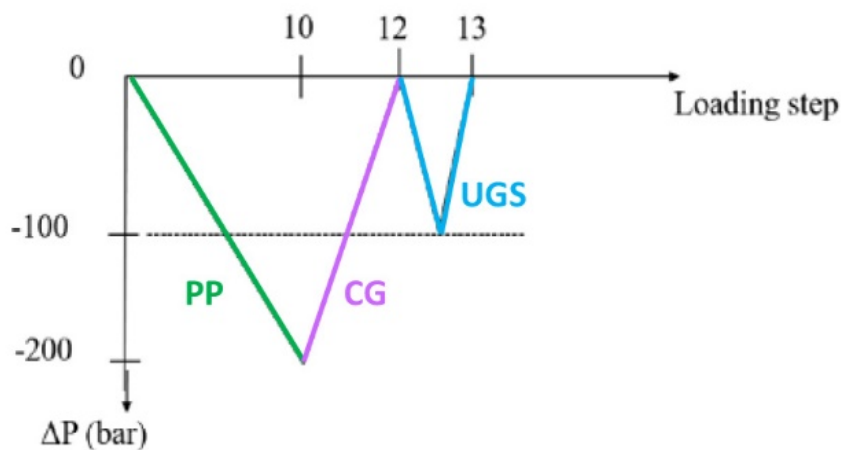
0

$$K = \int_{\Omega} \nabla^s \mathbf{N}_u : \frac{\partial \boldsymbol{\sigma}}{\partial \boldsymbol{\epsilon}} : \nabla^s \mathbf{N}_u d\Omega$$

$$E = \int_{\Gamma_f^{sl}} \llbracket \mathbf{N}_u \rrbracket \cdot \frac{\partial \mathbf{t}_T^*}{\partial \mathbf{u}} \cdot \llbracket \mathbf{N}_u \rrbracket d\Gamma$$

$$C = \int_{\Gamma_f^{st}} \llbracket \mathbf{N}_u \rrbracket \cdot \mathbf{N}_t d\Gamma + \int_{\Gamma_f^{sl}} \llbracket \mathbf{N}_u \rrbracket \cdot (\mathbf{n}_f \otimes \mathbf{n}_f) \cdot \mathbf{N}_t d\Gamma$$

$$F = \int_{\Gamma_f^{sl}} \llbracket \mathbf{N}_u \rrbracket \cdot \frac{\partial \mathbf{t}_T^*}{\partial \mathbf{u}} \cdot \mathbf{N}_t d\Gamma$$



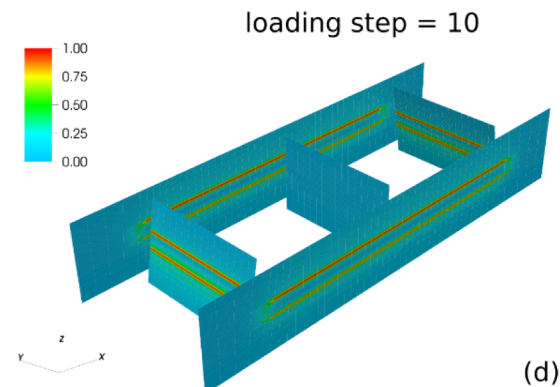
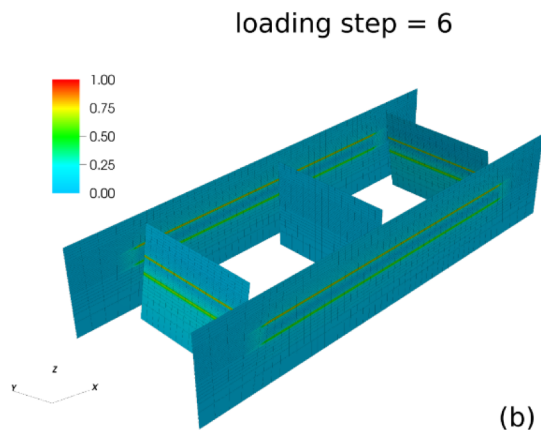
- ❑ Pressure history prescribed in the UGS field
- ❑ A loading step corresponds to a 1-year time interval
- ❑ Critical index:

$$\chi = \frac{\|\mathbf{t}_T\|_2}{\|\mathbf{t}_T^*\|_2} \leq 1$$

- ❑ Sensitivity analysis on different configuration parameters:
 - Block offset (0-200 m)
 - Biot coefficient (0.6 – 1.0)
 - Central fault dip (-25° - 25°)
 - Initial stress regime ($M_1=M_2=0.4$)
 - Reservoir and caprock stiffness
 - Pressure variation (0-200 bar)
 - Fault properties
 - Fault pressure (0-200 bar)

