

Predicting Extreme Events for Passive Scalar Turbulence through Reduced-Order Models

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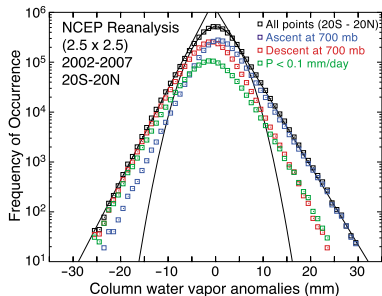
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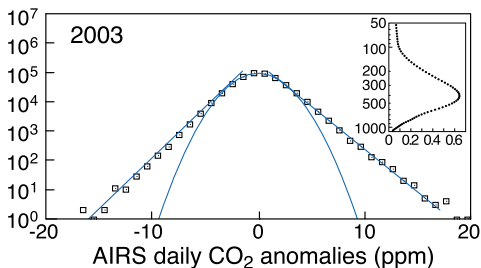
Long tails of natural tracers from observation

Intermittency is a key feature of turbulent flows and correctly taking it into account in turbulence models is of critical importance.

- Passive tracer distributions exhibit approximately **exponential tails**.
- Such tails are indeed **ubiquitous** in observation, model, and reanalysis data sets for a variety of tracers.¹



(a) Distribution of water vapor anomalies



(b) Distribution of CO2 anomalies

¹Neelin et al, *Geophys. Res. Lett.*, 2010

Simplified passive scalar models

● **Simplified models in tracer turbulence:**

- ▶ Understanding the **fundamental mechanisms** that can potentially lead to such behavior is a question of great theoretical interest.
- ▶ The analysis of **simple systems with a Gaussian core** can shed some light on similar behaviors observed in more complex models.

● **Major Questions:**

- ▶ What structure is needed for a velocity field so that the PDF for a passive scalar exhibits a transition from a Gaussian PDF to a broader than Gaussian shape.
- ▶ How to develop explicit reduced-order models with unambiguous behavior for intermittency of scalar PDFs.

Outline

- 1 **Turbulent diffusion models with mean gradient formulation**
 - Elementary models with intermittency for passive scalars
 - Rigorous intermittency in a random resonance regime
- 2 **Two-layer baroclinic turbulence as advection flow**
 - Two-layer baroclinic turbulence in ocean and atmosphere regimes
 - A reduced-order stochastic model with consistency and sensitivity
- 3 **Predicting intermittent PDFs with low-order stochastic models**
 - Model calibration in Gaussian velocity field
 - Prediction skill of low-order stochastic models

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Passive tracer with a mean gradient

Passive tracer of turbulent advection, diffusion, and damping

$$\frac{\partial T}{\partial t} + \mathbf{v} \cdot \nabla T = -d_T T + \kappa \Delta T.$$

- The tracer field is assumed to have a background mean gradient

$$T(\mathbf{x}, t) = \alpha_1 x + \alpha_2 y + T'(x, y, t).$$

- The model for the velocity field is two dimensional and periodic in space

$$\mathbf{v}(\mathbf{x}, t) = (U(t), v(x, t)),$$

with the *cross-sweep*, $U(t)$, and the *shear flow*, $v(x, t)$.

- $T'(x, t)$ denotes fluctuations around the mean gradient

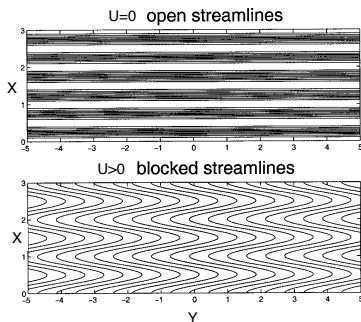
$$\frac{\partial T'}{\partial t} + U(t) \frac{\partial T'}{\partial x} = -\alpha v(x, t) - d_T T' + \kappa \Delta T'.$$

Elementary model with intermittency²

It is the time modulation in the transverse sweep, more precisely the fact that *the sweep crosses zero periodically*, that leads to the intermittency.

$$\frac{\partial T'}{\partial t} + \text{Pe}U(t) \frac{\partial T'}{\partial x} - \frac{\partial^2 T'}{\partial x^2} = -\alpha \text{Pe}v(x,t), \quad \boxed{\text{Pe} = VL/\kappa}.$$

$$T = \alpha y + T', \quad U(t) = \sin \omega t, \quad v(x,t) = \sin 2\pi x.$$



- when $U = 0$, the open streamlines in the horizontal direction, along the mean gradient αy , promote strong mixing by diffusion.
- when $U \neq 0$, the transverse sweep corresponds to blocked streamlines, little transport along the gradient, little opportunity for mixing by diffusion.

