

Input-Output Analysis of Channel Flows and Implications for Flow Control

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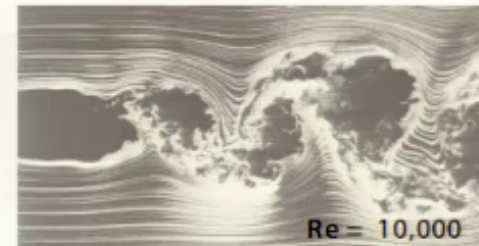
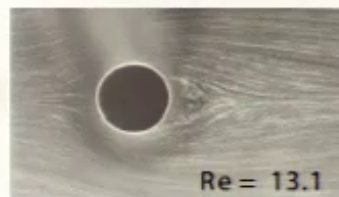
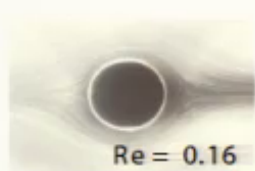


Transition & Turbulence as Natural Phenomena

All fluid flows *transition* (as $0 \xrightarrow{R} \infty$) from laminar to turbulent flows

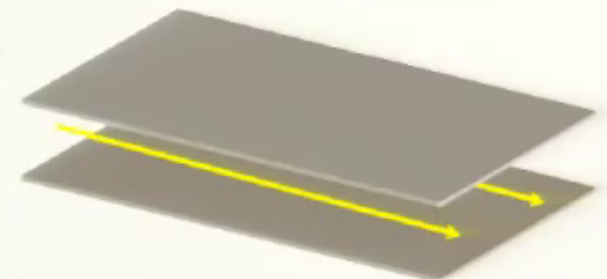
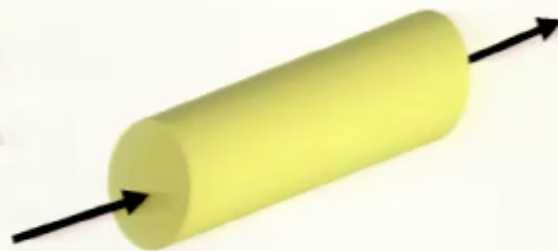
- Bluff bodies

dominant phenomenon: *separation*



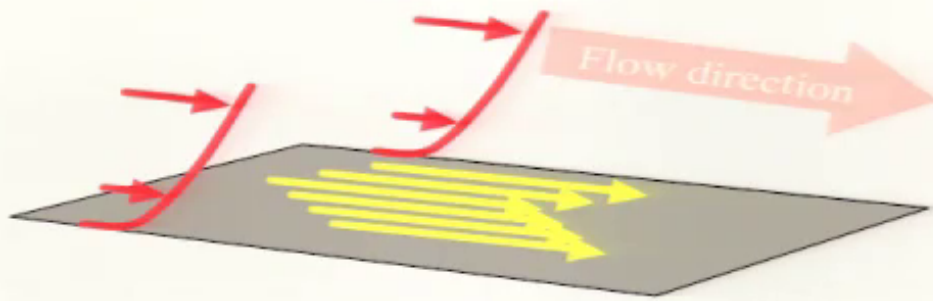
- Streamlined bodies

dominant phenomenon: *friction with walls*

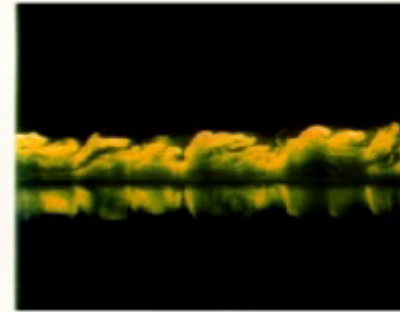


wall-bounded shear flows

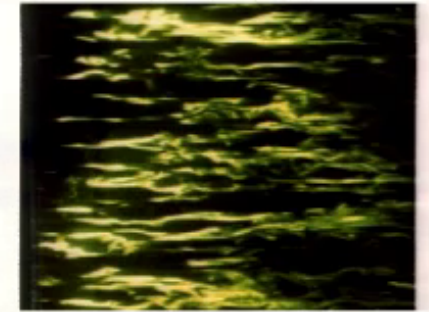
Boundary Layer Turbulence



boundary layer turbulence



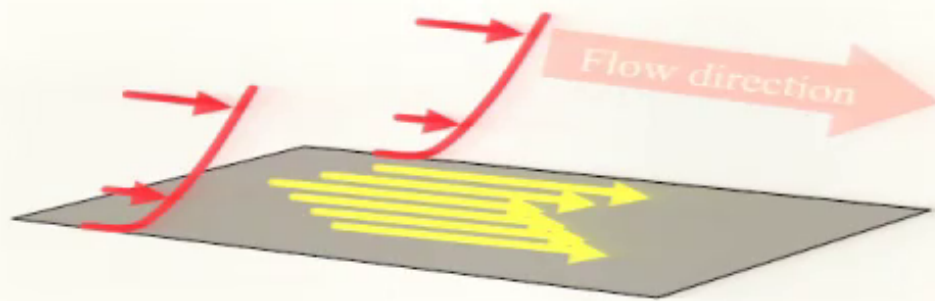
side view



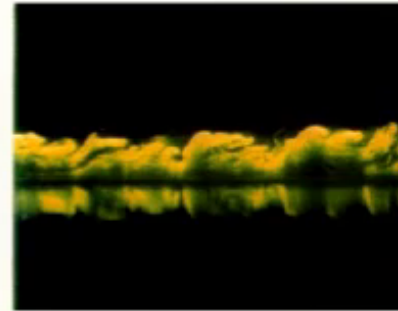
top view

- Transition & Turbulence in Boundary Layer and Channel Flows

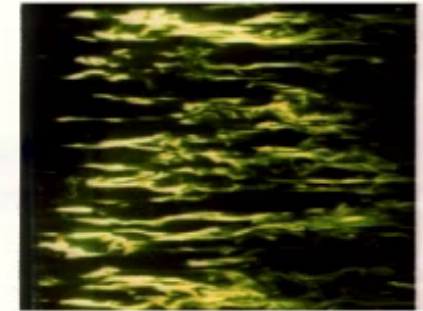
Boundary Layer Turbulence



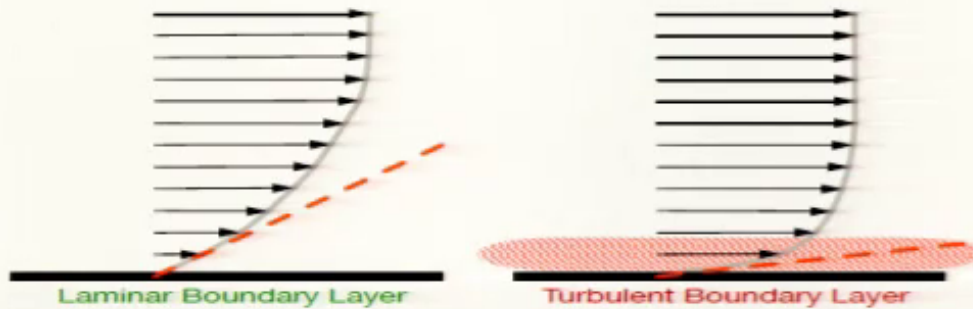
boundary layer turbulence



side view



top view



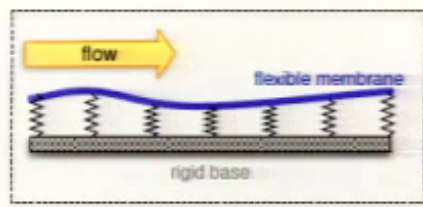
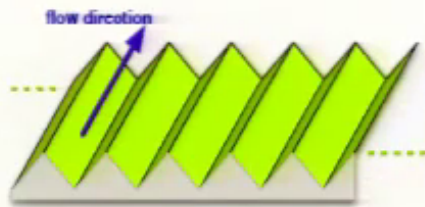
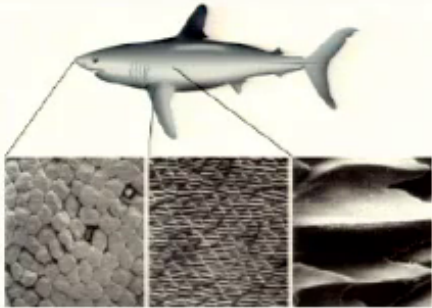
skin-friction drag: laminar vs. turbulent



- Transition & Turbulence in Boundary Layer and Channel Flows
- Technologically Important: *Skin-Friction Drag*