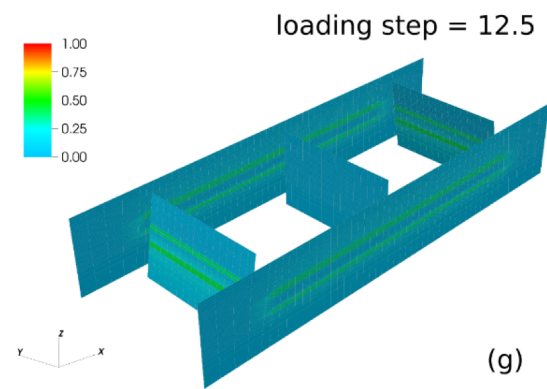
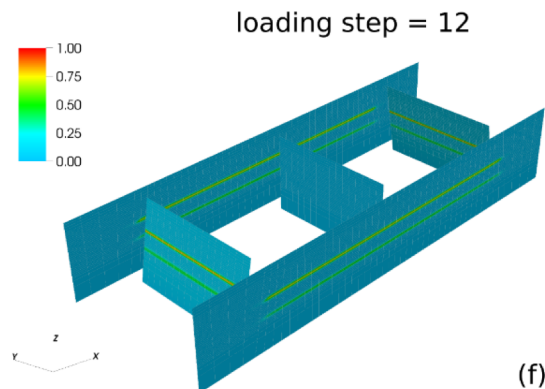