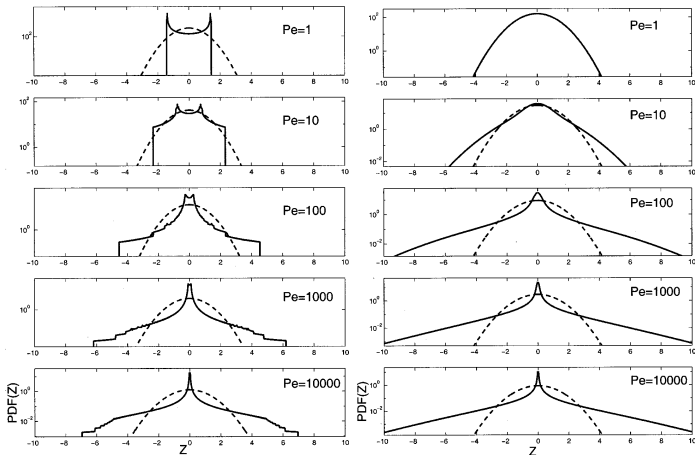
²Bourlioux & Majda, *Physics of Fluids*, 2002

Elementary model with intermittency

$$\frac{\partial T'}{\partial t} + \text{Pe}U(t) \frac{\partial T'}{\partial x} - \frac{\partial^2 T'}{\partial x^2} = -\alpha \text{Pe}v(x, t),$$

$$i) v = \sin 2\pi x,$$

$$ii) v = \xi(t), \langle \xi(t) \xi(0) \rangle = R(t).$$



(c) deterministic

(d) random Gaussian

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The two-layer flow as advection flow

The **two-layer quasi-geostrophic model** is one simple but fully nonlinear fluid model capable in capturing the essential physics.

Two-layer model

$$\begin{aligned}\frac{\partial q_\psi}{\partial t} + J(\psi, q_\psi) + J(\tau, q_\tau) + \beta \frac{\partial \psi}{\partial x} + U \frac{\partial \Delta \tau}{\partial x} &= -\frac{\kappa}{2} \Delta(\psi - \tau) - \nu \Delta^s q_\psi, \\ \frac{\partial q_\tau}{\partial t} + J(\psi, q_\tau) + J(\tau, q_\psi) + \beta \frac{\partial \tau}{\partial x} + U \frac{\partial}{\partial x} (\Delta \psi + k_d^2 \psi) &= -\frac{\kappa}{2} \Delta(\tau - \psi) - \nu \Delta^s q_\tau.\end{aligned}$$

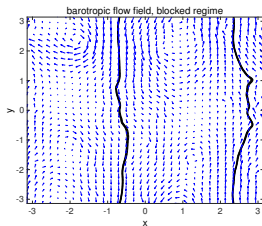
barotropic and baroclinic modes:

$$\begin{aligned}q_\psi &= \Delta \psi, \quad \psi = \frac{1}{2} (\psi_1 + \psi_2), \\ q_\tau &= \Delta \tau - k_d^2 \tau, \quad \tau = \frac{1}{2} (\psi_1 - \psi_2).\end{aligned}$$

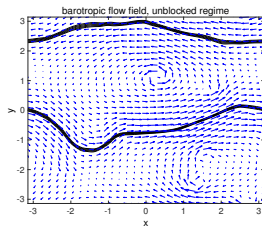
energy and heat flux:

$$\begin{aligned}E &= \frac{1}{2} \int |\nabla \psi|^2 + |\nabla \tau|^2 + k_d^2 \tau^2, \\ H_f &= k_d^2 U \int \nu \tau.\end{aligned}$$