Control of Boundary Layer Turbulence

“passive” control



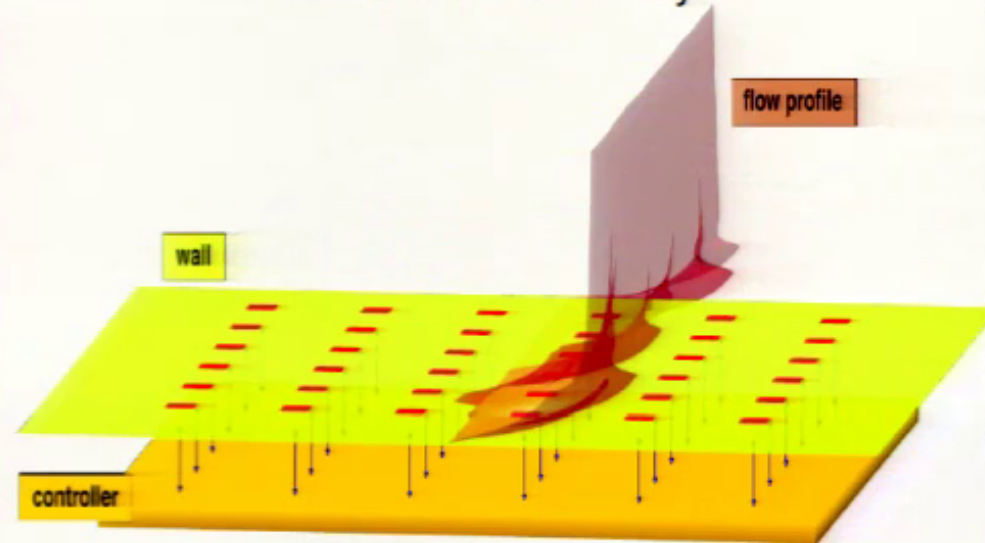
corrugated skin

compliant skin

● Other “open loop” schemes:

- ▶ Oscillating walls
- ▶ Body force traveling waves

active control with
sensor/actuator arrays



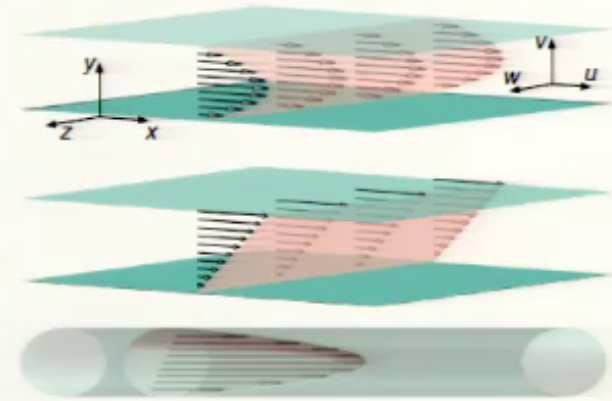
Caveat: *Plant's dynamics are not well understood*

obstacles { not only device technology
also: *dynamical modeling and control design*

Mathematical Modeling of Transition: Hydrodynamic Stability

The Navier-Stokes (NS) equations:

$$\begin{aligned}\partial_t \mathbf{u} &= -\nabla_{\mathbf{u}} \mathbf{u} - \text{grad } p + \frac{1}{R} \Delta \mathbf{u} \\ 0 &= \text{div } \mathbf{u}\end{aligned}$$

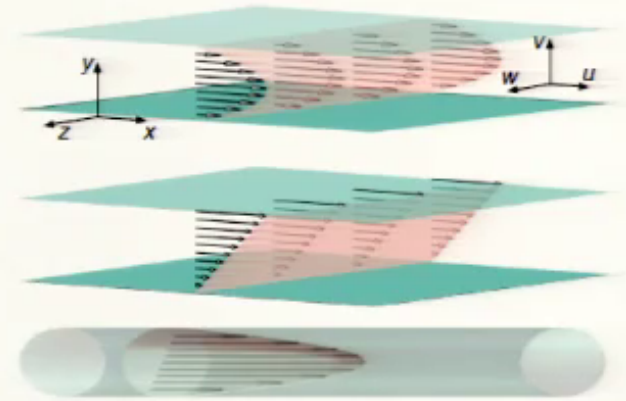


- Hydrodynamic Stability: view NS as a dynamical system
- *laminar flow* $\bar{\mathbf{u}}_R :=$ a stationary solution of the NS equations (an *equilibrium*)

Mathematical Modeling of Transition: Hydrodynamic Stability

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- Hydrodynamic Stability: view NS as a dynamical system
- laminar flow $\bar{\mathbf{u}}_R :=$ a stationary solution of the NS equations (an *equilibrium*)

laminar flow $\bar{\mathbf{u}}_R$ stable \iff i.c. $\mathbf{u}(0) \neq \bar{\mathbf{u}}_R$,
 $\mathbf{u}(t) \xrightarrow{t \rightarrow \infty} \bar{\mathbf{u}}_R$

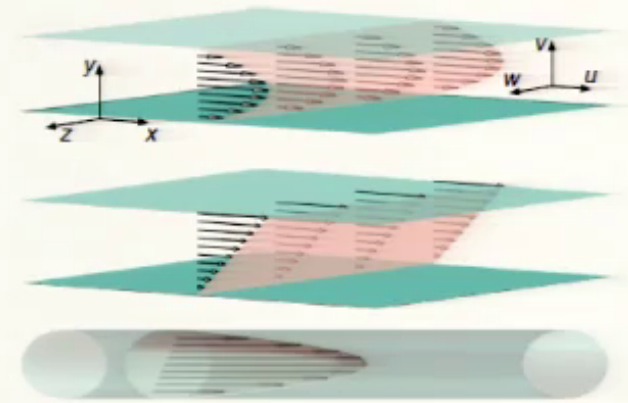
- ▶ typically done with dynamics linearized about $\bar{\mathbf{u}}_R$
- ▶ various methods to track further “non-linear behavior”



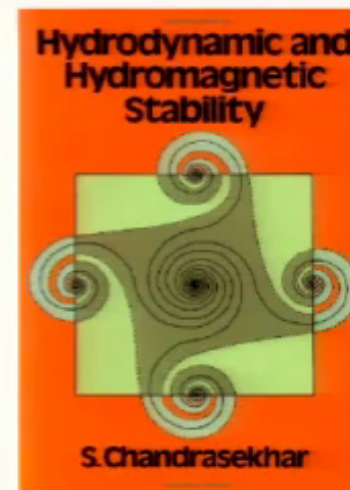
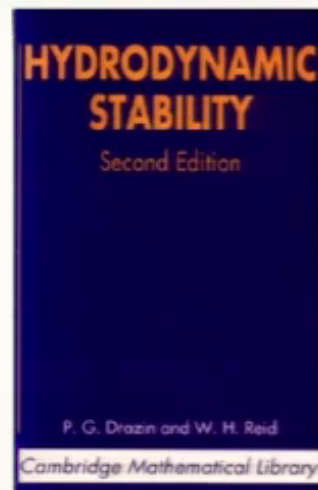
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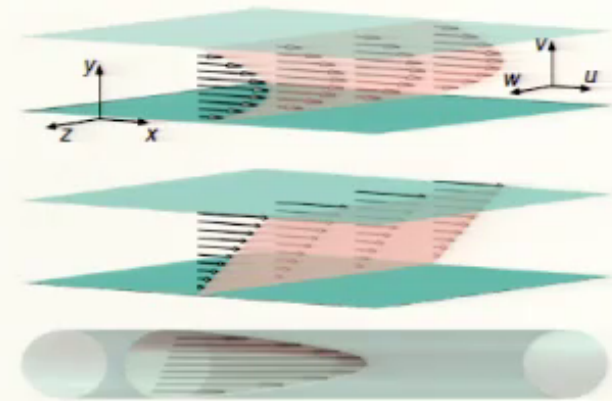
- Hydrodynamic Stability: view NS as a dynamical system
- A very successful (*phenomenologically predictive*) approach for many decades



Mathematical Modeling of Transition: Hydrodynamic Stability

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- Hydrodynamic Stability: view NS as a dynamical system
- however:** problematic for wall-bounded shear flows

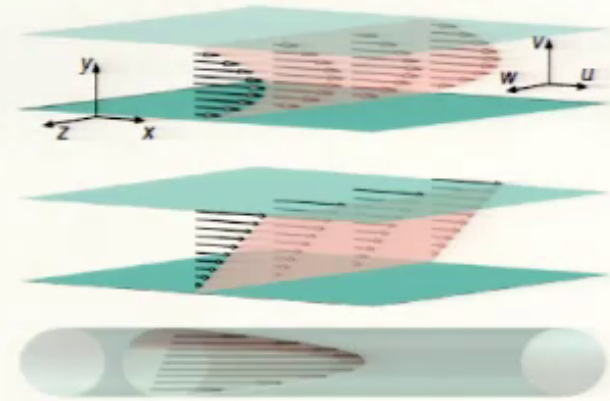
does not predict transition *Reynolds numbers*

Flow type	Classical linear theory R_c	Experimental R_c
Channel Flow	5772	$\approx 1,000-2,000$
Plane Couette	∞	≈ 350
Pipe Flow	∞	$\approx 2,200-100,000$

Mathematical Modeling of Transition: Hydrodynamic Stability

The Navier-Stokes (NS) equations:

