

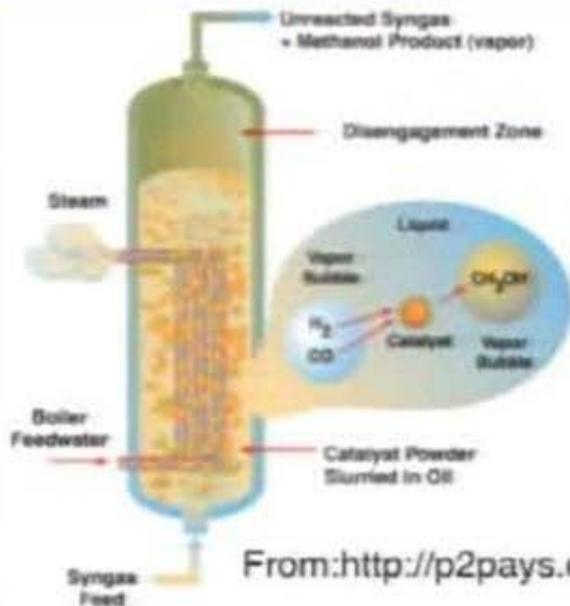


Direct Numerical Simulations of Multiphase Flows. Now What?

Grétar Tryggvason,
University of Notre Dame

SIAM Conference on Computational Science and Engineering (CSE15)
Salt Lake City, March 14-18

Work supported by NSF & DOE (CASL & PSAAP)



Liquid Phase Methanol (LPMEOH™) Process

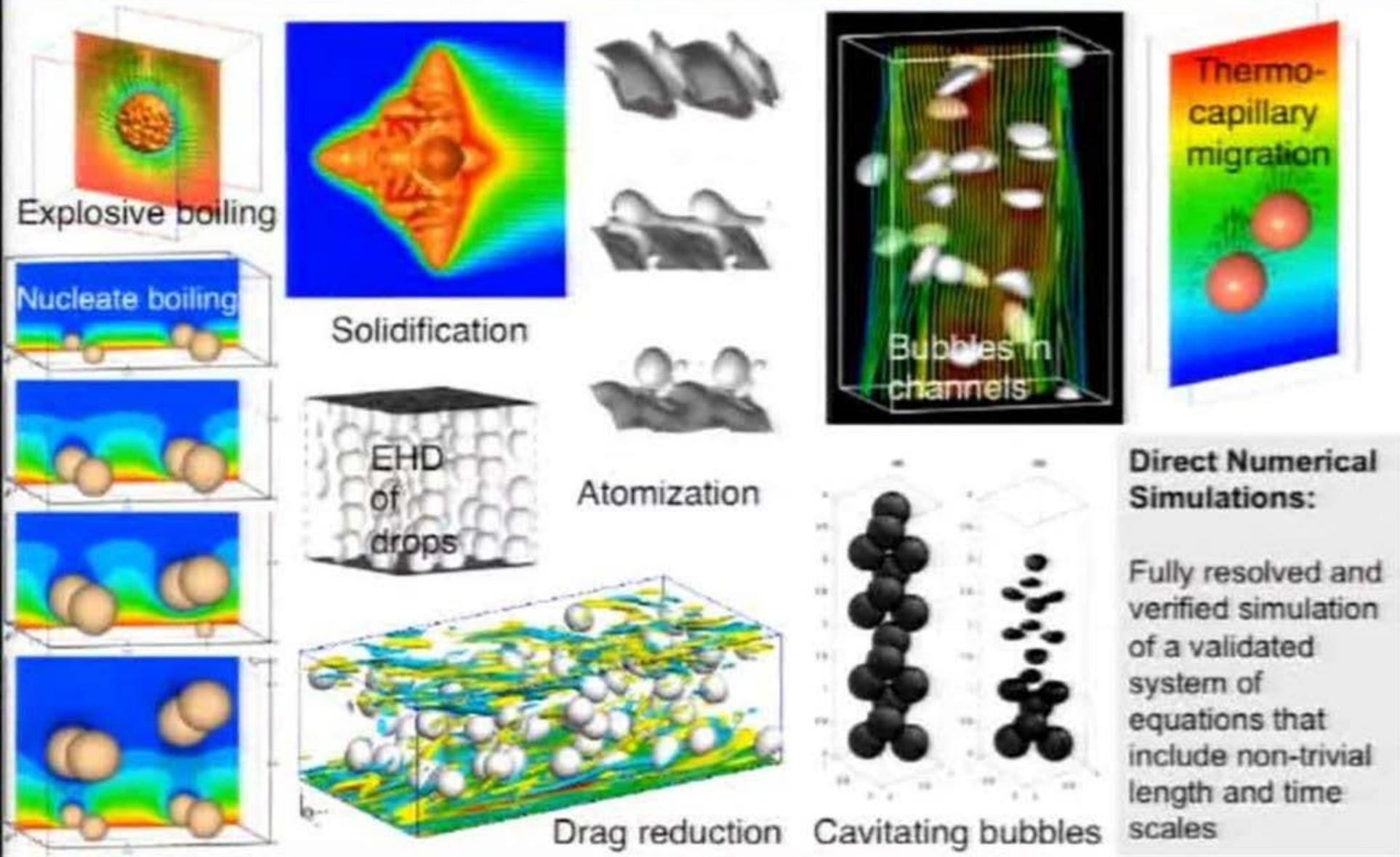
From: <http://p2pays.org/ref/16/15865.pdf>

Examples of Applications of Bubble Columns

Process	Methods and/or Reactants
Acetone	Oxidation of cumene
Acetic acid	Oxidation of acetaldehyde
	Oxidation of sec-butanol
	Carbonylation of methanol
Acetic anhydride	Oxidation of acetaldehyde
Acetaldehyde	Partial oxidation of ethylene
Acetophenone	Oxidation of ethylbenzene
Barium chloride	Barium sulfide and chlorine
Benzic acid	Oxidation of toluene
Bleaching powder	Aqueous calcium oxide and chlorine
Bromine	Aqueous sodium bromide and chlorine
Butene	Absorption in aqueous solutions of sulfuric acid
Carbon Dioxide	Absorption in ammoniated brine
Carbon tetrachloride	Carbon disulfide and chlorine
Copper oxychloride	Oxidation of cuprous chloride
Cumene	Oxidation of phenol
Cupric chloride	Copper and cupric acid or hydrochloric acid
Dichlorination	Oxychlorination of ethylene
Ethyl benzene	Benzene and ethylene
Hexachlorobenzene	Benzene and chlorine
Hydrogen peroxide	Oxidation of hydroquinone
Isobutylene	Absorption in aqueous solutions of sulfuric acid
Phthalic acid	Oxidation of xylene
Phenol	Oxidation of cumene
Potassium bicarbonate	Aqueous potassium carbonate
Sodium bicarbonate	Aqueous sodium carbonate
Sodium metabisulfides	Carbon dioxide, aqueous sodium carbonate, and sulfur dioxide
Thiourea disulfides	Dithiocarbamates, chlorine, and air
Vinyl acetate	Oxidation of ethylene in acetic acid solutions
Water	Wet oxidation of waste water



After: S. Furusaki, L.-S. Fan, J. Garside.
The Expanding World of Chemical Engineering (2nd ed), Taylor & Francis
2001



Direct Numerical Simulations:

Fully resolved and verified simulation of a validated system of equations that include non-trivial length and time scales

CFD of Multiphase Flows—one slide history

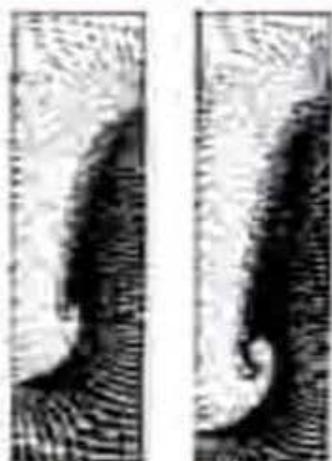
BC: Birkhoff and boundary integral methods for the Rayleigh-Taylor Instability

65' Harlow and colleagues at Los Alamos: The MAC method

75' Boundary integral methods for Stokes flow and potential flow

85' Alternative approaches (body fitted, unstructured, etc.)