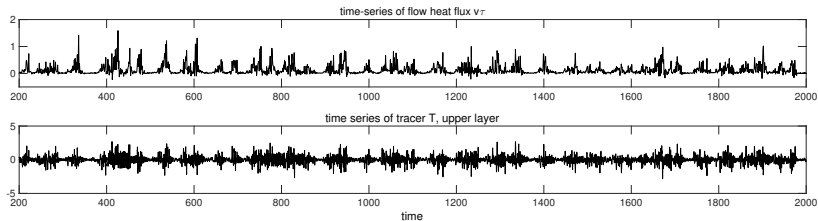
Illustration about tracer intermittency with heat flux



(a) blocked regime



(b) unblocked regime



(c) time-series of heat flux and scalar tracer

Stochastic formulations for the passive tracer and turbulent flow field

- **reduced-order advection flow equations** $\vec{q}_{\mathbf{k}} = (\hat{q}_{1,\mathbf{k}}, \hat{q}_{2,\mathbf{k}})^T$

$$d\vec{q}_{M,\mathbf{k}} = -(\gamma_{q,\mathbf{k}} + i\omega_{q,\mathbf{k}})\vec{q}_{M,\mathbf{k}}dt - D_{q,\mathbf{k}}^M \vec{q}_{M,\mathbf{k}}dt + \Sigma_{q,\mathbf{k}}^M d\vec{W}_{q,\mathbf{k}}, \quad 1 \leq |\mathbf{k}| \leq M.$$

$$\mathbf{v}_M = \nabla^\perp \vec{\psi}_M, \quad \vec{q}_{M,\mathbf{k}} = H_{\mathbf{k}} \vec{\psi}_{M,\mathbf{k}},$$

$$-\frac{1}{2} \sum_{\mathbf{m}+\mathbf{n}=\mathbf{k}} (A_{\mathbf{k}\mathbf{m}} \vec{q}_{\mathbf{m}} \circ \vec{q}_{\mathbf{n}} + A_{\mathbf{k}\mathbf{n}} \vec{q}_{\mathbf{n}} \circ \vec{q}_{\mathbf{m}}) \rightarrow -D_{q,\mathbf{k}}^M \vec{q}_{M,\mathbf{k}}dt + \Sigma_{q,\mathbf{k}}^M d\vec{W}_{q,\mathbf{k}}.$$

- **reduced-order passive tracer equations** $\vec{T}_{\mathbf{k}} = (\hat{T}_{1,\mathbf{k}}, \hat{T}_{2,\mathbf{k}})^T$

$$d\vec{T}_{M,\mathbf{k}} + \left(\mathbf{v}_M \cdot \nabla \vec{T}_M \right)_{\mathbf{k}} d\tilde{t} = \alpha \Gamma_{\mathbf{k}} \vec{\psi}_{M,\mathbf{k}} d\tilde{t} - (\gamma_{T,\mathbf{k}} + i\omega_{T,\mathbf{k}}) \vec{T}_{M,\mathbf{k}} d\tilde{t}, \quad 1 \leq |\mathbf{k}| \leq M.$$

$$\mathbf{v}_M = \sum_{|\mathbf{k}| \leq M_1} i\mathbf{k}^\perp \vec{\psi}_{M,\mathbf{k}} e^{i\mathbf{k} \cdot \mathbf{x}}.$$

Model calibration in the advection flow field

- The imperfect model calibration parameters (D_q^M, Σ_q^M) need properly reflect the true nonlinear energy mechanism.
- The statistical equations for the fluctuation covariance matrix $R_{\mathbf{k}}^q = \langle \vec{q}_{\mathbf{k}} \vec{q}_{\mathbf{k}}^* \rangle$ become⁴

$$\frac{dR_{\mathbf{k}}^q}{dt} + Q_F^q = (\mathcal{L}_{\mathbf{k}}^q + \mathcal{D}_{\mathbf{k}}^q) R_{\mathbf{k}}^q + R_{\mathbf{k}}^q (\mathcal{L}_{\mathbf{k}}^q + \mathcal{D}_{\mathbf{k}}^q)^*, \quad |\mathbf{k}| \leq N,$$

$$Q_F^q(\vec{q}_{\mathbf{k}}) = \frac{1}{2} \sum_{\mathbf{m}+\mathbf{n}=\mathbf{k}} \langle (A_{\mathbf{k}\mathbf{m}} \vec{q}_{\mathbf{m}} \circ \vec{q}_{\mathbf{n}} + A_{\mathbf{k}\mathbf{n}} \vec{q}_{\mathbf{n}} \circ \vec{q}_{\mathbf{m}}) \vec{q}_{\mathbf{k}}^* \rangle.$$