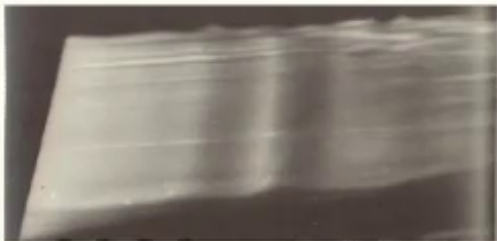
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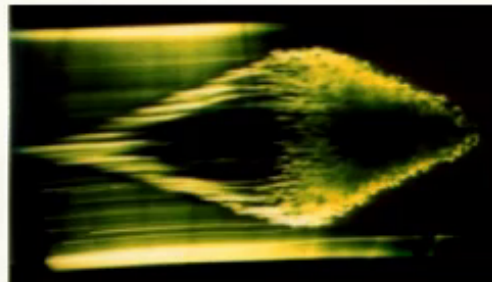
- Hydrodynamic Stability: view NS as a dynamical system
- **however:** problematic for wall-bounded shear flows

does not predict most transition *flow structures*

prediction: T-S waves



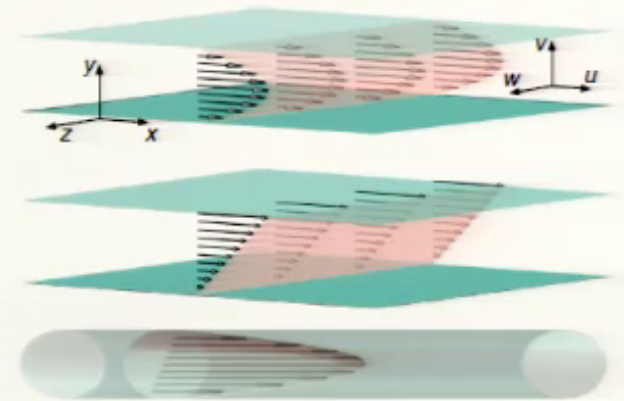
more common: bypass transition



Mathematical Modeling of Transition: Hydrodynamic Stability

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- Hydrodynamic Stability: view NS as a dynamical system
- **however:** problematic for wall-bounded shear flows

- ▶ was widely believed: *this theory fails because it is linear and “nonlinear effects” are important even for infinitesimal i.c.*
- ▶ *however, since 90’s: story is actually more interesting than that*

Nonmodal Stability Theory, Schmid, ARFM '07

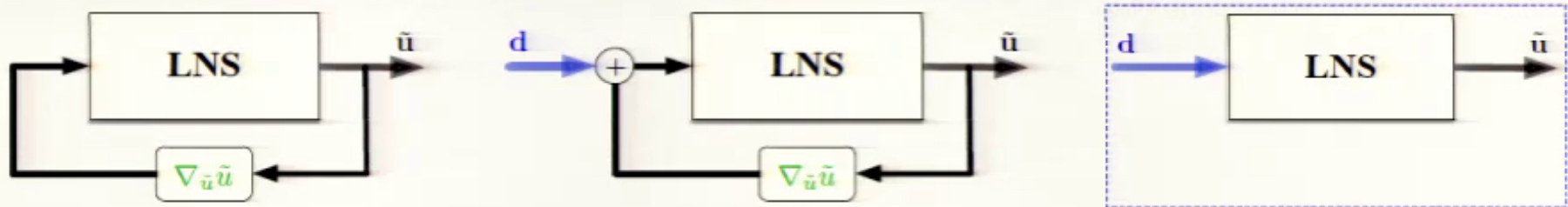
Mathematical Modeling of Transition: Incorporating Uncertainty

- Decompose the fields as

$$\mathbf{u} = \underbrace{\bar{\mathbf{u}}}_{\text{laminar}} + \underbrace{\tilde{\mathbf{u}}}_{\text{fluctuations}}$$

- Add a time-varying *exogenous disturbance* field \mathbf{d} (e.g. random body forces)

$$\begin{aligned} \partial_t \tilde{\mathbf{u}} &= -\nabla_{\bar{\mathbf{u}}} \tilde{\mathbf{u}} - \nabla_{\tilde{\mathbf{u}}} \bar{\mathbf{u}} - \text{grad } \tilde{p} + \frac{1}{R} \Delta \tilde{\mathbf{u}} - \nabla_{\tilde{\mathbf{u}}} \tilde{\mathbf{u}} + \mathbf{d} \\ 0 &= \text{div } \tilde{\mathbf{u}} \end{aligned}$$



Input-Output view of the Linearized NS Equations

Farrell, Ioannou, '93 PoF

BB, Dahleh, '01 PoF

Jovanovic, BB, '05 JFM

Mathematical Modeling of Transition: Incorporating Uncertainty

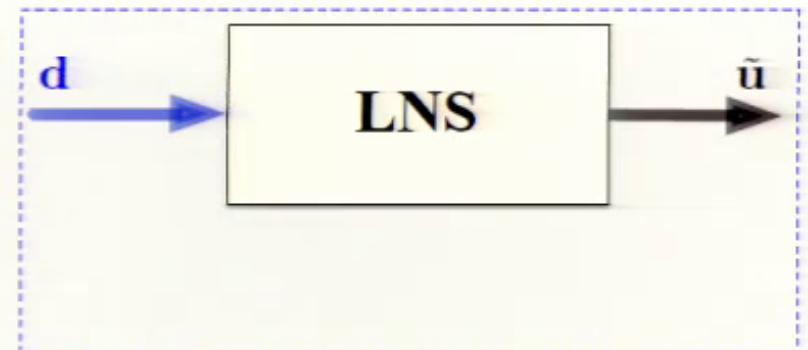
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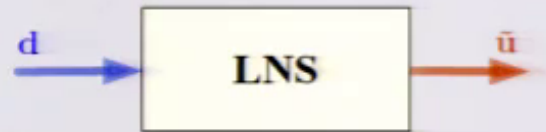
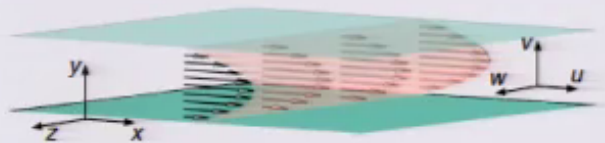
- Neglect the feedback $\nabla_{\tilde{\mathbf{u}}} \tilde{\mathbf{u}}$



Input-Output Analysis of the Linearized NS Equations

$$\partial_t \begin{bmatrix} \Delta \tilde{v} \\ \tilde{\omega} \end{bmatrix} = \begin{bmatrix} U'' \partial_x - U \Delta \partial_x + \frac{1}{R} \Delta^2 & 0 \\ -U' \partial_z & -U \partial_x + \frac{1}{R} \Delta \end{bmatrix} \begin{bmatrix} \tilde{v} \\ \tilde{\omega} \end{bmatrix} + \begin{bmatrix} -\partial_{xy} & \partial_x^2 + \partial_z^2 & -\partial_{zy} \\ \partial_z & 0 & -\partial_x \end{bmatrix} \begin{bmatrix} d_x \\ d_y \\ d_z \end{bmatrix}$$

$$\begin{bmatrix} \tilde{u} \\ \tilde{v} \\ \tilde{w} \end{bmatrix} = (\partial_x^2 + \partial_z^2)^{-1} \begin{bmatrix} \partial_{xy} & -\partial_z \\ \partial_x^2 + \partial_z^2 & 0 \\ \partial_{zy} & \partial_x \end{bmatrix} \begin{bmatrix} \tilde{v} \\ \tilde{\omega} \end{bmatrix}$$



$$\begin{aligned} \partial_t \Psi &= \mathcal{A} \Psi + \mathcal{B} \mathbf{d} \\ \tilde{\mathbf{u}} &= \mathcal{C} \Psi \end{aligned}$$

Surprises:

- Even when \mathcal{A} is stable

the gain $\mathbf{d} \rightarrow \tilde{\mathbf{u}}$ can be very large

(H^2 norm)² scales with R^3)

- Input-output resonances

very different from least-damped modes of \mathcal{A}

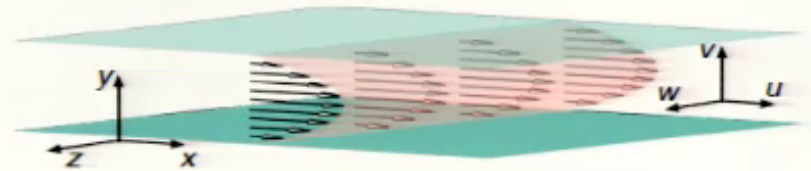


has important implications for control design!