95' Beginning of DNS of multiphase flow. Return of the "one-fluid" approach and development of other techniques



From G. Zang et al.



E.H. Merzban, I.P. Skerjanc
J.E. Wozniak, Liquid mixing
by computer, Science 199
1995, 1082-1092.

Simulation of cavitation using 13 trillion grid points to resolve the collapse of 15,000 bubbles, using 1.6 million cores of Sequoia, reaching 55% of its nominal peak performance

2013 ACM Gordon Bell Award



11 PFLOP/s Simulations of Cloud Cavitation Collapse

SC '13 Proceedings of the International Conference on High Performance Computing, Networking, Storage and Analysis, Article No. 3

Diego Rossinelli, Babak Hejazialhosseini, Panagiotis Hadjidoukas, Costas Bekas, Alessandro Curioni, Adam Bertsch, Scott Futral, Steffen J. Schmidt, Nikolaus A. Adams and Petros Koumoutsakos

In many cases, incompressible isothermal problems with immiscible flows describe the flows of interest. In this case, the equations are

Conservation of Momentum

$$\rho \frac{\partial \mathbf{u}}{\partial t} + \rho \nabla \cdot \mathbf{u} \mathbf{u} = -\nabla p + \mathbf{f} + \nabla \cdot \mu (\nabla \mathbf{u} + \nabla^T \mathbf{u}) + \int_F \sigma \kappa \mathbf{n} \delta(\mathbf{x} - \mathbf{x}_f) da$$

Singular interface term

Conservation of Mass

$$\nabla \cdot \mathbf{u} = 0 \quad \text{Incompressible flow}$$

Equation of State:

$$\frac{D\rho}{Dt} = 0; \quad \frac{D\mu}{Dt} = 0 \quad \text{Constant properties}$$



Oscillating drop: pressure

The conservation equations can be solved by relatively standard methods. In most cases regular structured fixed grid are used

The main challenge is how to advect the marker function

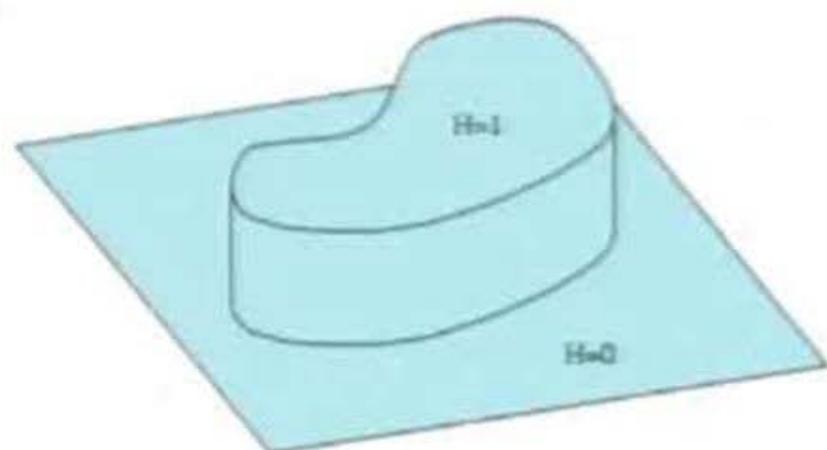
The conservation equations for mass, momentum and energy can be solved by standard techniques. The main challenge is the advection of the marker function

$$H = \begin{cases} 1 & \text{in fluid 1} \\ 0 & \text{Otherwise} \end{cases}$$

$$\frac{\partial H}{\partial t} + \mathbf{u} \cdot \nabla H = 0$$

Other challenges include

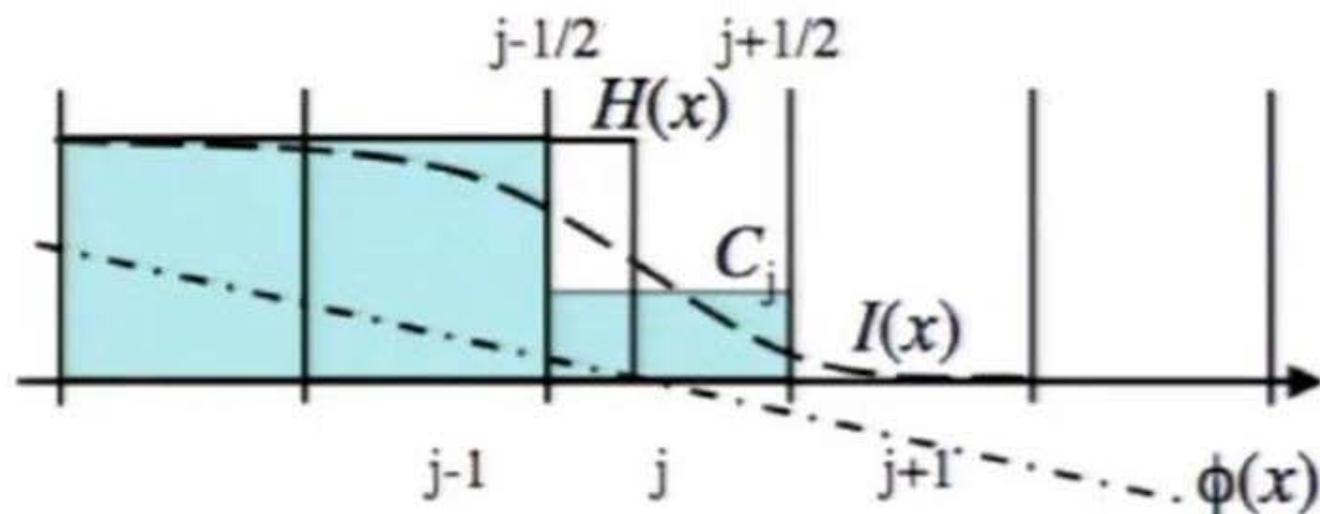
- computing surface tension,
- handling large density and viscosity differences, and
- topology changes.



Updating H —in spite of its apparent simplicity—is one of the hard problems in CFD!



The sharp marker function H can be approximated in several different ways for computational purposes.



