Conclusions

Mathematical relations between LCS detection methods SIAM Conference on Applications of Dynamical Systems

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Daniel Karrasch, joint with Johannes Keller (TUM) Mathematical relations between LCS methods ▲□▶ ▲□▶ ▲臣▶ ▲臣▶ 三臣 - のへで

Motivation: Study transport and mixing of scalar quantities

Hypotheses:

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- ► In purely advective flows, stretching is irrelevant, because advection is reversible.
- ► Caring about stretching implies that we care about advection-diffusion.

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Fokker–Planck equation (FPE) (Eulerian/spatial evolution equation):

$$\partial_t u + \operatorname{div}(u \cdot v) = \varepsilon \Delta_g u$$

- transport/advection along velocity field v
- mixing/diffusion according to spatial geometry/distances modeled by metric g

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Changing perspectives: looking from the material





Pullback geometry: stretching and diffusivity

 length measurement via Cauchy–Green strain tensor C



Pullback geometry: stretching and diffusivity

- length measurement via Cauchy–Green strain tensor C
- directional diffusivity via inverse CG strain tensor C⁻¹



Pullback geometry: stretching and diffusivity

- length measurement via Cauchy–Green strain tensor C
- directional diffusivity via inverse CG strain tensor C⁻¹
- ► strong stretching ↔ weak Lagrangian diffusion
- diffusion in Lagrangian coordinates is anisotropic



Seek closed curves γ , parametrized by r, that stationarize averaged tangential strain

$$\delta Q(\gamma) \stackrel{!}{=} 0, \qquad \qquad Q(\gamma) = rac{1}{|\gamma|} \int_{\gamma} \sqrt{rac{\langle r', Cr'
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$$\langle r', Cr' \rangle = \langle r', \Omega^{\top} C^{-1} \Omega r' \rangle = \underbrace{\langle \Omega r', C^{-1} \Omega r' \rangle}_{\text{diffusive flux through } \gamma}, \qquad \Omega = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \quad \det(C) = 1$$

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specific solution: $r' = \frac{1}{\sqrt{1 + \lambda_{\min}(C)}} v_{\min}(C) \pm \frac{1}{\sqrt{1 + \lambda_{\max}(C)}} v_{\max}(C)$

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Seek closed curves γ , parametrized by r, that stationarize averaged tangential strain

$$\begin{split} \delta Q(\gamma) \stackrel{!}{=} 0, \qquad Q(\gamma) = \frac{1}{|\gamma|} \int_{\gamma} \sqrt{\frac{\langle r', Cr' \rangle}{\langle r', r' \rangle}} &= \frac{1}{|\gamma|} \int_{\gamma} \frac{\|DF \cdot r'\|}{\|r'\|} \\ \langle r', Cr' \rangle = \langle r', \Omega^{\top} C^{-1} \Omega r' \rangle &= \underbrace{\langle \Omega r', C^{-1} \Omega r' \rangle}_{\text{diffusive flux through } \gamma}, \qquad \Omega = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \quad \det(C) = 1 \\ \text{specific solution:} \quad r' = \frac{1}{\sqrt{1 + \lambda_{\min}(C)}} v_{\min}(C) \pm \frac{1}{\sqrt{1 + \lambda_{\max}(C)}} v_{\max}(C) \\ \text{If, as often observed "around" vortices, } \lambda_{\max}(C) \gg 1, \text{ then } \lambda_{\min}(C) \ll 1 \text{ and} \\ r' \wr v_{\min}(C), \qquad \text{or equivalently} \qquad \gamma^{\perp} \wr v_{\max}(C) \end{split}$$

Black hole vortices admit very small (Lagrangian) cross-diffusivity.

Lagrangian model of advection-diffusion

Eulerian:

$$\partial_t u + \operatorname{div}(u \cdot v) = \varepsilon \Delta_g u$$

Equivalent Lagrangian FPE (material evolution equation) [Thiffeault 2003]

$$\partial_t w = \varepsilon \Delta_{g(t)} w$$

 $g(t) \coloneqq \Phi(t)^* g = C$ —pullback metric, CG strain tensor, $\Delta_{g(t)}$ —Laplace–Beltrami operator. $\Rightarrow \text{ evolving (material) manifold}$

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Definition (Lagrangian Coherent Structures)

maximal material sets with minimal diffusive flux, or, metastable/almost-invariant sets under Lagrangian FPE

The dynamic Laplacian as a proxy for $\Delta_{g(t)}$

Lagrangian FPE: time-dependent & anisotropic diffusion(-only) equation

$$\partial_t w = \varepsilon \Delta_{g(t)} w$$

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Possible simplification: Autonomize $t \mapsto \Delta_{g(t)}$ by time-averaging

$$\left|\overline{\Delta} = \frac{1}{T} \int_0^T \Delta_{g(t)} dt - \text{dynamic Laplacian}_{[Froyland2015]}\right|$$

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Lemma (Froyland 2015, DK & Keller 2016, Froyland & Kwok 2016)

 $\overline{\Delta}$ is an elliptic, nonpositive, selfajoint, 2nd-order differential operator.

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Theorem (spectral convergence to probabilistic transfer operator $\mathcal{L}_{\varepsilon}$, DK & Keller 2016)

$$\mathcal{L}_{\varepsilon}^{*}\mathcal{L}_{\varepsilon} = I + c\varepsilon^{2}\overline{\Delta} + \mathcal{O}\left(\varepsilon^{4}
ight)$$

Small example: transient double gyre [Mosovsky & Meiss, 2011], figure courtesy of [Froyland, Nonlinearity, 2015]



solid: level set of 2nd eigenfunction of dynamic Laplacian, minimizing

 $h^{\mathrm{D}}(\Gamma) = 8.2$

Small example: transient double gyre [Mosovsky & Meiss, 2011], figure courtesy of [Froyland, Nonlinearity, 2015]



solid: level set of 2nd eigenfunction of dynamic Laplacian, minimizing $h^{D}(\Gamma) = 8.2$ dashed: outermost closed λ -line $h^{D}(\Gamma) = 6.7$

Lagrangian heat flow, initialized within coherent structure

Cylinder flow [Froyland et al., 2010]

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Lagrangian heat flow, initialized within mixing region

Animation 1000 times slower than previous one!

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Transport and mixing	Lagrangian perspective	Lagrangian diffusion equation	Conclusions
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- ▶ math. advection-diffusion-based relation b/w BH vortices and probabilistic TO
- ▶ within this framework strongly related objectives: minimizing Lagr. cross-diffusion
- discretization and implementation details may matter

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D. Karrasch and J. Keller, preprint, 2016, available on ResearchGate & arXiv.