Spatio-temporal Impulse and Frequency Responses

Translation invariance in x & z implies

- *Impulse Response* (Green's Function)



$$\tilde{\mathbf{u}}(t, x, y, z) = \int G(t - \tau, x - \xi, \mathbf{y}, \mathbf{y}', z - \zeta) \mathbf{d}(\tau, \xi, \mathbf{y}', \zeta) d\tau d\xi dy' d\zeta$$

$$\tilde{\mathbf{u}}(t, x, \cdot, z) = \int \mathcal{G}(t - \tau, x - \xi, z - \zeta) \mathbf{d}(\tau, \xi, \cdot, \zeta) d\tau d\xi d\zeta$$

$\mathcal{G}(t, x, z)$: Operator-valued impulse response

- *Frequency Response*

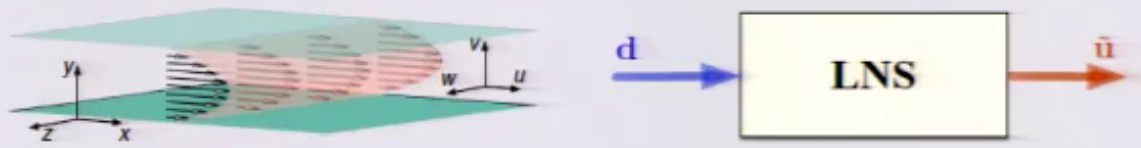
$$\tilde{\mathbf{u}}(\omega, k_x, k_z) = \mathcal{G}(\omega, k_x, k_z) \mathbf{d}(\omega, k_x, k_z)$$

$\mathcal{G}(\omega, k_x, k_z)$: Operator-valued frequency response (Packs lots of information!)

- *Spectrum of \mathcal{A} :*

$$\sigma(\mathcal{A}) = \overline{\bigcup_{k_x, k_z} \sigma(\hat{\mathcal{A}}(k_x, k_z))}$$

Modal vs. Input-Output Analysis

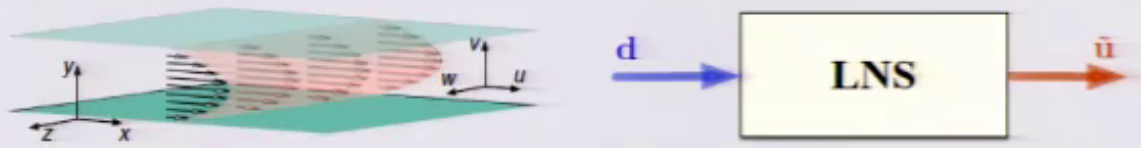


- $\partial_t \Psi = \mathcal{A} \Psi + \mathcal{B} d$
- $\tilde{u} = \mathcal{C} \Psi$
- IR: $\mathcal{G}(t, x, z)$
- FR: $\mathcal{G}(\omega, k_x, k_z)$

Modal Analysis: Look for unstable eigs of $\mathcal{A} \left(\bigcup_{k_x, k_z} \sigma(\hat{\mathcal{A}}(k_x, k_z)) \right)$

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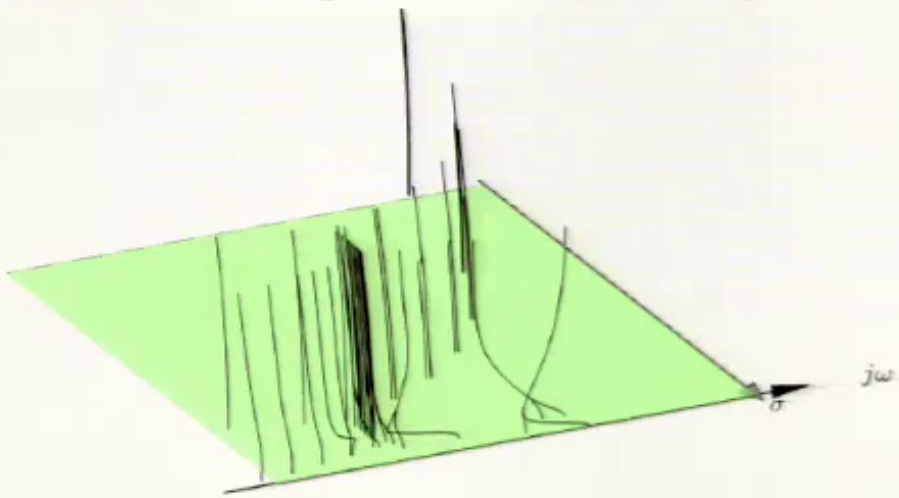
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Modal Analysis: Look for unstable eigs of \mathcal{A}

$$\left(\bigcup_{k_x, k_z} \sigma \left(\hat{\mathcal{A}}(k_x, k_z) \right) \right)$$

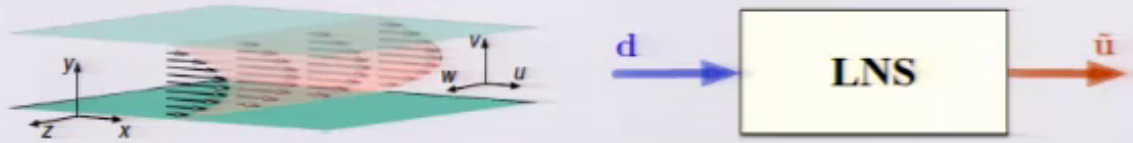
- Channel Flow @ $R = 2000, k_x = 1,$

($k_z =$ vertical dimension):



top view

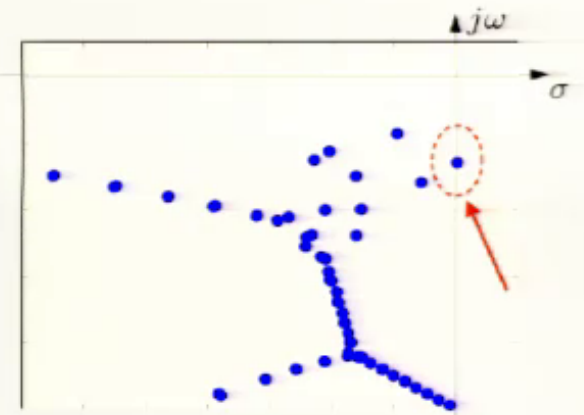
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Modal Analysis: Look for unstable eigs of $\mathcal{A} \left(U_{k_x, k_z} \sigma \left(\hat{\mathcal{A}}(k_x, k_z) \right) \right)$

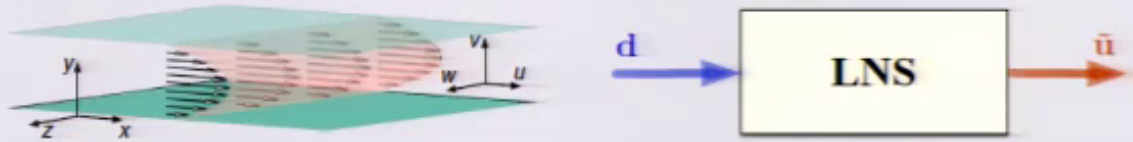
- Channel Flow @ $R = 6000, k_x = 1, k_z = 0$:



- Flow structure of corresponding eigenfunction:
Tollmein-Schlichting (TS) waves

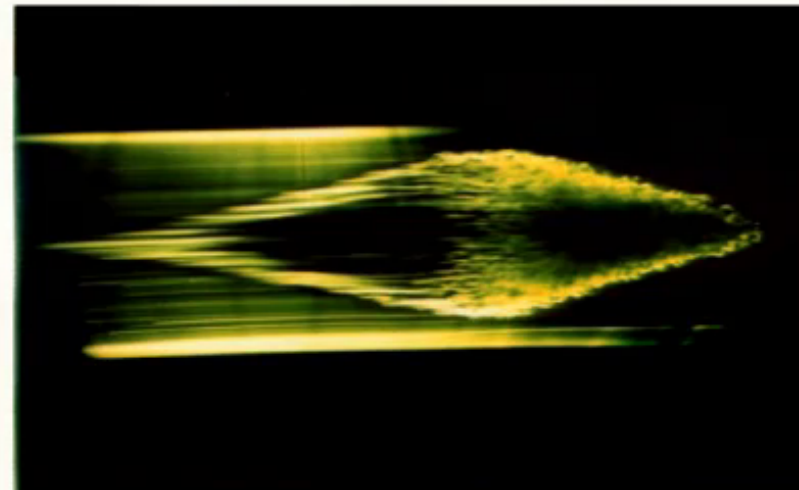
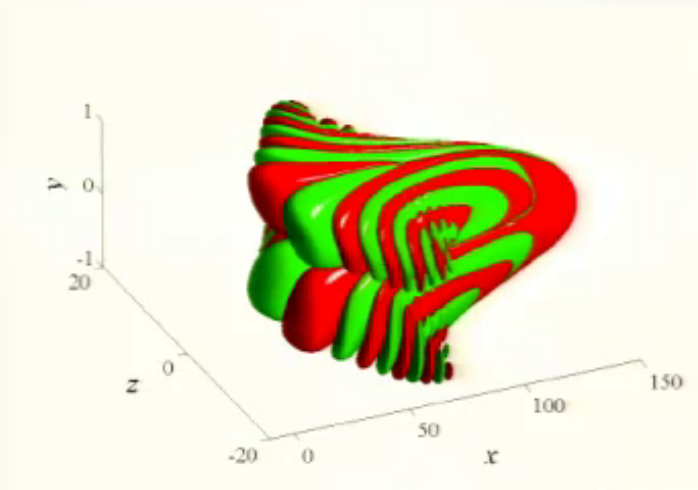


Modal vs. Input-Output Analysis



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Impulse Response Analysis: Channel Flow @ $R = 2000$