Marker Points

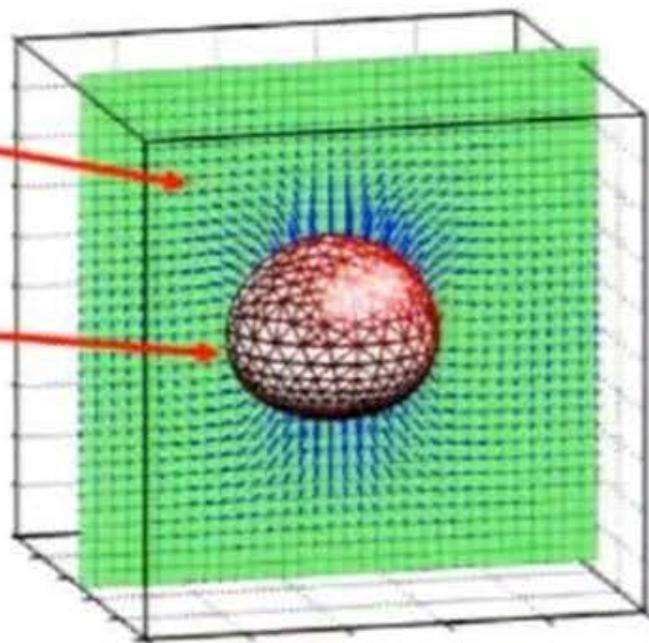
Marker Function

Front Tracking

Fixed grid used for the solution of the Navier-Stokes equations

Tracked front to advect the fluid interface and find surface tension

The method has been used to simulate many problems and extensively tested and validated



Direct Numerical Simulations
of Gas-Liquid Multiphase Flows

Gertur Tryggvason,
Rudolf Scardovelli and Stéphane Zaleski



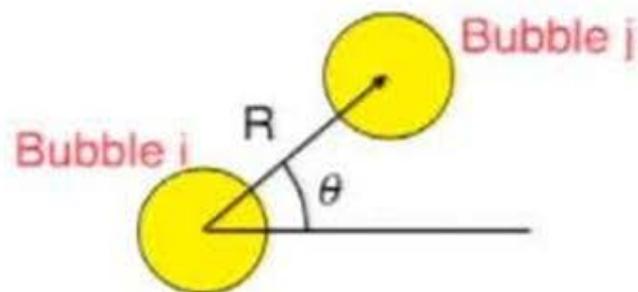
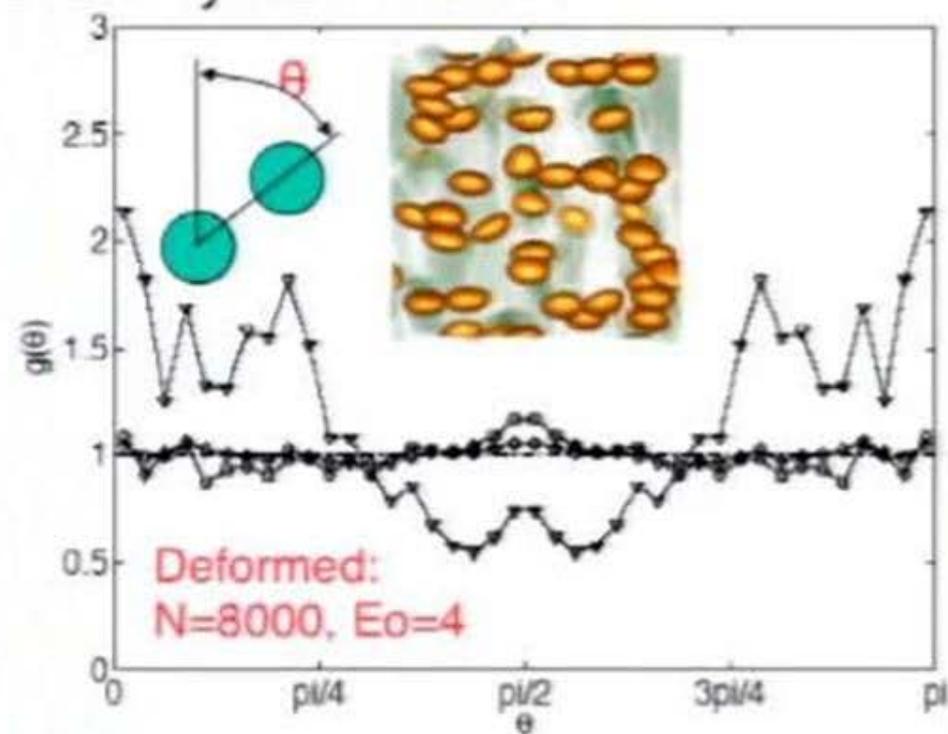
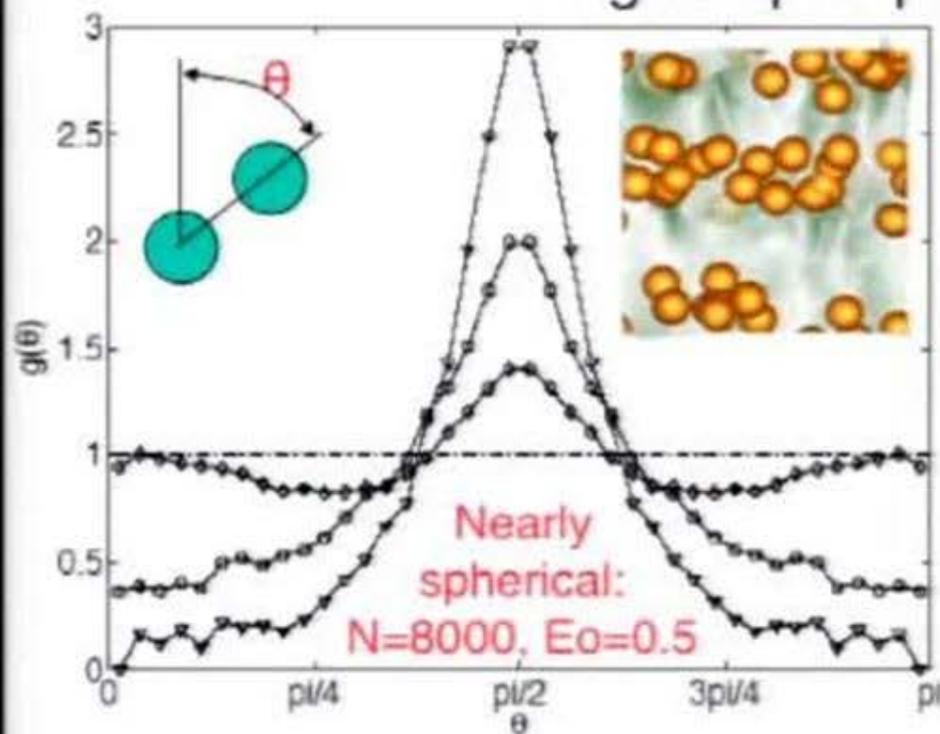
$$\rho \frac{\partial \mathbf{u}}{\partial t} + \rho \nabla \cdot \mathbf{u} \mathbf{u} = -\nabla p + \mathbf{f} + \nabla \cdot \mu (\nabla \mathbf{u} + \nabla^T \mathbf{u}) + \int_F \sigma \kappa \mathbf{n} \delta(\mathbf{x} - \mathbf{x}_f) da$$

Singular interface term

$$\nabla \cdot \mathbf{u} = 0 \quad \frac{D\rho}{Dt} = 0; \quad \frac{D\mu}{Dt} = 0$$