Primary
Production

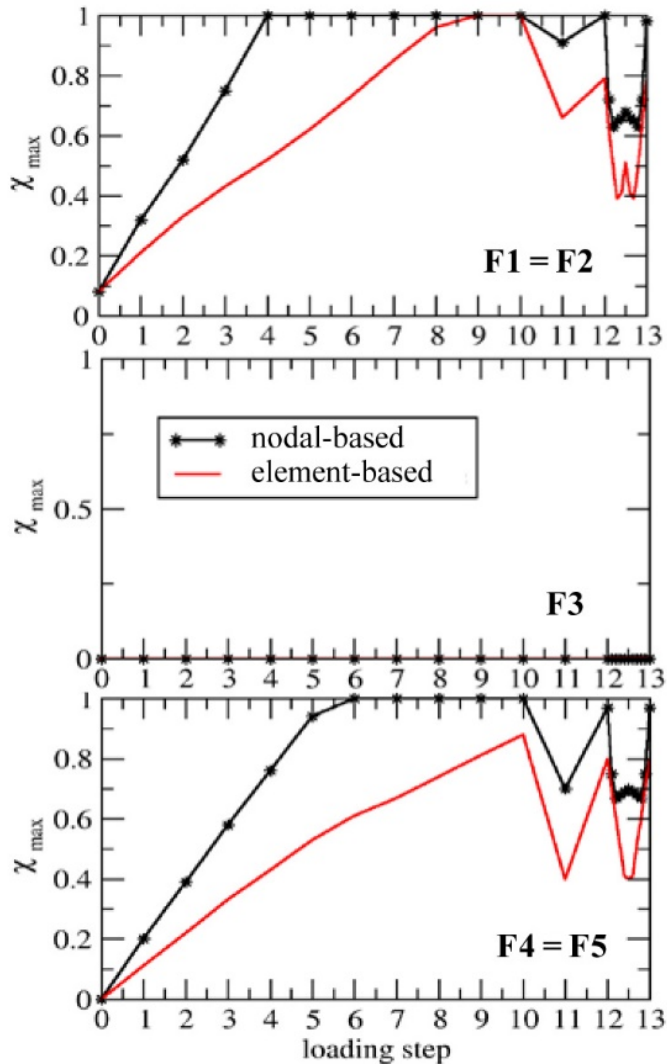


Primary
Production

Cushion
Gas

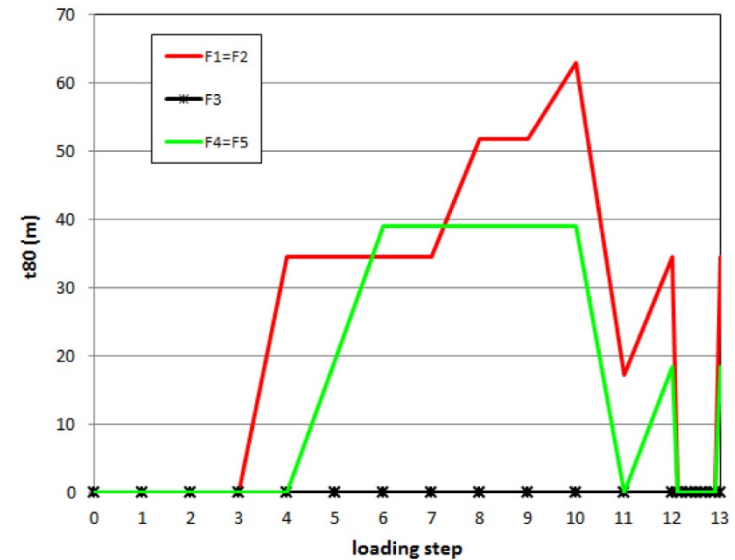


Gas
Storage



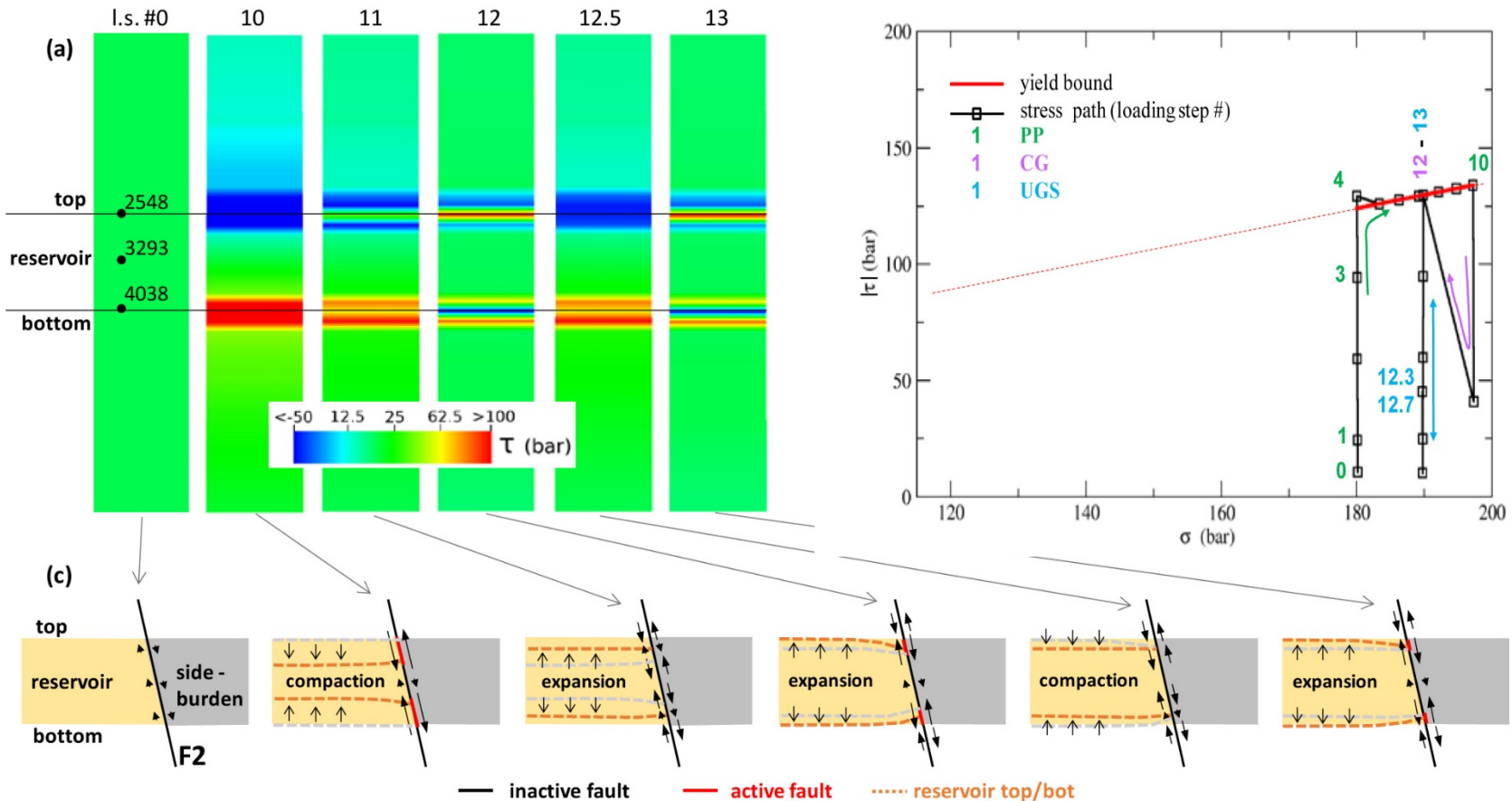
□ Critical index on the bounding faults vs the loading steps