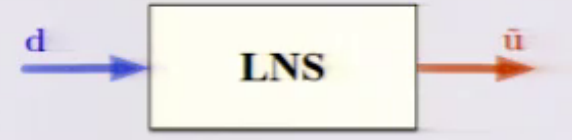
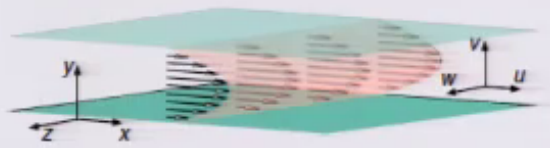


cf. "turbulent spots"

Jovanovic, BB, '01 ACC,

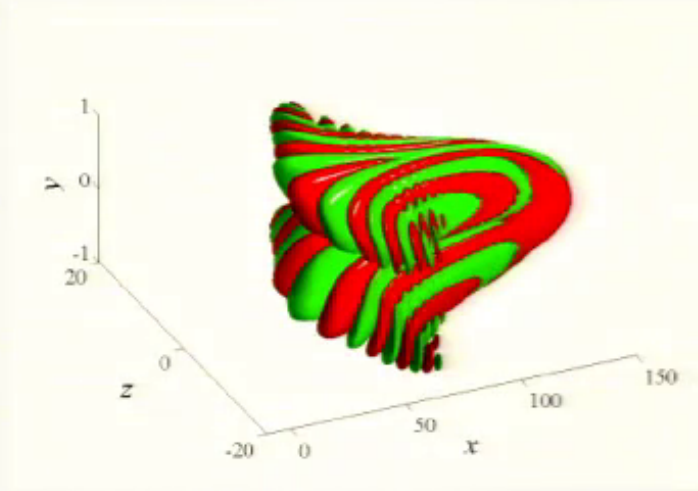
more movies and pics at http://engineering.ucsb.edu/~bamieh/pics/impulse_page.html

Modal vs. Input-Output Analysis

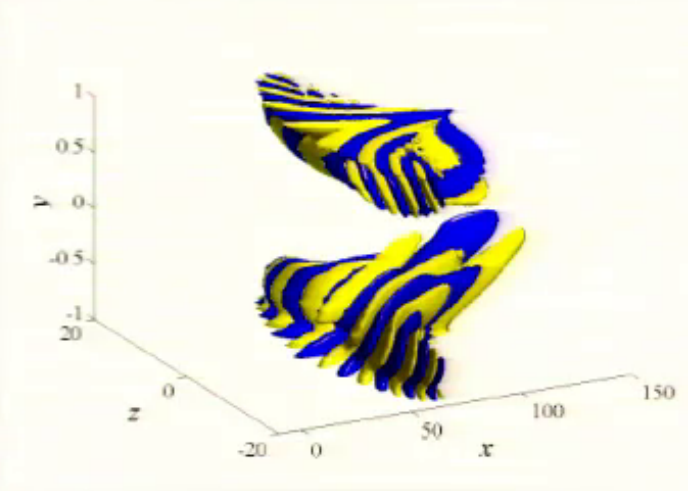


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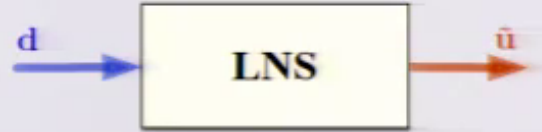
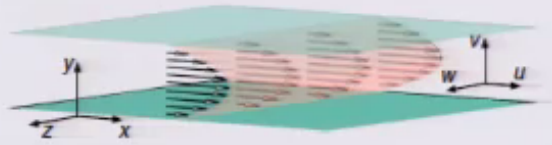


streamwise velocity



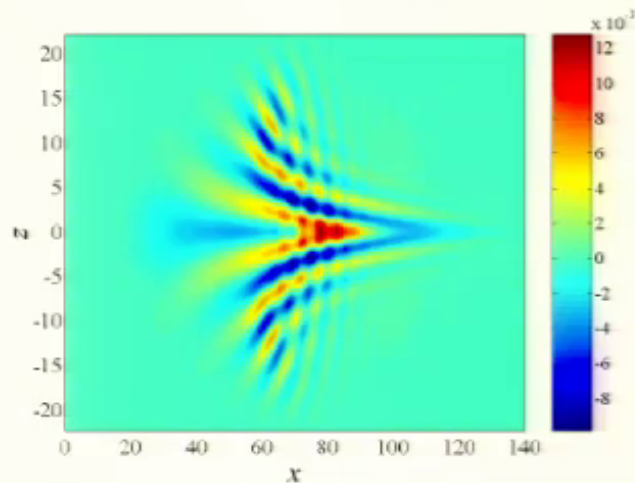
streamwise vorticity

Modal vs. Input-Output Analysis

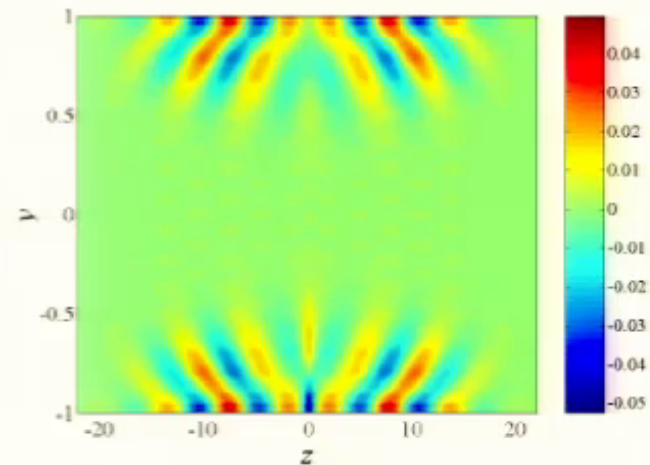


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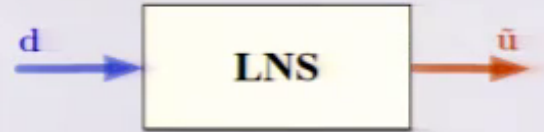
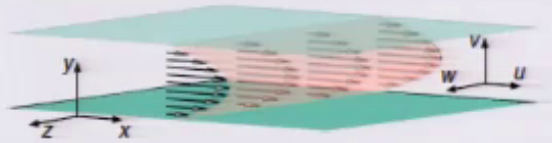


u in a horizontal plane



u in a vertical plane

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- *Frequency Response*

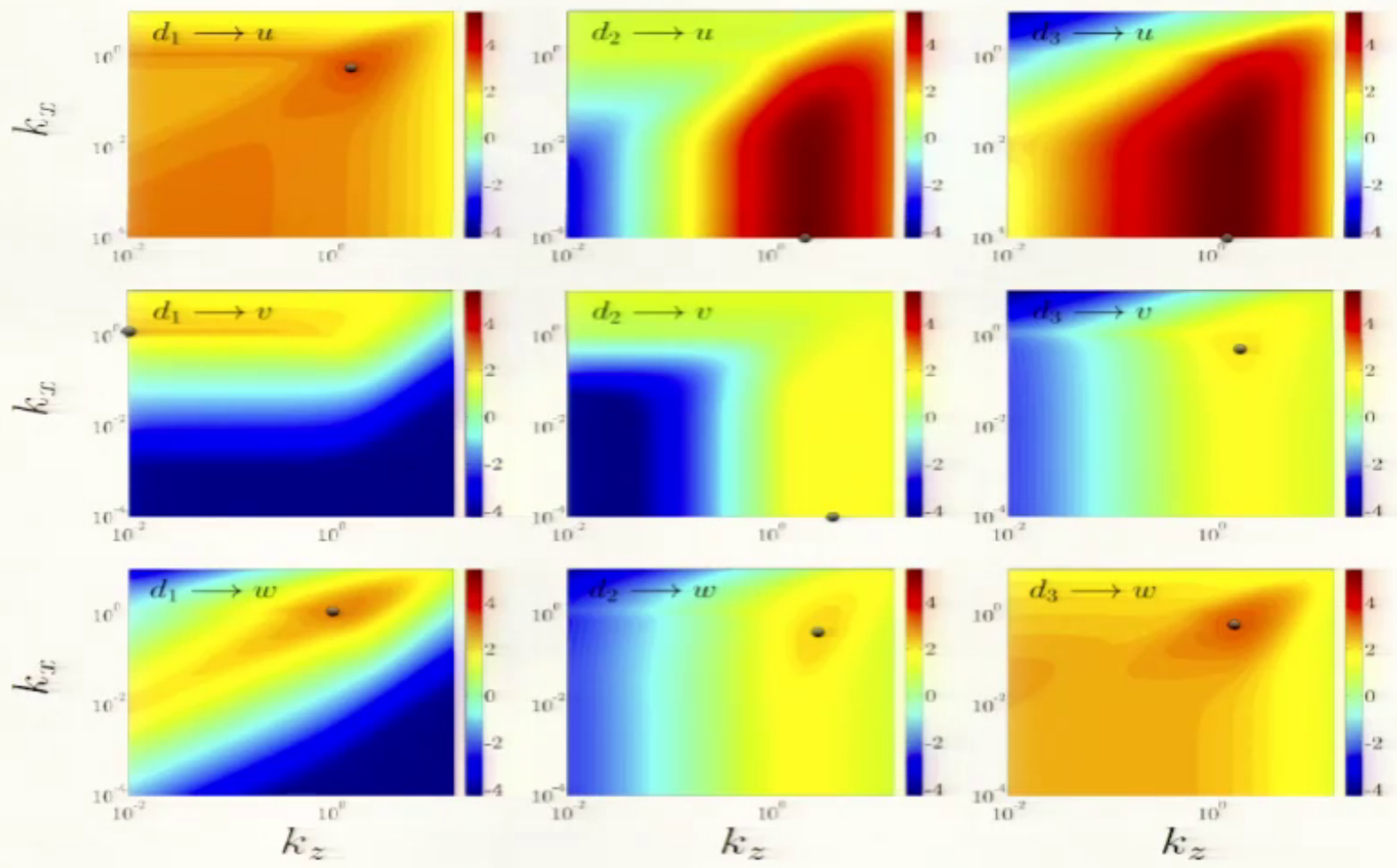
$$\tilde{\mathbf{u}}(\omega, k_x, k_z) = \mathcal{G}(\omega, k_x, k_z) \mathbf{d}(\omega, k_x, k_z)$$