- In the statistical steady state, $dR_{\mathbf{k}}^q/dt = 0$, then the nonlinear fluxes can be calculated at equilibrium as $t \rightarrow \infty$

$$Q_{F,\text{eq}}^q = (\mathcal{L}_{\mathbf{k}}^q + \mathcal{D}_{\mathbf{k}}^q) R_{\mathbf{k},\text{eq}}^q + R_{\mathbf{k},\text{eq}}^q (\mathcal{L}_{\mathbf{k}}^q + \mathcal{D}_{\mathbf{k}}^q)^*.$$

⁴Sapsis & Majda, *Physica D*, 2013

Imperfect model correction from equilibrium statistics and additional terms

- the first proposal for the linear damping and Gaussian random noise correction can be introduced as

$$D_{q,\mathbf{k}}^{\text{eq}} = -\frac{1}{2}Q_{F,\text{eq},\mathbf{k}}^{q,-} \left(R_{\mathbf{k},\text{eq}}^q\right)^{-1}, \quad \Sigma_{q,\mathbf{k}}^{\text{eq}} = \left(Q_{F,\text{eq},\mathbf{k}}^{q,+}\right)^{1/2};$$

- a further correction for the noise and damping with a simple constant damping and noise

$$Q_{M,\mathbf{k}}^{\text{add}} = -D_M^{\text{add}}R_{M,\mathbf{k}} + \left(\Sigma_{M,\mathbf{k}}^{\text{add}}\right)^2, \quad D_M^{\text{add}} = \begin{bmatrix} d_M + i\omega_M & \\ & d_M - i\omega_M \end{bmatrix}.$$

- Combining the ideas, we propose the additional damping and noise corrections for the reduced-order flow vorticity mode

$$D_{q,\mathbf{k}}^M = -\frac{1}{2}Q_{F,\text{eq},\mathbf{k}}^{q,-} \left(R_{\mathbf{k},\text{eq}}^q\right)^{-1} - D_M^{\text{add}}, \quad \Sigma_{q,\mathbf{k}}^M = \left(Q_{F,\text{eq},\mathbf{k}}^{q,+} + \left(\Sigma_{M,\mathbf{k}}^{\text{add}}\right)^2\right)^{1/2}.$$

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Gaussian velocity stochastic models for passive tracer statistics

- Apply the Galerkin truncation strategy by resolving only the **first M leading modes in spectral space**

$$\frac{d\hat{u}_k^M}{dt} = (-\gamma_{u_k} + i\omega_{u_k}) \hat{u}_k^M + \sigma_{u_k} \dot{W}_k, \quad |k| \leq M.$$

- One of the simplest and most direct way to estimate the undetermined coefficients $\gamma_{u_k}, \omega_{u_k}, \sigma_{u_k}$ is through the **mean stochastic model (MSM)**:

$$E_k \equiv \text{var}(\hat{u}_k(t)) = \langle |\hat{u}_k(t) - \langle \hat{u}_k \rangle|^2 \rangle,$$
$$R_k = \int \mathcal{R}_k(t) \equiv \int \frac{\langle (\hat{u}_k(\tau) - \langle \hat{u}_k \rangle) (\hat{u}_k(\tau+t) - \langle \hat{u}_k \rangle)^* \rangle}{\text{var}(\hat{u}_k(\tau))}.$$

Spectral Information criterion for improving imperfect model prediction skill

A natural way of measuring the lack of information is the relative entropy

$$\mathcal{P}(\pi, \pi^M) = \int \pi \log \frac{\pi}{\pi^M}.$$

Khinchin's formula: if the autocorrelation function $\mathcal{R}(t)$ is smooth and rapid-decay,

$$u(t) = \int_{-\infty}^{\infty} e^{i\lambda t} \hat{Z}(d\lambda), \quad \mathcal{R}(t) = \int_{-\infty}^{\infty} e^{i\lambda t} dF(\lambda) = \int_{-\infty}^{\infty} e^{i\lambda t} E(\lambda) d\lambda,$$