□ Fault thickness with $\chi > 0.8$





Stress on a bounding fault during Primary Production, Cushion Gas e Gas Storage





- ❑ Modifying the mechanical and geometrical properties can anticipate or delay the fault activation, or increase the amount of slippage and the activated area
 - the initial stress regime can play a major role: decreasing significantly the horizontal principal components favors an early fault reactivation, with a large area critically stressed and significant sliding
 - factors increasing the activation risk: (1) a reduced friction angle; (2) an offset between producing blocks; (3) stiffness contrasts between reservoir, caprock, sideburden, and underburden; (4) uneven pressure change in adjacent compartments
 - most critical configuration: 200-m compartment offset, a relatively small friction angle, and a viscous caprock



- ❑ A 3D modeling study has been developed on a conceptual configuration representative of the Dutch UGS fields to evaluate how and when faults bounding a compartmentalized reservoir can be reactivated during the UGS activities
- ❑ The fault mechanics is simulated with the aid of a Lagrangian formulation with a non-linear weakening behavior for the friction angle, a viscous salt caprock and partially sealing faults
- ❑ The investigation has been carried out in detail on a reference benchmark and by changing mechanical and geometrical configurations in an extensive sensitivity analysis



- ❑ Fault reactivation may occur “unexpectedly” during Cushion Gas and Underground Gas Storage stages, following more expected reactivations during Primary Production
- ❑ Activation during Primary Production leads to a stress redistribution and a new "equilibrated" configuration that is re-loaded, in the opposite direction, when the pressure variation changes the sign at the Cushion Gas injection
- ❑ The settings more prone to activation during Primary Production are also the most critical ones during Cushion Gas injection and Underground Gas Storage

- ❑ A preliminary indication may rely on limiting the pressure recovery during the storage operations if a reactivation has been already experienced during the Primary Production



UNIVERSITÀ
DEGLI STUDI
DI PADOVA

**Dept. of Civil, Environmental and
Architectural Engineering**

Thank you for your attention

**SIAM Conference on Mathematical and Computational Issues in the
Geosciences (GS19)**

March 11-14, 2019 – Houston, TX



- ❑ Solving a linear problem with a saddle-point matrix is a common issue in several applications, e.g., flow problems in mixed form, coupled consolidation, Navier-Stokes equations, etc.
- ❑ The most effective approach proceeds as follows:

- Solve for the incremental displacements $\delta \mathbf{u}$ in the first equation

$$\delta \mathbf{u} = (K + E)^{-1} [\mathbf{f} - (C - F) \delta \boldsymbol{\lambda}]$$

- Replace $\delta \mathbf{u}$ in the second equation and solve for $\delta \boldsymbol{\lambda}$

$$\delta \boldsymbol{\lambda} = \left[C^T (K + E)^{-1} (C - F) \right]^{-1} C^T (K + E)^{-1} \mathbf{f}$$