$\mathcal{G}(\omega, k_x, k_z)$: **Operator-valued frequency response** (Packs lots of information!)

Spatio-temporal Frequency Response

$\mathcal{G}(\omega, k_x, k_z)$ is a *LARGE* object! (very “data rich”!)

one visualization: $\sup_{\omega} \sigma_{\max} \left(\mathcal{G}(\omega, k_x, k_z) \right)$ (max over temporal frequency & wall-normal dimension)



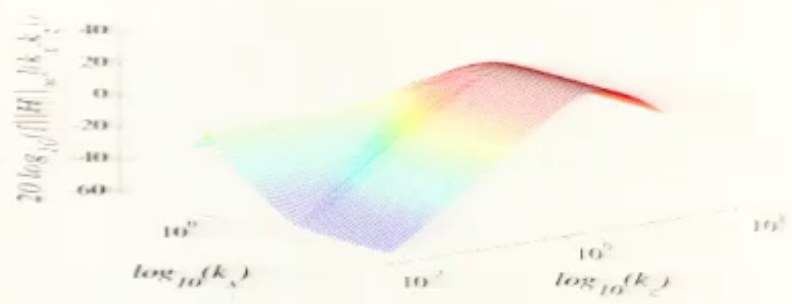
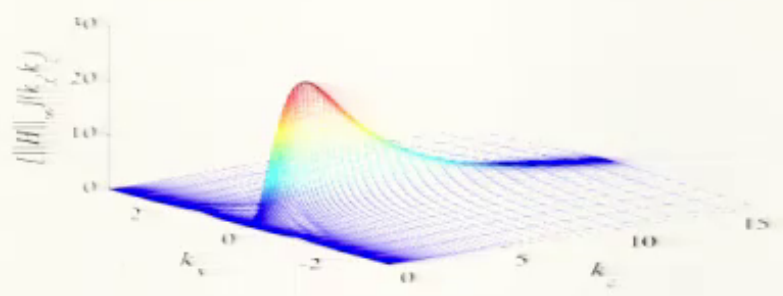
Jovanovic, BB, '05 JFM

Spatio-temporal Frequency Response

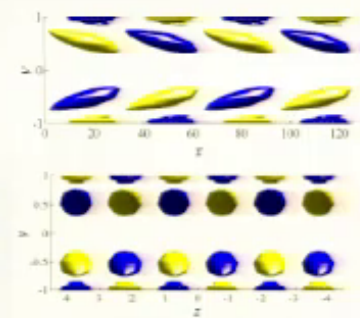
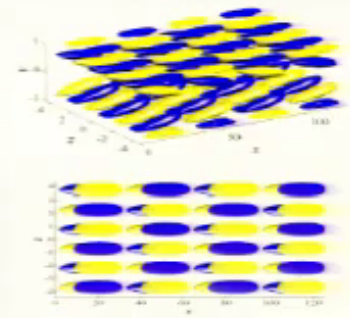
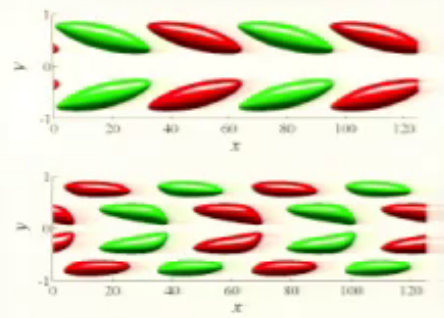
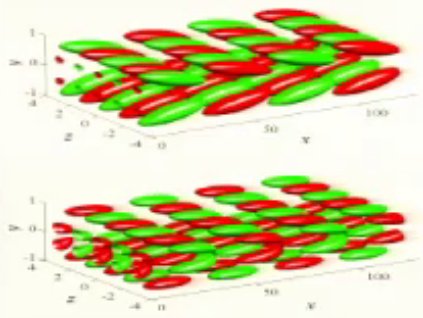
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(max over *temporal frequency* & *wall-normal dimension*)



flow structure corresponding to peak:



streamwise velocity isosurfaces

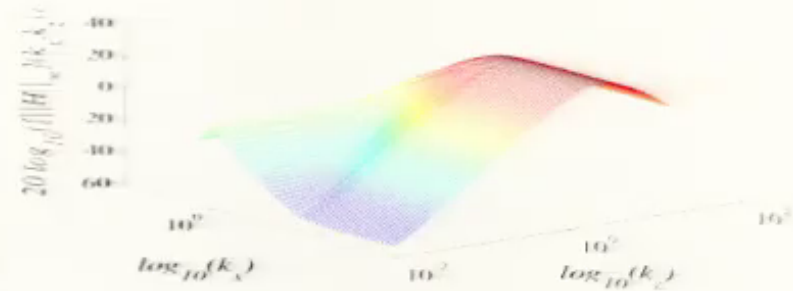
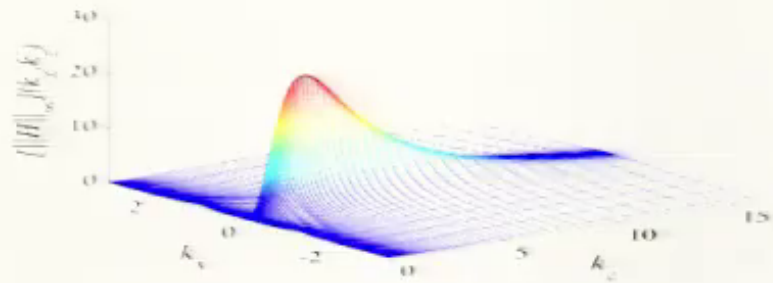
streamwise vorticity isosurfaces

Spatio-temporal Frequency Response

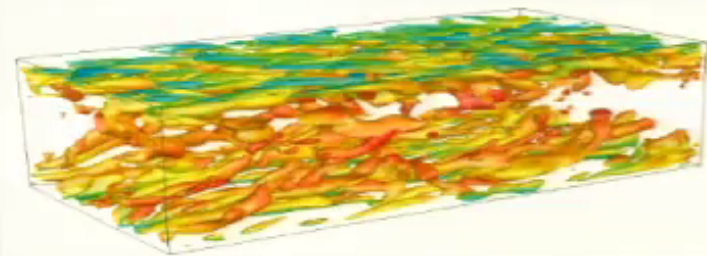
$\mathcal{G}(\omega, k_x, k_z)$ is a *LARGE* object! (very “data rich”!)

one visualization: $\sup_{\omega} \sigma_{\max} \left(\mathcal{G}(\omega, k_x, k_z) \right)$

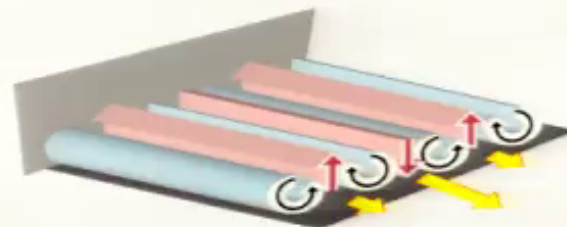
(max over *temporal frequency*
& *wall-normal dimension*)