Energy spectrum can be represented by $E(\lambda)$ or $dF(\lambda)$

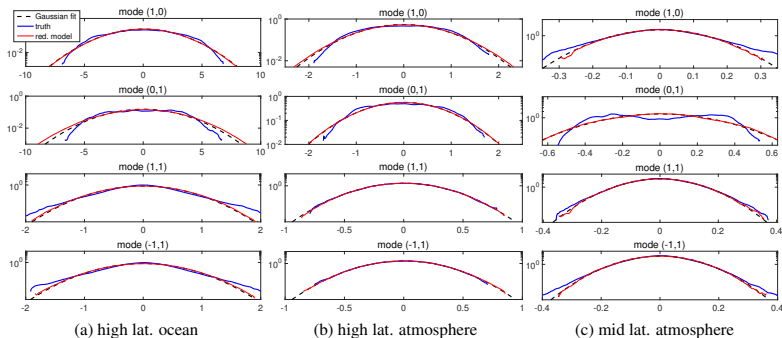
$$dF(\lambda) = E(\lambda) d\lambda = \mathbb{E} |\hat{Z}(d\lambda)|^2.$$

Spectral information criterion

- A) $\max |\mathcal{R}(t) - \mathcal{R}^M(t)| \leq 2\sqrt{3} (\int (E^2(\lambda) + E_M^2(\lambda)) d\lambda)^{1/2} \mathcal{P}(E(\lambda), E^M(\lambda))^{1/2};$
B) $\int |\mathcal{R}(t) - \mathcal{R}^M(t)|^2 dt \leq 12 \max |E^2(\lambda) + E_M^2(\lambda)| \mathcal{P}(E(\lambda), E^M(\lambda)).$

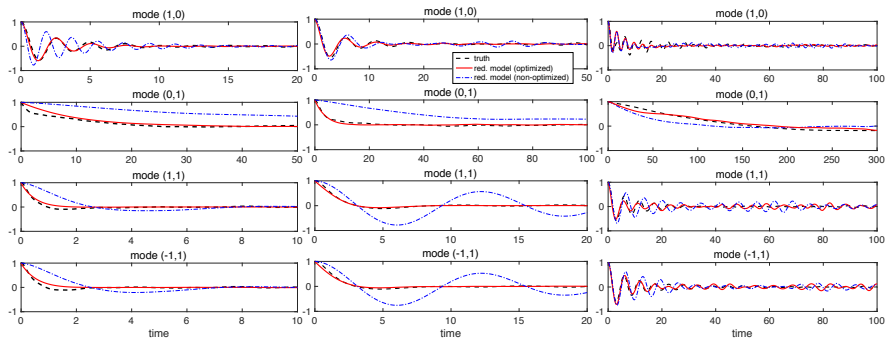
PDFs in stream functions of the advection flow

regime	N	β	k_d	U	r	d_T	κ	α
ocean, high lat.	128	2	40	0.1	0.1	0.5	0.001	1
atmosphere, high lat.	128	2	4	0.2	0.2	0.1	0.001	1
atmosphere, mid lat.	128	1	4	0.2	0.1	0.1	0.001	1



Autocorrelation functions for representative modes

$$\mathcal{R}_k(t) = \frac{\langle (\hat{u}_k(\tau) - \langle \hat{u}_k \rangle)(\hat{u}_k(\tau+t) - \langle \hat{u}_k \rangle)^* \rangle}{\text{var}(\hat{u}_k(\tau))}$$