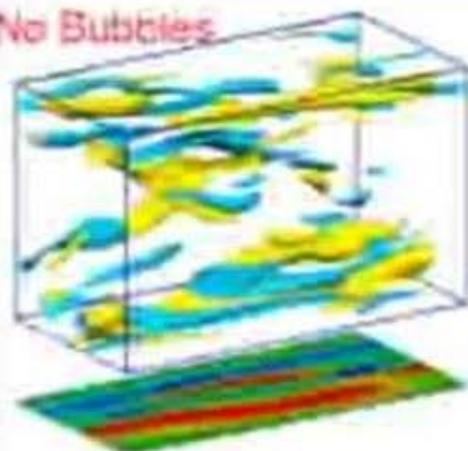
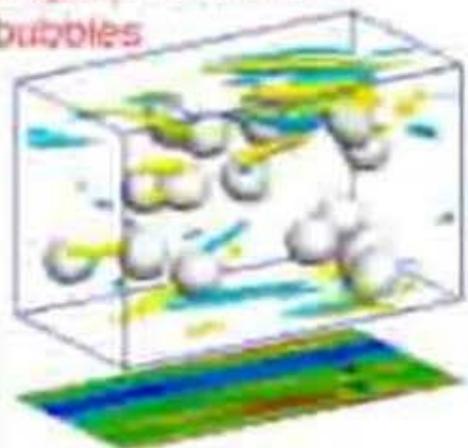


Angular pair probability distribution

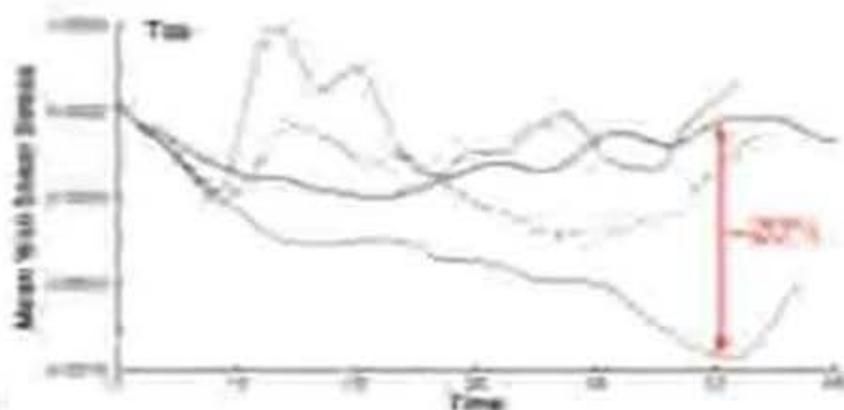


$$G(\vec{r}) = \frac{V}{N(N-1)} \left\langle \sum_i \sum_{j \neq i} \delta(\vec{r} - \vec{r}_{ij}) \right\rangle$$

No Bubbles


 Slightly deformable
 bubbles


DNS of bubbles injected near the wall in a turbulent channel flow show that the deformability of the bubbles plays a major role. Bubbles with a deformability comparable to what is seen experimentally (S. L. Ceccio, taken in the LCC) can lead to drag significant drag reduction, but only for a short time. The simulations clarified the turbulent modification and showed that less deformable bubbles can lead to an increase in the wall drag.



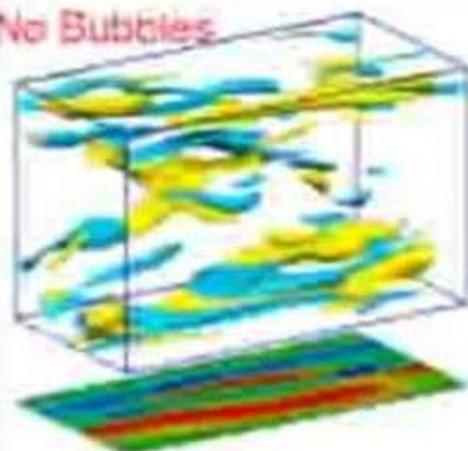
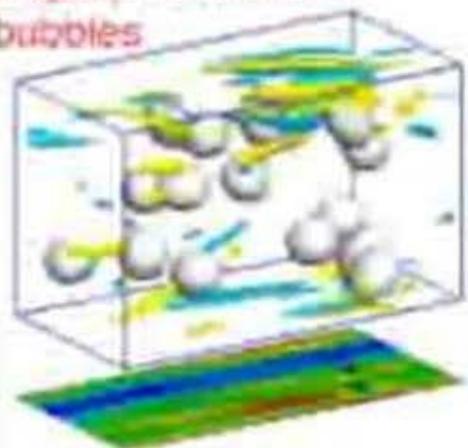
Experiment:

 $d^+ = 120$; $We = 0.686$

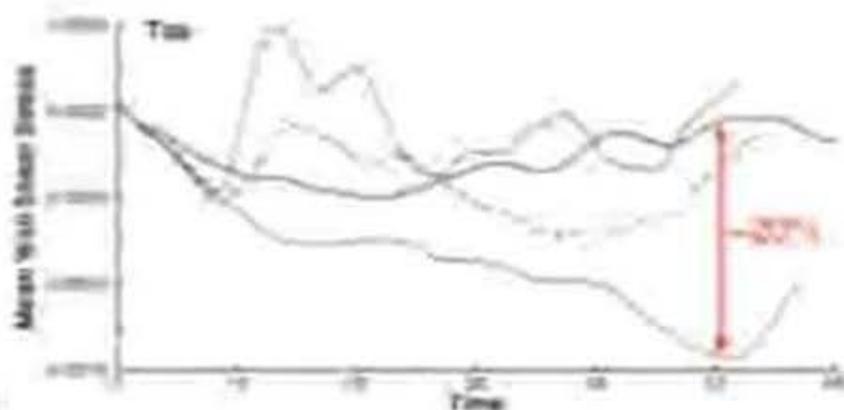
Simulations:

 $d^+ = 54$; $We = 0.203 - 0.405$


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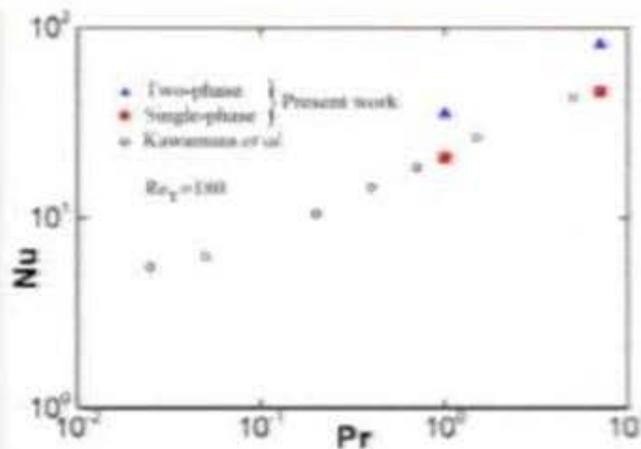
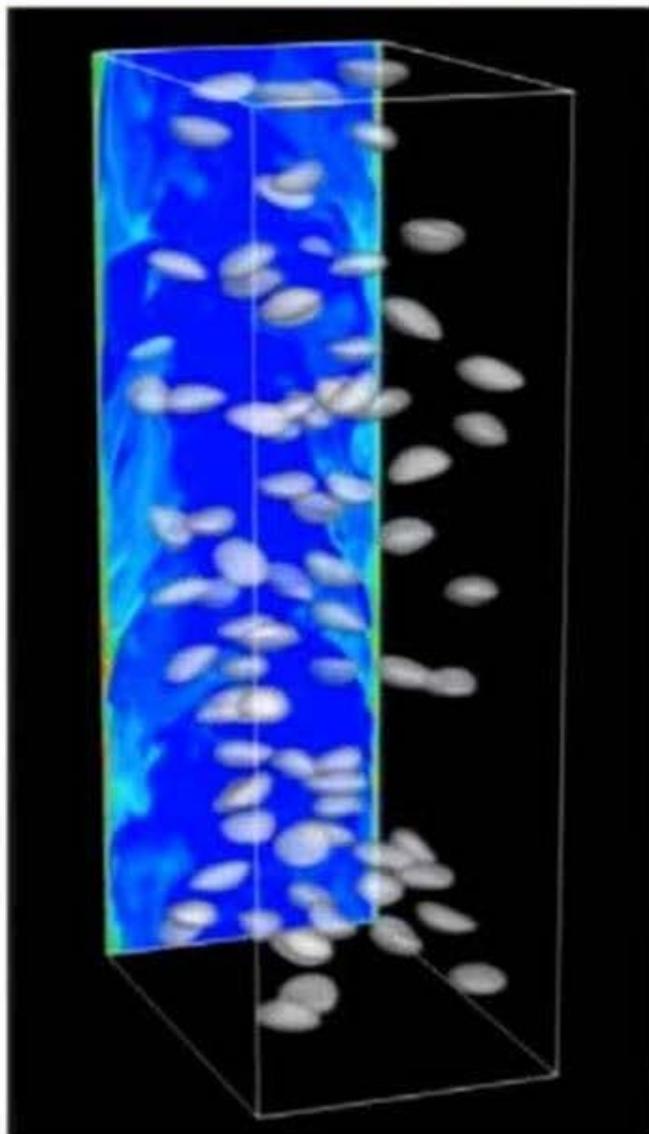


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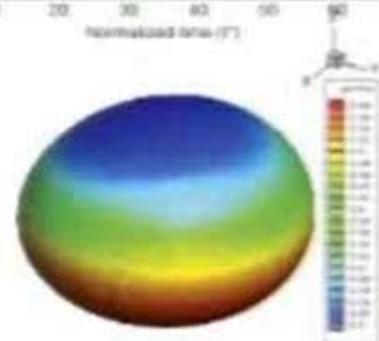
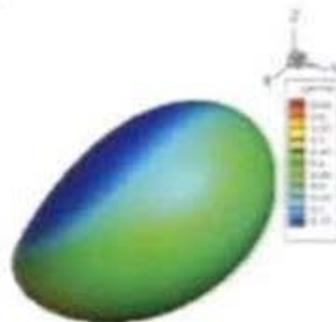
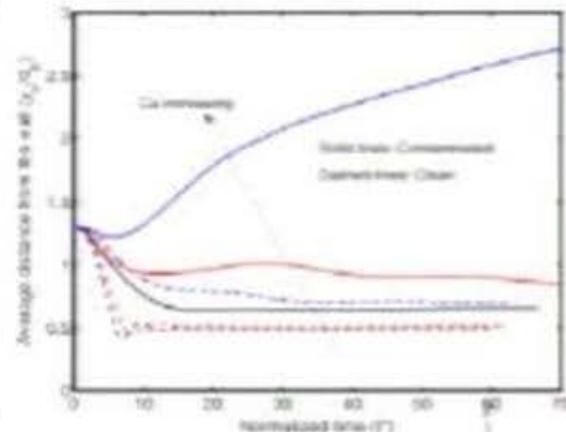
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S. Dabiri and G. Tryggvason.
Heat transfer in turbulent
bubbly flow in vertical
channels. *Chemical
Engineering Science*. 122
(2015), 106-113.

M. Muradoglu and G.
Tryggvason. Simulations of
Soluble Surfactants in 3D
Multiphase Flow. *Journal of
Computational Physics*. 274
(2014), 737-757.





Mining DNS Data for Closure Relations

With
Ming Ma & Jiakai Lu

A simple description of the average flow can be derived by integrating the vertical momentum equation and assuming that the density and viscosity of the gas is zero

Void fraction and phase averaged velocity

$$\alpha_i = \frac{1}{A_{zy}} \int \chi_i da \quad \langle v \rangle = \frac{1}{\alpha_i A_{zy}} \int \chi v da$$

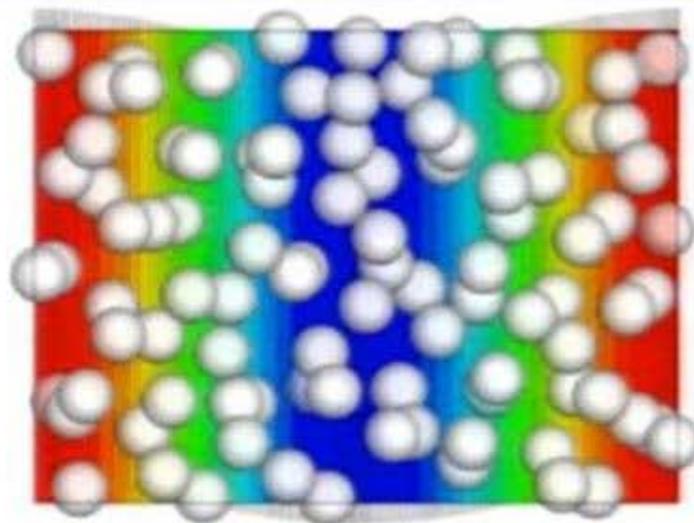
Horizontal flux of bubbles

$$\frac{\partial \alpha_i}{\partial t} + \frac{\partial F_i}{\partial x} = 0 \quad F_i = \frac{1}{\alpha A_{zy}} \int \chi_i u_i da$$

Averaged vertical momentum of the liquid:

$$\begin{aligned} \frac{\partial}{\partial t} \alpha_l \langle v_l \rangle + \frac{\partial}{\partial x} \alpha_l \langle v_l \rangle \langle u_l \rangle = & \\ - \frac{1}{\rho_l} \frac{dp}{dy} - \frac{1}{\rho_l} \alpha_l \rho_l g + & - \frac{1}{\rho_l} \frac{\partial}{\partial x} \left(\alpha_l \mu_l \frac{\partial v_l}{\partial x} \right) - \frac{\partial}{\partial x} \left(\alpha_l \langle uv \rangle_l \right) \end{aligned}$$