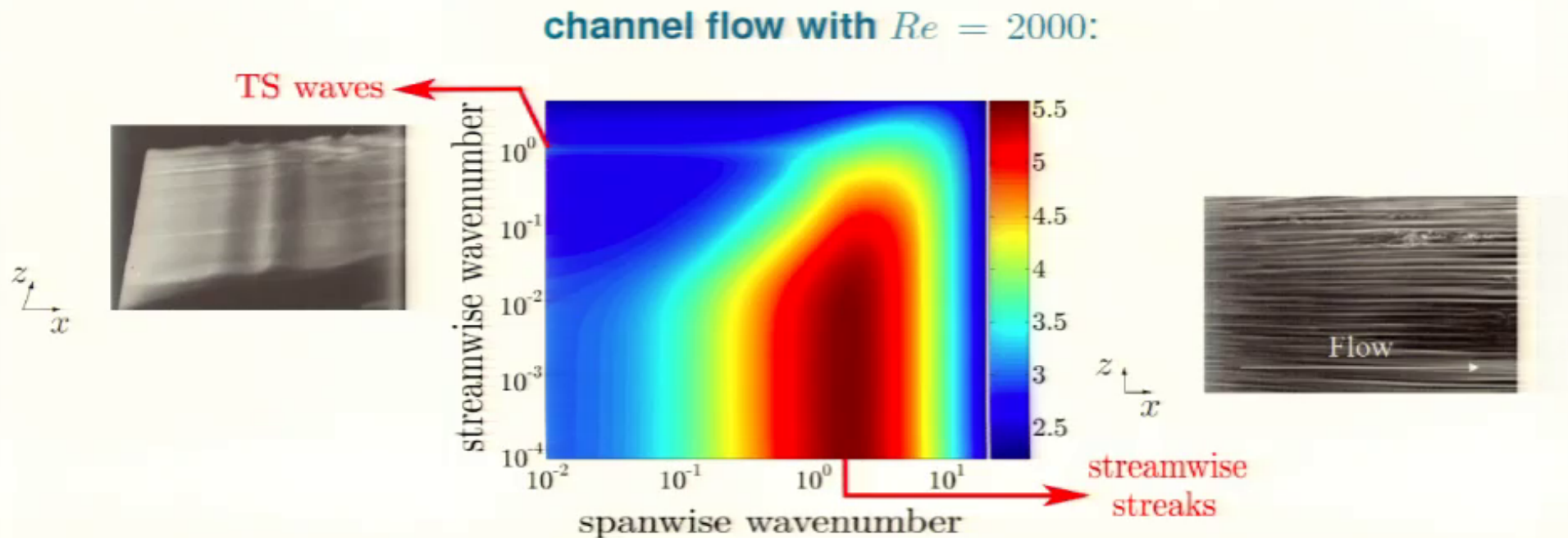
flow structure corresponding to peak:



similar to observed near-wall structures



Underdamped Modes vs. Large IO Resonances



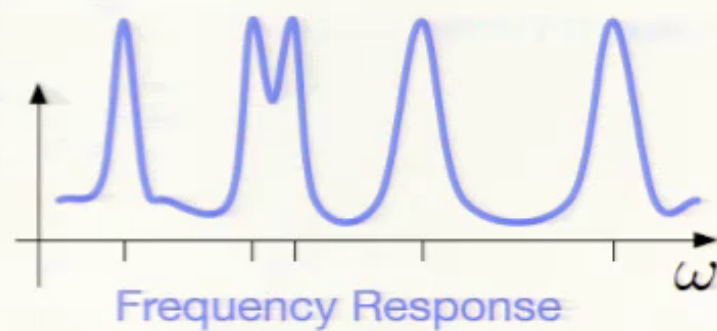
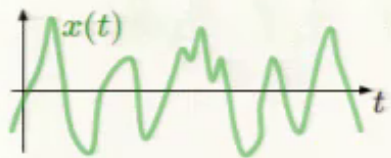
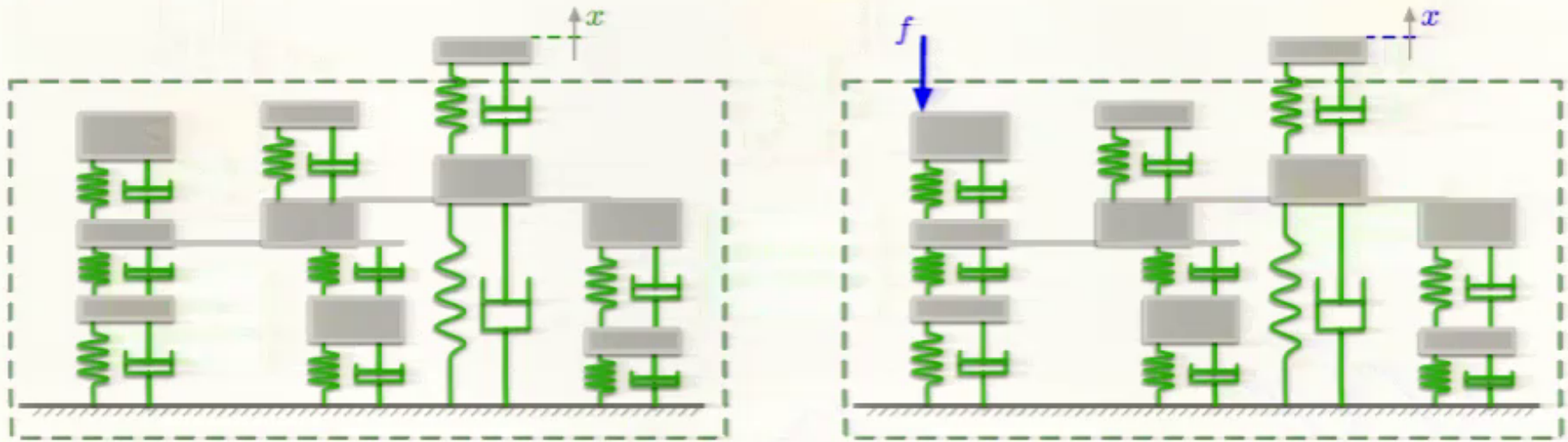
- highly underdamped TS modes barely register in the IO response
- large peaks of IO response *do not* correspond to any particular pure mode

Internal Modes vs. External Resonances

A Detour

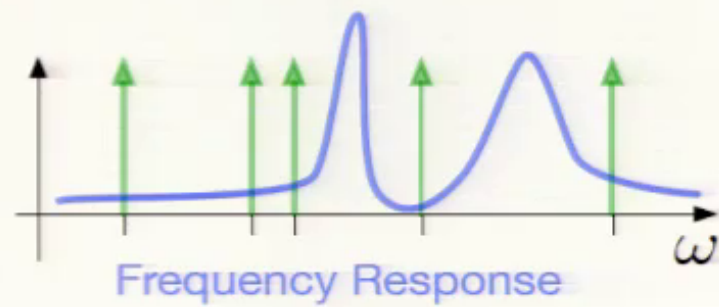
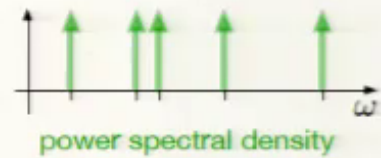
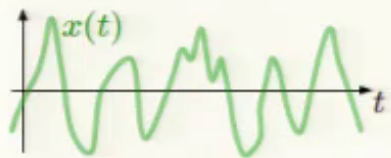
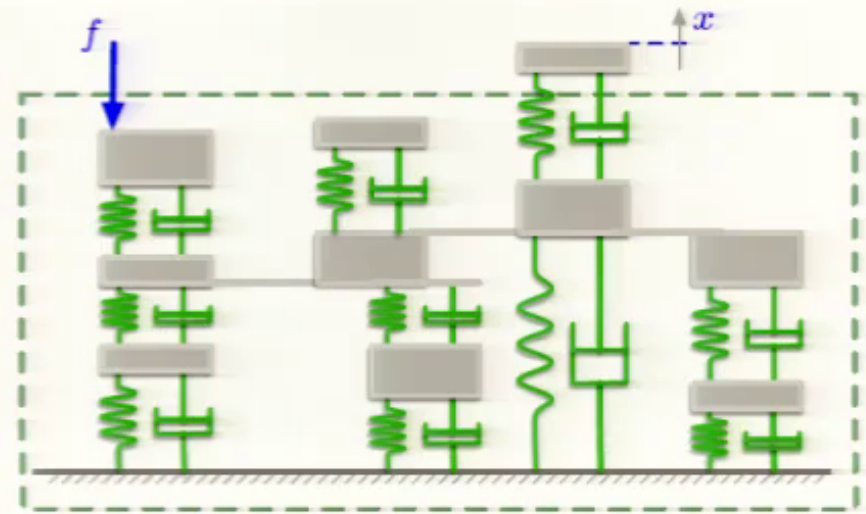
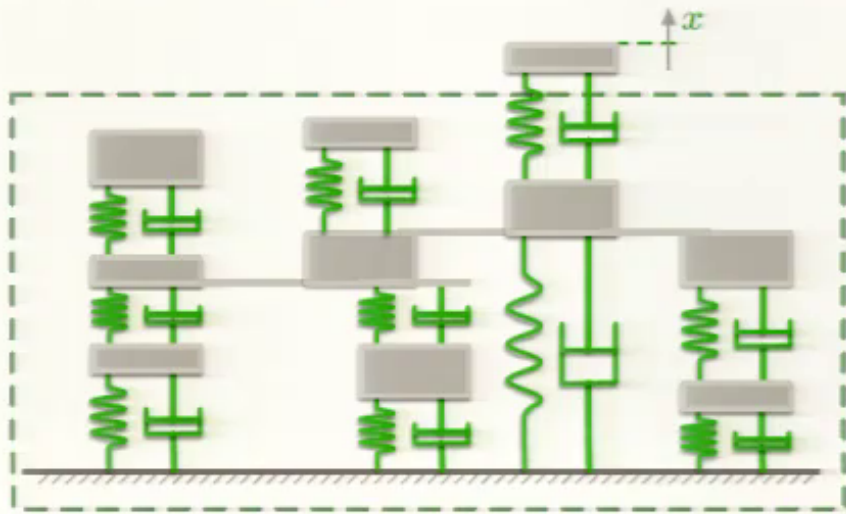
Internal Modes vs. External Resonances

Does this correspondence hold for large-scale systems?



