

# Energy scaling law for a single disclination in a thin elastic sheet

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# Overview

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- 2 The single disclination – setup and main result
- 3 Sketch of the proof
- 4 More results

# Energy focusing in thin elastic sheets

- Experimental setup: Take a thin elastic sheet and subject it to boundary conditions that force it to form sharp folds or vertices.
- Rigorous formulation: simply connected Lipschitz domain  $\Omega \subset \mathbb{R}^2$ : elastic sheet in the undeformed configuration;  
 $y : \Omega \rightarrow \mathbb{R}^3$ : deformation map;  
 elastic energy:



$$\begin{aligned}
 I_h(y, \Omega) &:= \int_{\Omega} \left| Dy^T Dy - g_0 \right|^2 + h^2 \left| D^2 y \right|^2 \\
 &= \int_{\Omega} \left\| g_y - g_0 \right\|^2 + h^2 \left| D^2 y \right|^2
 \end{aligned}$$

( $g_y = y^* e^{(3)}$ ,  $e^{(3)}$  = Euclidean metric on  $\mathbb{R}^3$ ,  $g_0$  = reference metric.)

## Tensile boundary conditions

- Looking for lower and upper bounds, both as expansions in the small parameter  $h$
- Crumpling: *too hard*. Alternative problem: Impose *tensile boundary conditions*. I.e., choose Dirichlet boundary conditions such that there exists a unique (Lipschitz) isometric immersion, which is singular (infinite bending energy)
- Every deviation from the singular configuration is penalized by the membrane term; the balance between the two leads to an (optimal) lower bound
- If *short maps* (i.e., maps  $y$  with  $Dy^T Dy < Id$ ) are permissible by the boundary conditions, this method of proof breaks down

## Why are lower bounds difficult? (for non-tensile b.c.)

### Theorem (Nash '54, Kuiper '59)

Let  $(M, g)$  be a Riemannian manifold of dimension  $n$ ,  $m \geq n + 1$ . Any short immersion of  $M$  into  $\mathbb{R}^m$  can be uniformly approximated by isometric immersions of class  $C^1$ .

**Thus:** For a given reference metric  $g$ , there exists a huge set of maps that are very close to isometric immersions. The theory behind this is Gromov's  $h$ -principle. How to exclude these degrees of freedom?



## The single disclination

- Experimental setup: Remove a sector from a circular thin sheet and glue the edges back together
- Instead of “removing a sector” we may equivalently “shorten the reference metric in the angular direction”,



$$\hat{y}(x) = \sqrt{1 - \Delta^2|x|}x + \Delta|x|e_z$$
$$g_\Delta = D\hat{y}^T D\hat{y}$$

$$\Rightarrow g_\Delta(x) = \text{Id}_{2 \times 2} - \Delta^2 \hat{x}^\perp \otimes \hat{x}^\perp$$

$0 < \Delta < 1$ , domain =  $B_1 = B(0, 1) \subset \mathbb{R}^2$  with  
 $\hat{x} = x/|x|$ ,

$$\hat{x}^\perp = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \hat{x}$$

## Main result

$$g_{\Delta} = \text{Id} - \Delta^2 \hat{x}^{\perp} \otimes \hat{x}^{\perp}, \quad g_y = Dy^T Dy$$
$$I_{h,\Delta}(y) = \int_{B_1} |g_y - g_{\Delta}|^2 + h^2 |D^2 y|^2 dx$$

### Theorem (O. '15)

There exists a constant  $C = C(\Delta) > 0$  such that

$$2\pi\Delta^2 h^2 \left( |\log h| - \frac{3}{2} \log |\log h| \right) \\ \leq \inf_{y \in W^{2,2}(B_1; \mathbb{R}^3)} I_{h,\Delta}(y) \leq 2\pi\Delta^2 h^2 |\log h| + C$$

for all small enough  $h$ .

# An ansatz for lower bounds for the single disclination

The “right” curvature like quantity is  $\mathcal{K}(y) := \sum_{i=1}^3 \det D^2 y_i$ .

## Main idea

- Use  $\mathcal{K}(y)$  as a control variable for lower bounds of both, the membrane and bending term.
- To deal with the membrane term, use the formulation of the determinant  $\det D^2 v$  as a *very weak Hessian*,  
$$\det D^2 v = (v_{,1} v_{,2})_{,12} - \frac{1}{2}(|v_{,1}|^2)_{,22} - \frac{1}{2}(|v_{,2}|^2)_{,11}$$
- Hence, we have

$$\sum_{i=1}^3 \det D^2 y_i = (y_{,1} \cdot y_{,2})_{,12} - \frac{1}{2}(|y_{,1}|^2)_{,22} - \frac{1}{2}(|y_{,2}|^2)_{,11}$$