Plus surface
tension term



Obviously:

$$F_b = \alpha_b \langle u_b \rangle$$



“Closure” variables
needed for models
of the average flow

Resolved
average
variables

Quantities
summarizing the state
of the unresolved flow

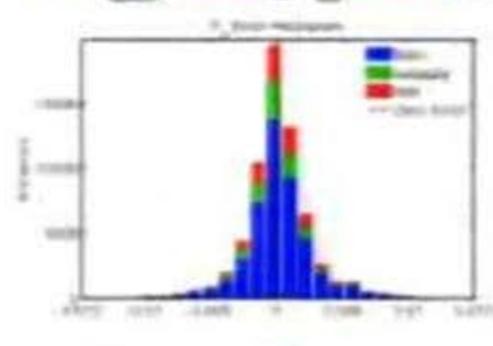
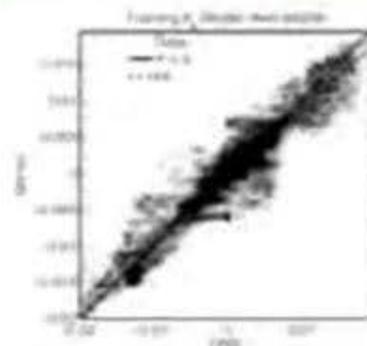
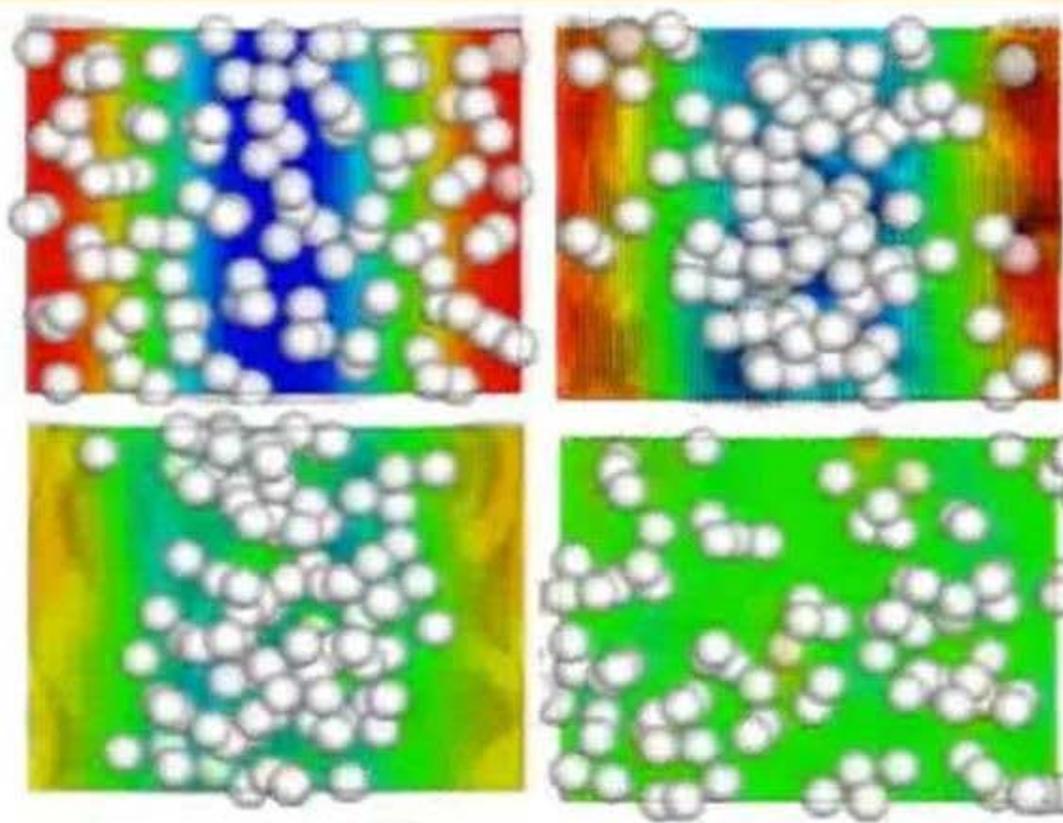
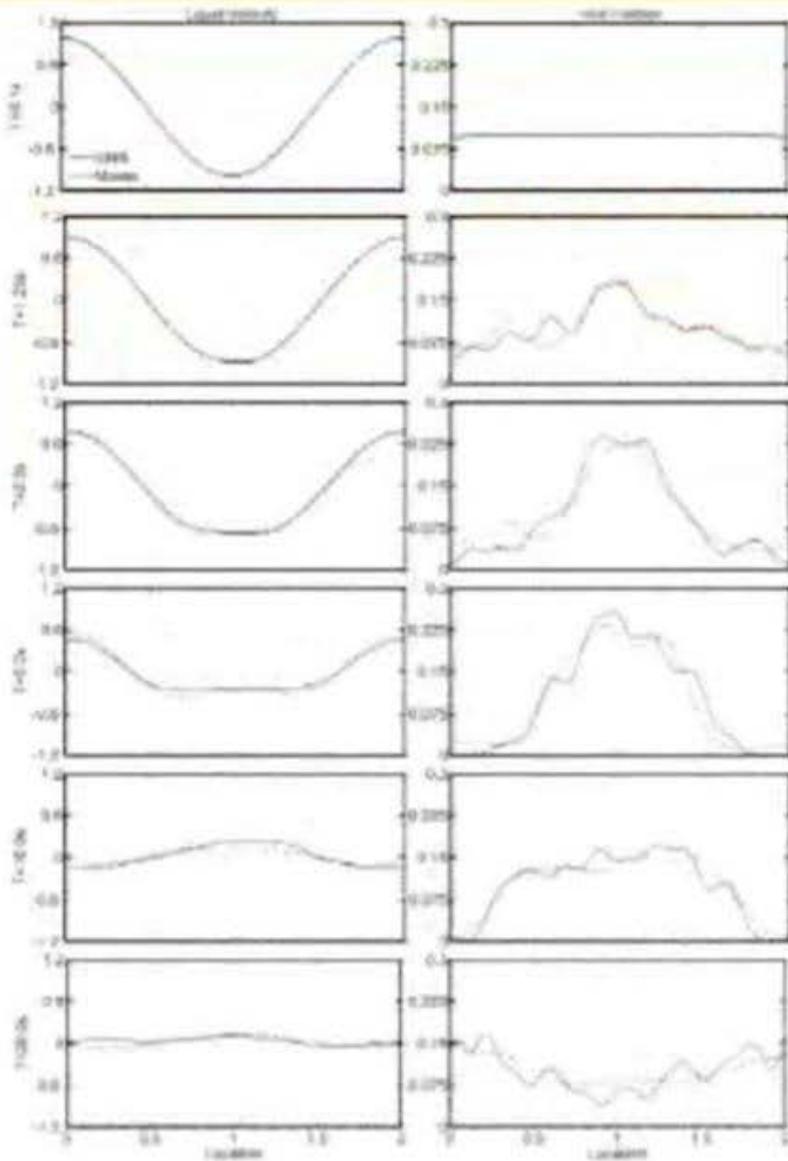
DNS data provides any quantity that we care to record or collect at any point in the flow, or as averages over any time or spatial region that we select. “Mining” the data allows us to correlate any variable with any others and quantify the accuracy of the fit.