- Approximate the application of the inverse of 1,1 block $(K+E)$ and the Schur complement S :

$$S = C^T (K + E)^{-1} (C - F)$$

- Use these equations to apply a «preconditioner» to J



- ❑ Numerical ingredients to define an effective preconditioner for a Krylov subspace solver:
 - Approximation of $(K+E)^{-1}$
 - Approximation of S and of S^{-1}
- ❑ Approximating the application of the inverse of $(K+E)$ is not an issue and can be done in several different ways, e.g., Incomplete Factorizations, Sparse Approximate Inverses, Algebraic Multigrid
- ❑ The difficulty lies in approximating the Schur complement S and its inverse
- ❑ Three approaches:
 - Factored Sparse Approximate Inverse (FSAI)
 - Block Diagonal Schur complement (BDS)
 - Least-Square Commutator (LSC)



Factored Sparse Approximate Inverse (FSAI)

- ❑ For the computation of S an explicit approximation of $(K+E)^{-1}$ is required
- ❑ Since $(K+E)$ is SPD, we can use an adaptive FSAI approximation [Janna et al. 2015]:

$$(K + E)^{-1} \approx GG^T \quad \Rightarrow \quad S \approx C^T GG^T (C - F) \approx H^T H$$

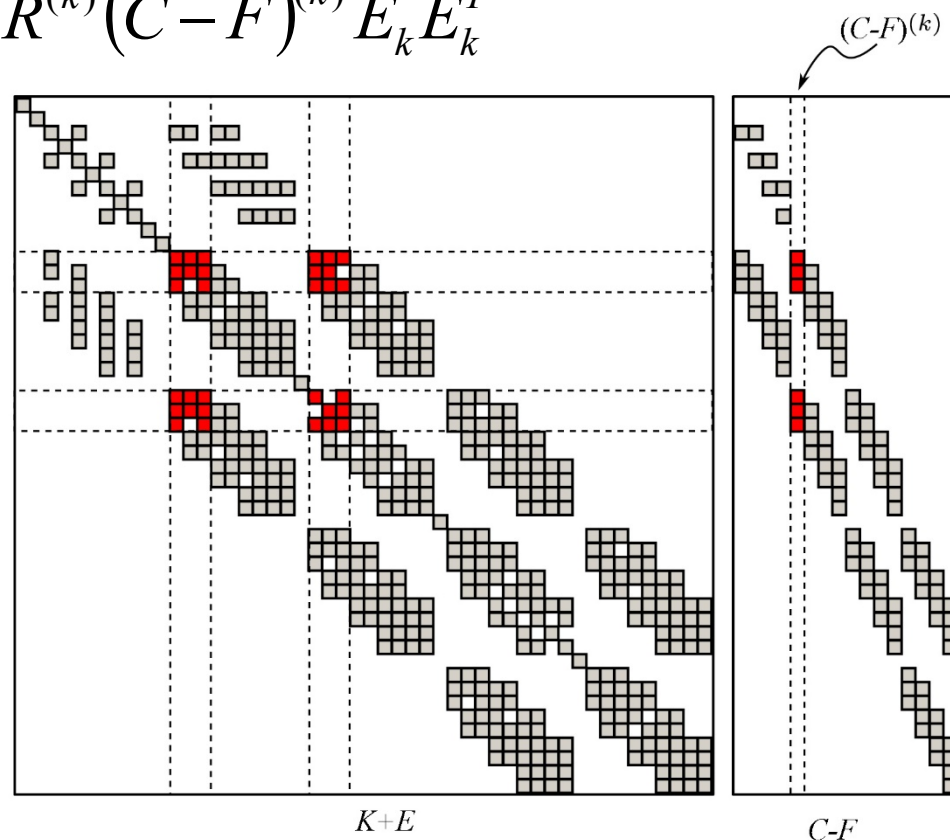
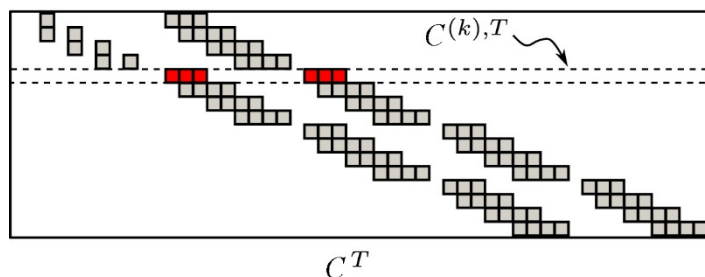
- ❑ The inverse of the Schur complement is then applied with a direct solver or another FSAI approximation