## Main idea continued

- $\implies \sum_i \det D^2 y_i - \det D^2 \hat{y}_i = (g_y - g_\Delta)_{12,12} - \frac{1}{2}(g_y - g_\Delta)_{11,22} - \frac{1}{2}(g_y - g_\Delta)_{22,11}$
- Explicit calculation yields  $\det D^2 \hat{y}_i = \pi \Delta^2 \delta_0$
- $\implies \|\sum_i \det D^2 y_i - \pi \Delta^2 \delta_0\|_{W^{-2,1}} \lesssim$  membrane energy, i.e., we obtain control over  $\sum_i \det D^2 y_i - \pi \Delta^2 \delta_0$  in  $W^{-2,1}$  through the membrane energy

## Lower bound for bending through $\int \sum_i \det D^2 y_i$

**Bending:** Note that by the coarea formula,

$\left| \int_{B_\rho} \sum_i \det D^2 y_i dx \right| \leq \sum_i \mathcal{L}^2(Dy_i(B_\rho))$ . By the isoperimetric inequality in  $\mathbb{R}^2$ , we can get a lower bound for  $\sum_i \int_{\partial B_\rho} |D^2 y_i|$ :

### Lemma

For  $v \in C^2(B_1)$  and  $0 \leq r \leq 1$ ,

$$\int_{\partial B_r} |D^2 v| d\mathcal{H}^1 \geq \sqrt{4\pi} \left| \int_{B_r} \det D^2 v dx \right|.$$

## More results: 3d elasticity

$$\begin{aligned} \Omega_h &:= B_1 \times [-h/2, h/2], \quad \hat{Y}(x, z) = \hat{y}(x) + z\nu_{\hat{y}} \\ g_{\Delta}^{(3)}(x, z) &= \text{Id}_{3 \times 3} - \Delta^2 \hat{x}^{\perp} \otimes \hat{x}^{\perp} \\ E_{h, \Delta}(Y) &= \int_{\Omega_h} \text{dist}^2 \left( DY(x), SO(3) \sqrt{g_{\Delta}^{(3)}(x)} \right) d\mathcal{L}^3(x). \end{aligned}$$

### Theorem (O., '15)

Let  $0 < \Delta < 1$ . There exists a constant  $C = C(\Delta)$  with the following property: For  $h$  small enough,

$$\frac{1}{C} |\log h| \leq \inf_{y \in W^{2,2}(\Omega_h; \mathbb{R}^3)} h^{-2} E_{h, \Delta}(y) \leq 2\pi \Delta^2 (|\log h| + C).$$

*Proof:* Translate lower bounds of  $I_{h, \Delta}(y)$  into bounds in 3d elasticity through Geometric Rigidity by Friesecke, James, Müller.

## More results: The Föppl-von Kármán model

$$I_{h,\Delta}^{\text{vK}} = \int_{B_1} \left| \text{sym } Du + \frac{1}{2} Dv \otimes Dv + \Delta e_\varphi \otimes e_\varphi \right|^2 + h^2 |D^2 v|^2 dx$$

### Theorem (O. '15)

There exists a constant  $C = C(\Delta) > 0$  with the following property: For  $h$  small enough,

$$4\pi\Delta (|\log h| - 2 \log |\log h|) - C \leq h^{-2} \inf_{(u,v) \in W^{1,2}(B_1; \mathbb{R}^2) \times W^{2,2}(B_1)} I_{h,\Delta}^{\text{vK}}(u, v) \leq 4\pi\Delta |\log h| + C.$$

## More results: Convergence of almost minimizers to the singular cone

$$\hat{y} : B_1 \rightarrow \mathbb{R}^3, \quad \hat{y}(x) = \sqrt{1 - \Delta^2}x + \Delta|x|e_3.$$

### Theorem (Müller, O., '15)

Let  $y^h \in W^{2,2}(B_1; \mathbb{R}^3)$  be a sequence with  $I_{h,\Delta}(y^h) \leq 2\pi\Delta^2 h^2(|\log h| + C)$ . Then up to Euclidean motions, we have for every  $0 < \rho < 1$ ,

$$y^h \rightharpoonup \hat{y} \quad \text{in } W^{2,2}(B_1 \setminus B_\rho; \mathbb{R}^3).$$

Important tools in the proof: Lower bounds for  $I_{h,\Delta}(y)$  + Structure result for flat  $W^{2,2}$  surfaces (Pakzad'04, Hornung '11)

## Summary/Outlook

- Main results: Lower bounds for a single disclination in a geometrically fully nonlinear plate model, in 3-dimensional nonlinear elasticity and in the Föppl-von Kármán model. Moreover, minimizers of the elastic energy converge to the singular cone as  $h \rightarrow 0$  weakly in  $W_{loc}^{2,2}(B_1 \setminus \{0\}; \mathbb{R}^3)$
- Key to the first result is to look at a curvature-like quantity, that gives control over certain properties of the graph of  $Dy$ . The lower bound follows by an isoperimetric inequality
- Can this idea be applied to other problems of a similar flavor? (D-cones? Non-conical geometries?)

Literature:

H. O., Energy scaling law for a single disclination in a thin elastic sheet. Preprint, 2015. Available at arXiv:1509.07378.

H. O., Energy scaling law for the regular cone. *J. Nonlinear Sci.*, 2015. Available at arXiv:1502.07013.