F_g	$\langle u'v' \rangle$	α_g	$\frac{\partial \alpha_g}{\partial x}$	$\frac{\partial \langle v \rangle_l}{\partial x}$	d_w	k_t	ε_t	A	

We have used several data mining techniques to find the fit, including regression and neural networks

$$F_b = f\left(\alpha, \frac{\partial \alpha}{\partial x}, \frac{\partial \langle v \rangle}{\partial x}\right); \quad \langle u'v' \rangle = g\left(\alpha, \frac{\partial \alpha}{\partial x}, \frac{\partial \langle v \rangle}{\partial x}\right);$$

$$f(\mathbf{x}) = b_0 + \sum_{i=1}^n b_i h_i(\mathbf{x}); \quad h_i(\mathbf{x}) = s\left(a_{0i} + \sum_{j=1}^n x_j a_{i,j}\right); \quad s(u) = \frac{2}{1 + e^{-2u}} - 1$$





Capturing small-scale
processes using
Embedded Analytical
Descriptions

With

Bahman Aboulhasanzadeh and Siju Thomas

Capturing isolated small-scale motion in simulations where the focus is on the larger scales can be done in many ways, such as by various grid refinement techniques (unstructured grids, AMR for Cartesian grids, wavelets, etc.) or reduced order models

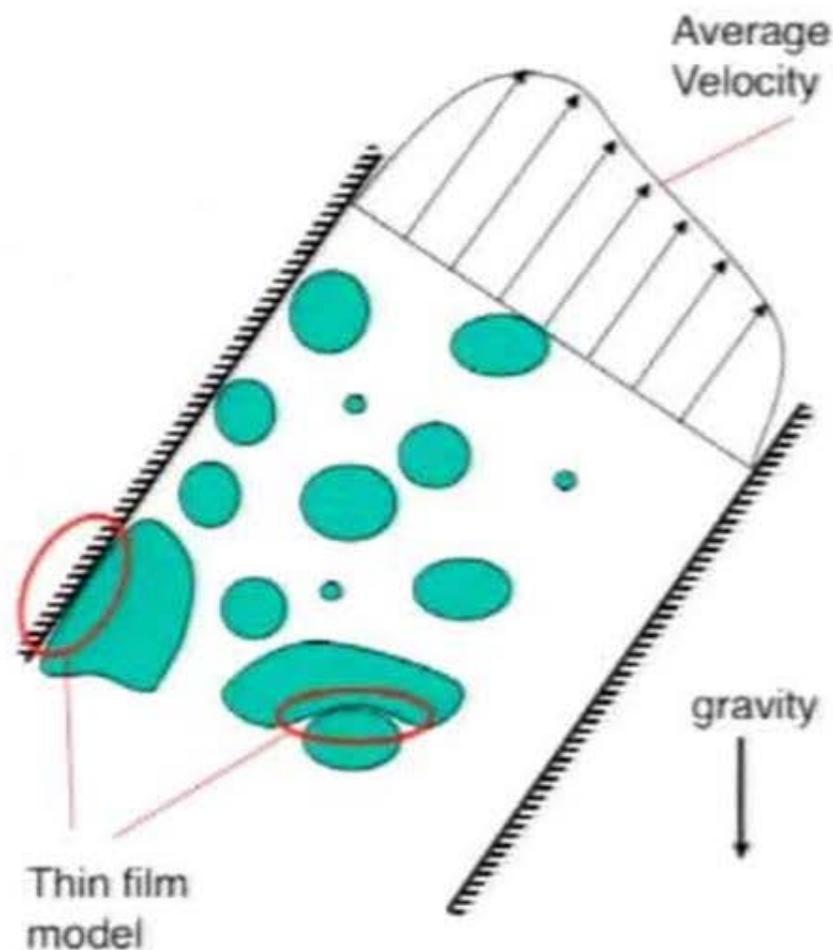
However:

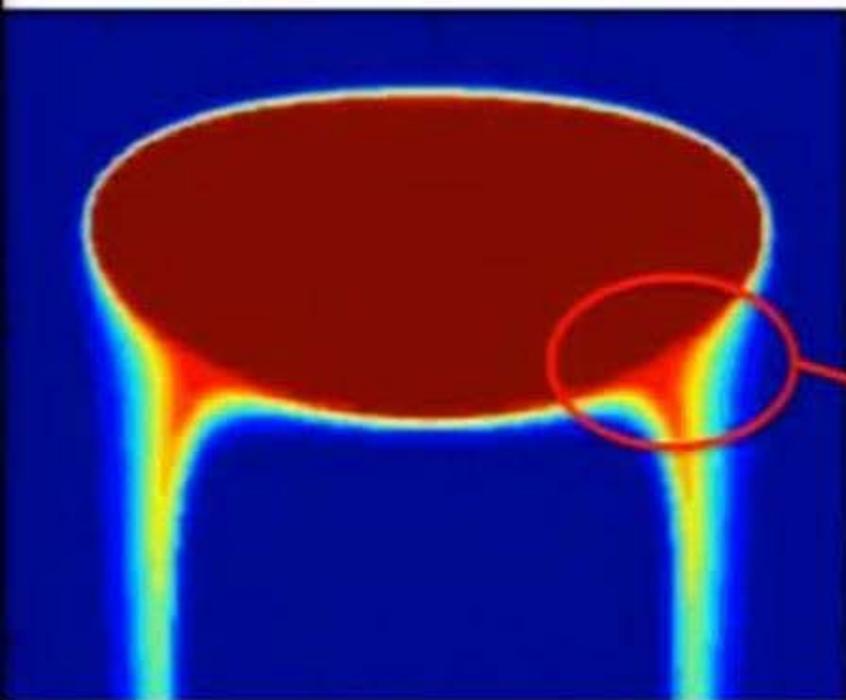
At small scales, the effect of surface tension is strong so interface geometries are simple

At small scales the effect of viscosity is strong so the flow is simple

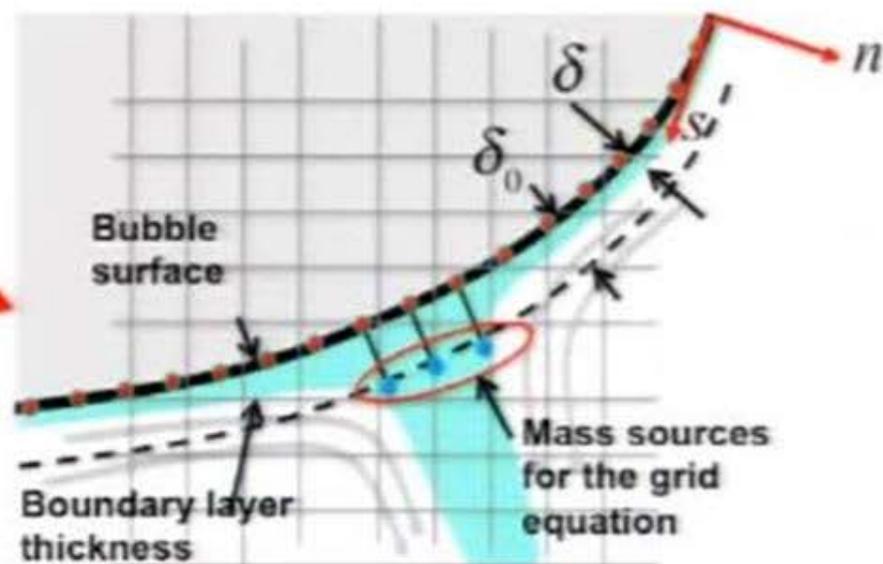
Those are exactly the situation that can be —and have been— handled analytically