Internal Modes vs. External Resonances

Does this correspondence hold for large-scale systems?

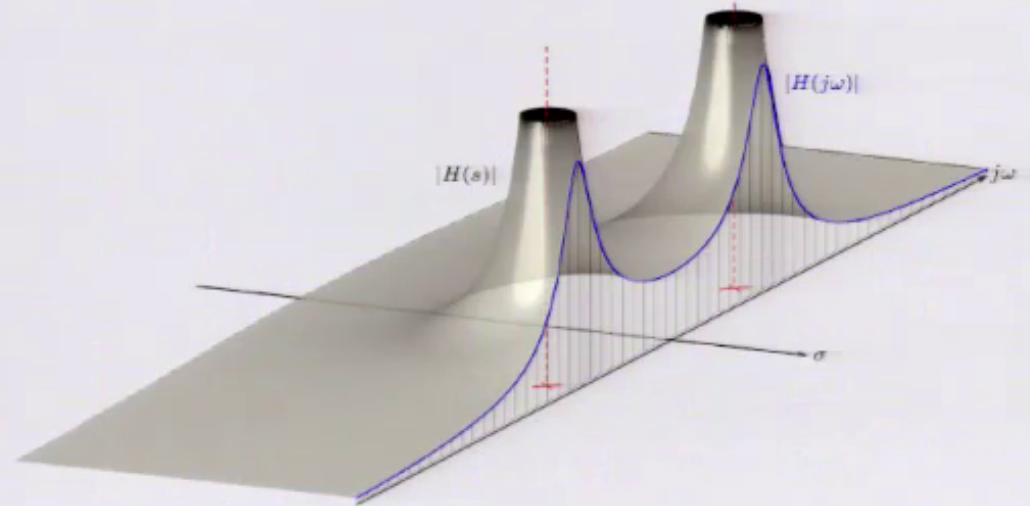


However: this may not hold in general
even in linear systems!

Modal vs. Input-Output Response

Typically: underdamped poles \longleftrightarrow frequency response peaks

cf. The “rubber sheet analogy”:



ODE (state space model)

$$\dot{\psi}(t) = A \psi(t) + B d(t)$$

$$\tilde{\mathbf{u}}(t) = C \psi(t)$$

Transfer Function

$$H(s) = C (sI - A)^{-1} B$$

- $\text{eigs}(A) = \text{poles}(H(s))$

Modal vs. Input-Output Response

However: Pole Locations \leftrightarrow Frequency Response Peaks

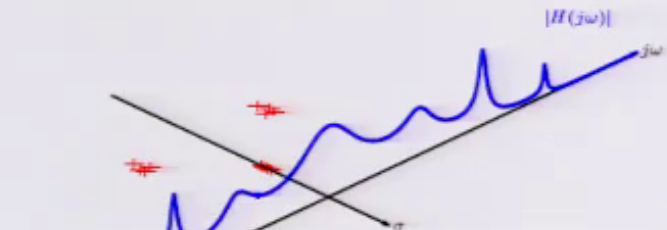
Theorem: Given any desired *pole locations*

$$z_1, \dots, z_n \in \mathbb{C}_- \text{ (LHP),}$$

and any *stable frequency response* $H(j\omega)$, arbitrarily close approximation is achievable with

$$\left\| H(s) - \left(\sum_{i=1}^{N_1} \frac{\alpha_{1,i}}{(s - z_1)^i} + \dots + \sum_{i=1}^{N_n} \frac{\alpha_{n,i}}{(s - z_n)^i} \right) \right\|_{\mathcal{H}^2} \leq \epsilon$$

by choosing any of the N_k 's large enough



Remarks:

- No necessary relation between *pole locations* and *system resonances*
- ($\epsilon \rightarrow 0 \Rightarrow N_k \rightarrow \infty$), i.e. this is a *large-scale systems* phenomenon
- **Large-scale systems:** IO behavior not always predictable from modal behavior

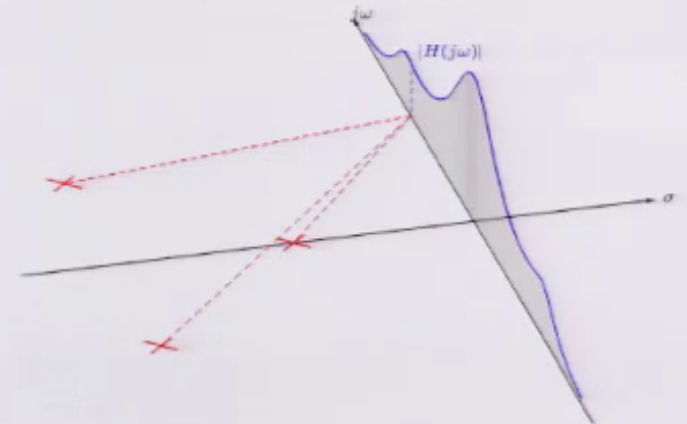
Modal vs. Input-Output Response

However: Pole Locations \leftrightarrow Frequency Response Peaks

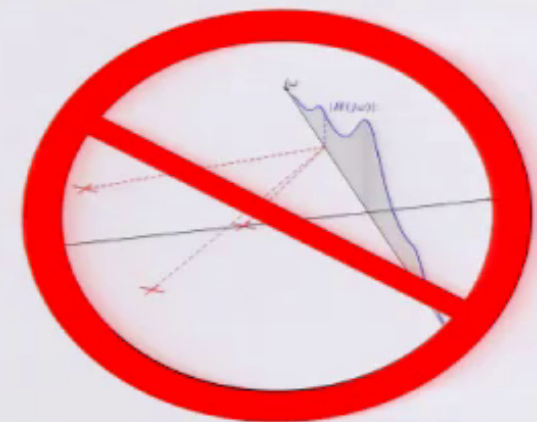
MIMO case: $H(s) = (sI - A)^{-1}$

- If A is *normal* (has orthogonal eigenvectors), then

$$\sigma_{\max} \left((j\omega I - A)^{-1} \right) = \frac{1}{\text{distance}(j\omega, \text{nearest pole})}$$

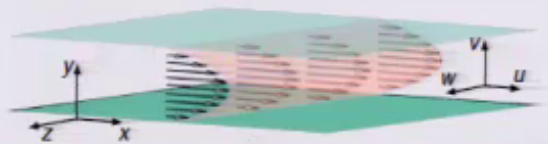


- If A is *non-normal*: no clear relation between singular value plot \leftrightarrow eigs(A)

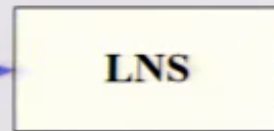


Implications for Turbulence Modeling

For large-scale systems: IO behavior not predictable from modal behavior



\mathbf{d}

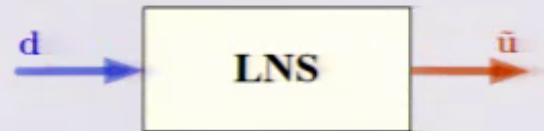
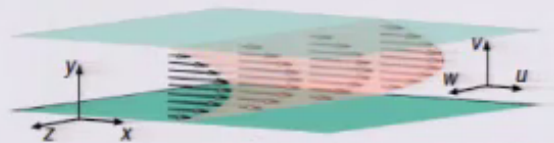


$\tilde{\mathbf{u}}$

- $\partial_t \Psi = \mathcal{A} \Psi + \mathcal{B} \mathbf{d}$
- $\tilde{\mathbf{u}} = \mathcal{C} \Psi$
- IR: $\mathcal{G}(t, x, z)$
- FR: $\mathcal{G}(\omega, k_x, k_z)$

Implications for Turbulence Modeling

For large-scale systems: IO behavior not predictable from modal behavior



- $\partial_t \Psi = \mathcal{A} \Psi + \mathcal{B} \mathbf{d}$
- $\tilde{\mathbf{u}} = \mathcal{C} \Psi$
- IR: $\mathcal{G}(t, x, z)$
- FR: $\mathcal{G}(\omega, k_x, k_z)$

- “modal behavior”: Stability due to *initial condition uncertainty*
- “IO behavior”: behavior in the presence of *ambient uncertainty*
 - ▶ forcing terms from wall roughness and/or vibrations
 - ▶ Free-stream disturbances in boundary layers
 - ▶ Thermal (Langevin) forces
 - ▶ uncertain dynamics