Block Diagonal Schur complement (BDS)

- ❑ From a physical point of view, $\delta\lambda$ ensures the continuity of the displacement through the fracture
- ❑ Such a continuity can be approximately prescribed on a node-by-node basis considering the local stiffness associated to each node only

- From a purely algebraic point of view, this means considering the restrictions $(K+E)^{(k)}$ of $(K+E)$ corresponding to the entries related to DoFs of the elements sharing the k -th node:

$$S \approx \sum_k C^{(k),T} R^{(k),T} K^{(k),-1} R^{(k)} (C-F)^{(k)} E_k E_k^T$$





Least-Square Commutator (LSC)

- This approach was originally introduced for Navier-Stokes problems in order to avoid the inversion of the 1,1 block for approximating the inverse of S [Elman et al. 2006]

- The objective is to find a commutator K_p such that:

$$(K + E)C \approx CK_p \quad \Rightarrow \quad K_p = (C^T C)^{-1} C^T (K + E)C$$

- From the previous approximation it follows also that:

$$CK_p^{-1} \approx (K + E)^{-1} C$$

so the inverse of the Schur complement can be written as:

$$\begin{aligned} S^{-1} &= \left[C^T (K + E)^{-1} (C - F) \right]^{-1} \approx \left[C^T (K + E)^{-1} C \right]^{-1} \approx \left[C^T CK_p^{-1} \right]^{-1} \approx \\ &\approx K_p (C^T C)^{-1} = \underline{(C^T C)^{-1} C^T (K + E)C (C^T C)^{-1}} \end{aligned}$$



- ❑ The advantage of the LSC approach is that the inverse of S is directly available using the blocks of J and $(K+E)^{-1}$ is not needed
- ❑ In our fault/fracture formulation, $C^T C$ is diagonal, so that the computation and application of its inverse is totally inexpensive
- ❑ The final Schur complement inverse can be also seen as:

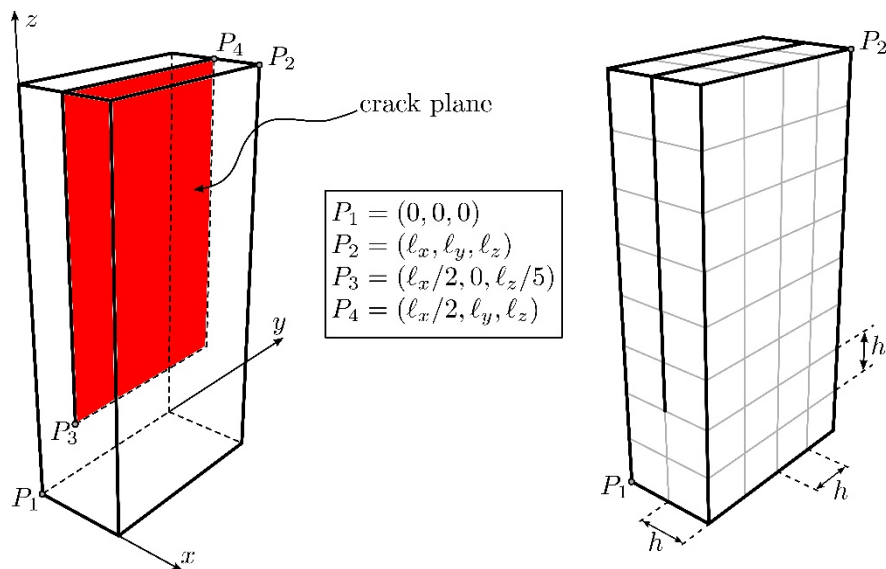
$$S^{-1} \approx C^+ (K + E)^{-1} C^{+,T}$$

where C^+ is the pseudo-inverse of C according to the Moore-Penrose definition

- ❑ *Theorem.* The eigenvalues λ of $S^{-1}S$, where S^{-1} is approximated by the LSC, are bounded from below by 1:

$$1 \leq \lambda \leq 1 + \|P\|^2 \|H\|$$

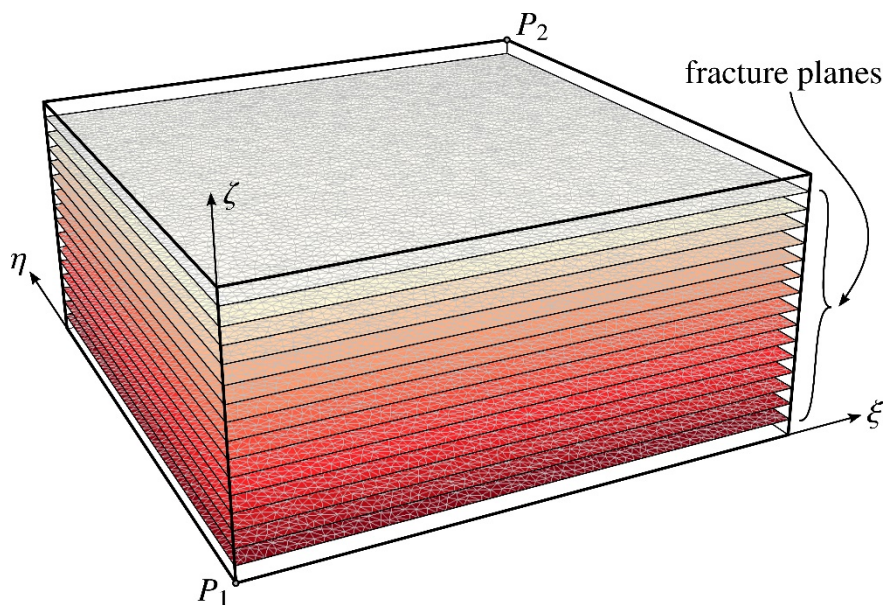
with P and H matrices that depend on the SVD decomposition of C and on K



h	# u dofs	# λ dofs
1/2	615	120
1/4	3,267	432
1/8	20,451	1,632
1/16	142,659	6,336
1/32	1,060,995	24,960