(a) high lat. ocean

(b) high lat. atmosphere

(c) mid lat. atmosphere

Skill of stochastic models in two-layer QG

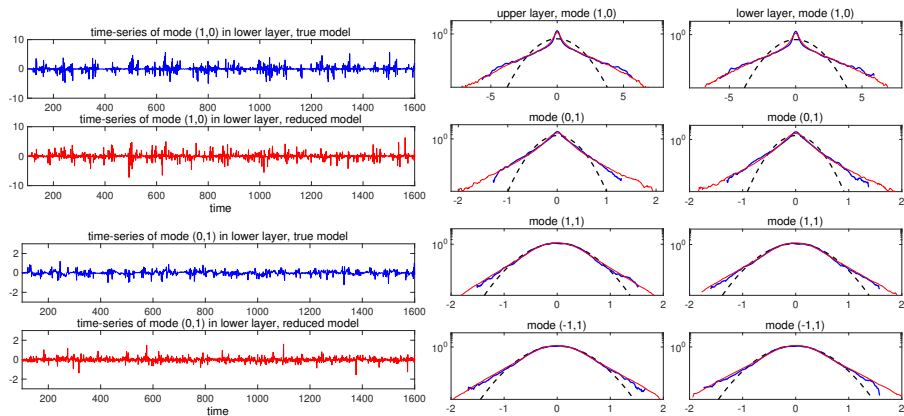


Figure: Prediction about tracer intermittency in high latitude ocean regime.

Skill of stochastic models in two-layer QG

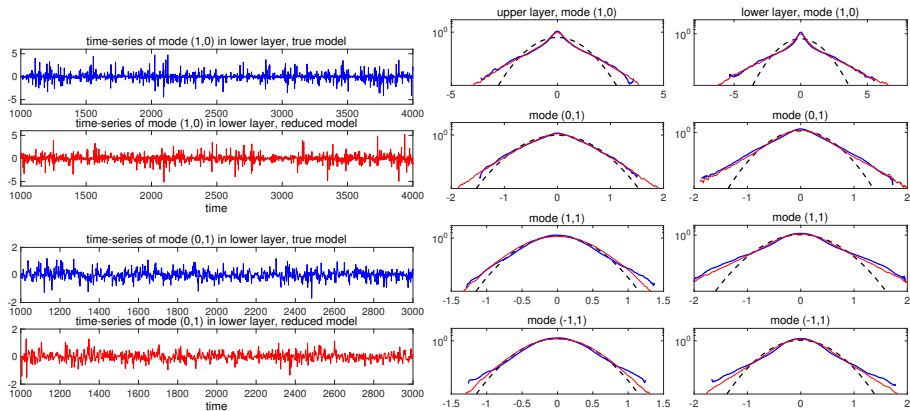


Figure: Prediction about tracer intermittency in high latitude atmosphere regime.

Skill of stochastic models in two-layer QG

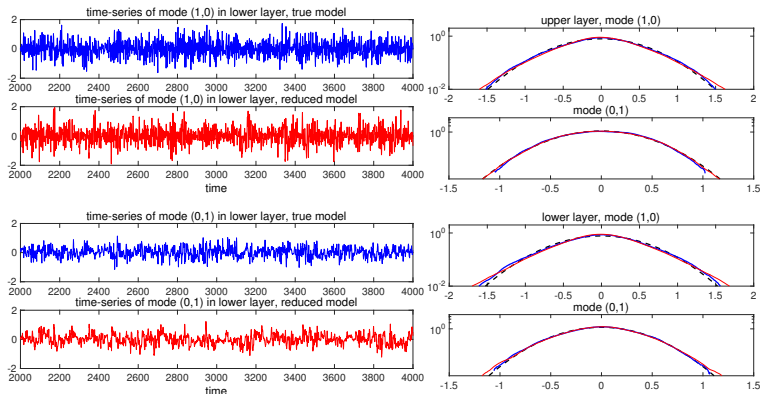


Figure: Prediction about tracer intermittency in mid latitude atmosphere regime with parameters.

Summary

- The single-point probability distribution function of a passive scalar has been the focus of much interest and requires better understanding.
 - The generation of tracer intermittency is related with the competition between the blocked and unblocked modes in advection flow.
 - The complex strongly turbulent dynamical system requires proper reduced-order modeling strategy of adopting simple advection flow models.
 - Reduced-order stochastic models in passive scalar turbulence are useful in predicting extreme events and intermittency.
-
- D. Qi and A. Majda, *Predicting fat-tailed Intermittent probability distributions in passive scalar turbulence with imperfect models through empirical information theory*, CMS, 2015.
 - D. Qi and A. Majda, *Predicting extreme events for passive scalar turbulence in two-layer baroclinic flows through reduced-order stochastic models*, submitted, 2017.

Thank you for your attention!