Buoyant bubbles in an inclined channel flow





Capturing the mass boundary layer



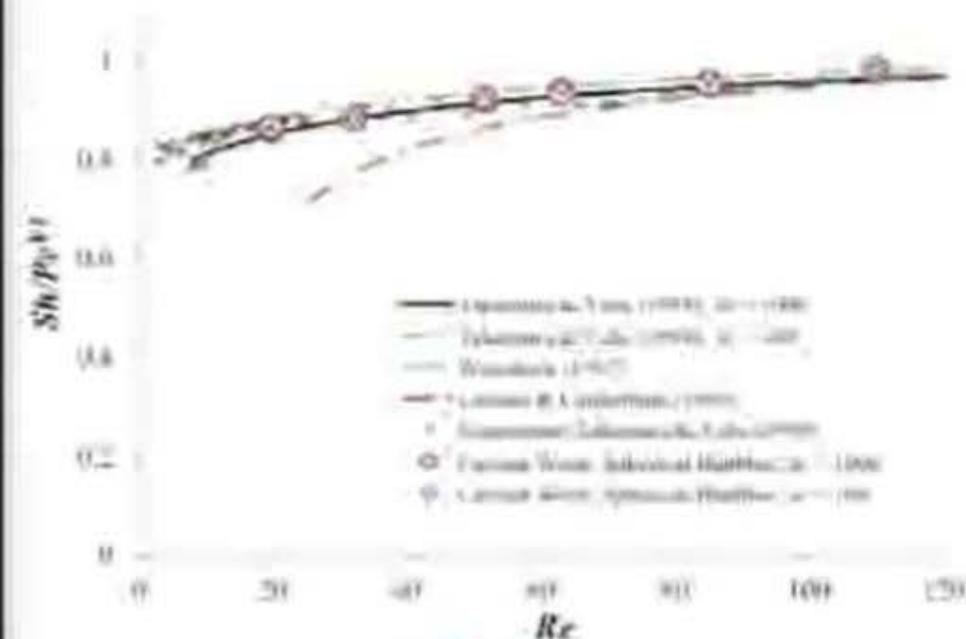
$$\frac{\partial f}{\partial t} = \sigma n \frac{\partial f}{\partial n} + D \frac{\partial^2 f}{\partial n^2}$$

$$M_0 = \int_0^\infty f \, dn$$

$$M_1 = \int_0^\infty n f \, dn$$

$$\frac{dM_0}{dt} = -\sigma (M_0 - \delta_o f_\delta) - D \left. \frac{\partial f}{\partial n} \right|_0$$

$$\frac{dM_1}{dt} = -\sigma (2M_1 - \delta_o^2 f_\delta) + D (f_o - f_\delta)$$

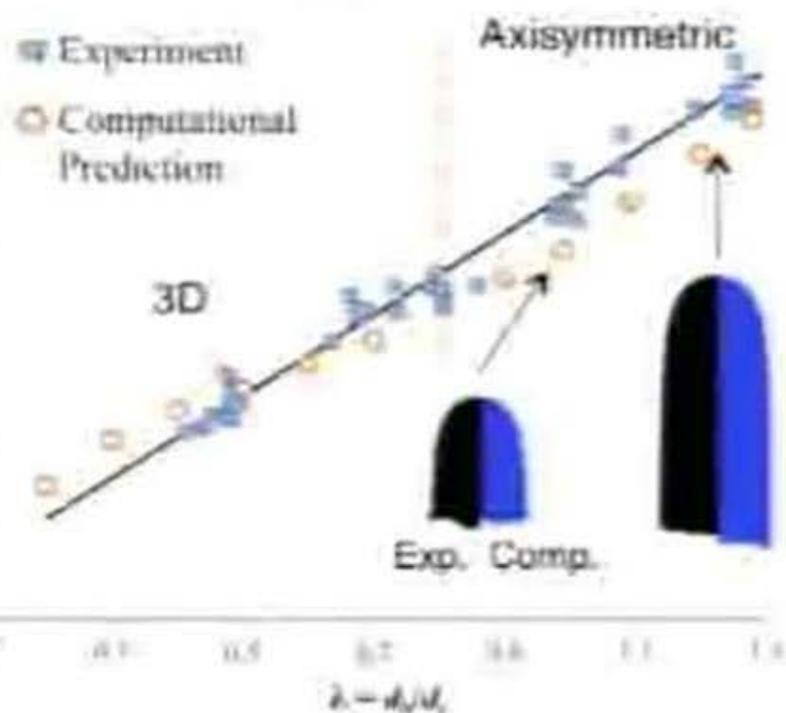


Comparison with experimental results from A. Tomiyama:

$$Eo = 24.7$$

$$Mo = 10^{-7.78} \text{ and}$$

$$Sc = 8260$$



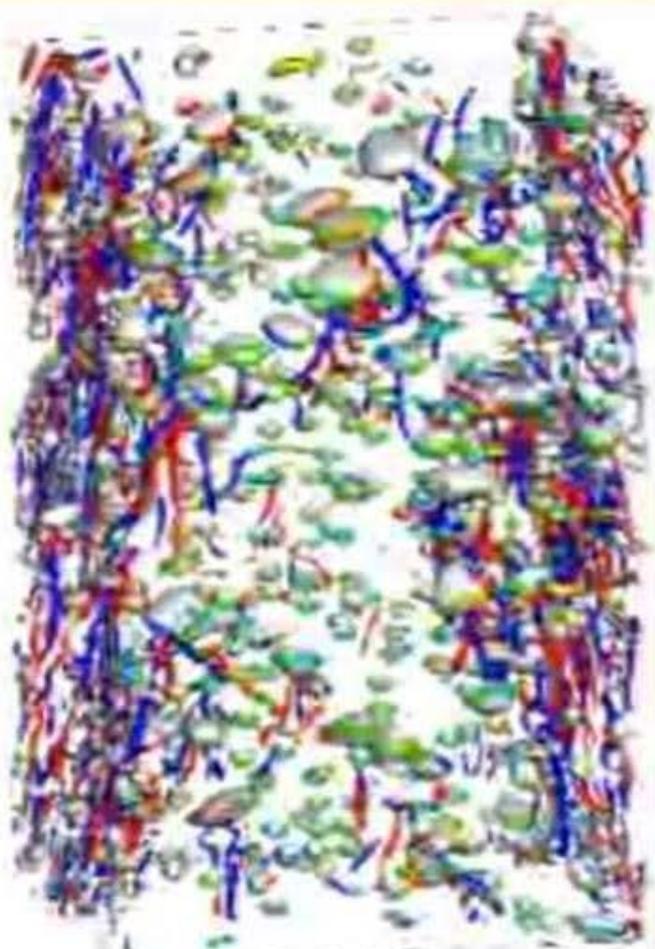
The mass transfer versus Re , for a single bubble in a large domain, along with the predictions of experimental correlations

DNS of several flows have been developed to the point that they should be able to help produce new models for "industrial" simulations. The availability of the data is putting new demands on the modeling of complex multiphase flows.

DNS needs to be extended to handle flows with more complex topology and those undergoing flow regime transitions

Complex isothermal flows and flows with phase change and other additional physics, such as mass transfer, need multiscale modeling that must be developed further and put on a rigorous theoretical basis. Need "almost" DNS

One of the biggest obstacle for more rapid increase in the use of DNS is the high "entry barrier" for new investigators. Many "things" to learn!



500 bubbles of different sizes
in channel flow with $Re^b=500$
1024 × 768 × 512 grid points
2048 processors on the Titan