Iteration count
to converge
vs. the mesh
size

l/h	FSAI(5,.01)	FSAI(20,.01)	BDS	LSC
2	22	20	27	22
4	29	25	34	27
8	35	20	40	32
16	42	36	48	39
32	49	43	56	46



Multiple fractured elastic
medium with 15 discontinuities

# u dofs	# λ dofs	Ratio
379,983	167,799	0.442

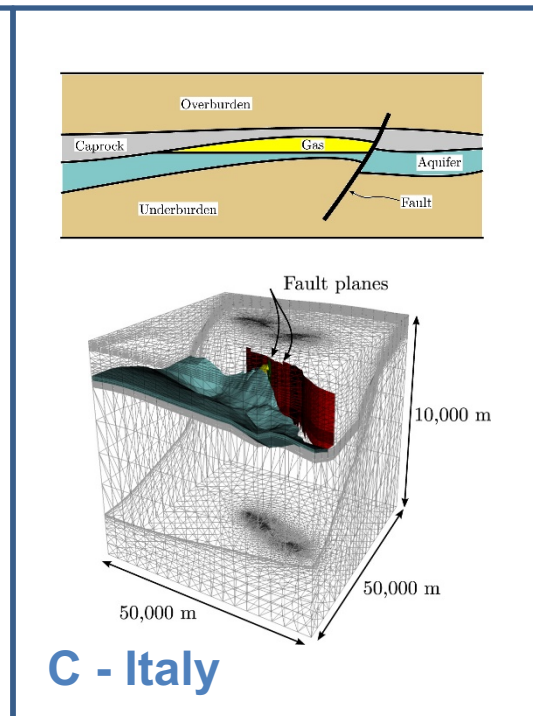
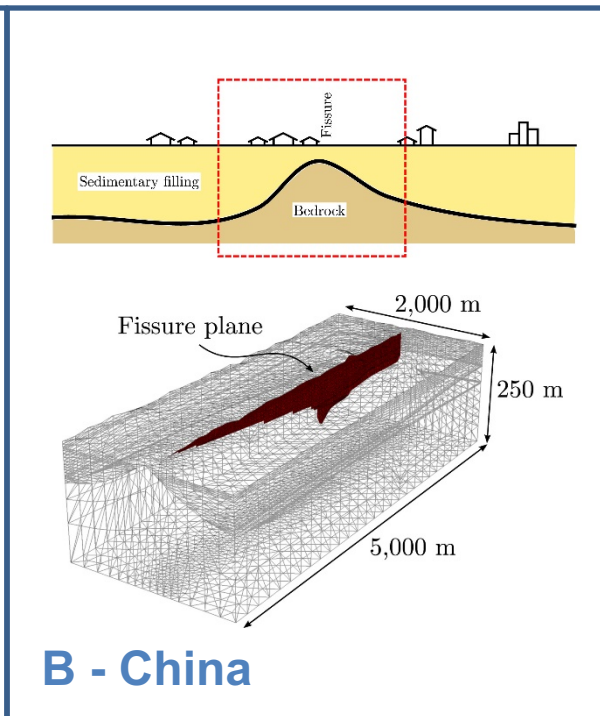
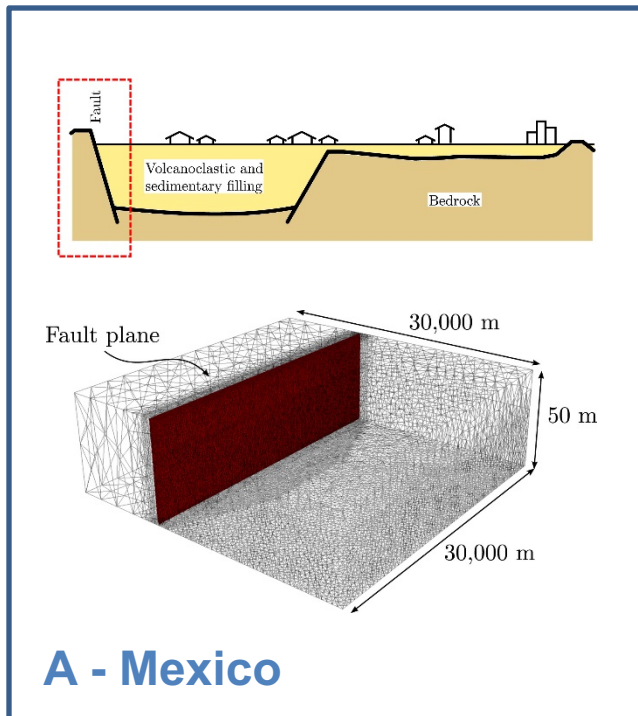
FSAI + DIR

FSAI + FSAI

BDS

LSC

$(K+E)^{-1}$	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]
FSAI	409	96.5	428	80.7	482	68.1	866	161.6
IC	229	73.5	227	67.9	574	80.6	407	77.8
exact	100	--	92	--	268	--	99	--



	# u dofs	# λ dofs	# tot dofs	Nnz(K+E)	Nnz(C-F)	Nnz(C ^T)
Case A	171,150	12,027	183,177	7,313,814	72,162	72,162
Case B	69,909	2,757	72,666	2,946,915	16,542	16,542
Case C	1,142,655	38,109	1,180,764	49,858,749	228,654	228,654



A - Mexico

FSAI + DIR

FSAI + FSAI

BDS

LSC

$(K+E)^{-1}$	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]
FSAI	123	9.6	126	9.1	255	13.0	175	12.3
IC	98	10.3	97	10.1	240	12.2	59	6.8
exact	35	--	34	--	126	--	23	--

B - China

FSAI + DIR

FSAI + FSAI

BDS

LSC

$(K+E)^{-1}$	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]
FSAI	72	1.7	73	2.0	371	5.5	77	2.6
IC	34	1.6	34	1.8	189	3.0	24	0.8
exact	21	--	24	--	129	--	12	--



C - Italy

FSAI + DIR

FSAI + FSAI

BDS

LSC

$(K+E)^{-1}$	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]	# iter	Time [s]
FSAI	496	156.5	623	172.7	--	--	768	287.1
IC	194	116.1	195	128.3	--	--	125	73.2
exact	50	--	50	--	276	--	35	--