

Reverse Engineering a PDE from an Image Inpainting Algorithm

Rob Hocking

University of Cambridge

Supervisor: Carola-Bibiane Schönlieb

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Outline of Talk

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 - Yes. In fact analysis of (a) continuum limit already exists.
 - Actually, there is a second possible continuum limit...
 - ...which is “closer” to the discrete solution.
 - analyzing the *closer* limit facilitates the design of better algorithms.

The Problem

Image Inpainting

Filling holes / removing objects from images.

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Figure

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Approach of Interest

Fill in Shells using by taking weighted averages

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```
while  $O \neq \emptyset$ 
  for pixels  $x \in \partial O$ 
    
$$u(x) \leftarrow \frac{\sum_{y \in B_\epsilon(x) \cap O^c} w(x, y) u(y)}{\sum_{y \in B_\epsilon(x) \cap O^c} w(x, y)}$$

  end
   $O \leftarrow O \setminus \partial O.$ 
end
```

Example: Uniform Weights



Figure

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Figure

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- This limit does not account for other behaviour we will see later.
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- It does.

Motivation

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- Yes - in fact, in the movie industry entire films are routinely converted.
- There are companies that exist solely for this purpose.
- One of them is Gener8, a company I worked for at the start of my PhD.



Figure

“Recent” Converted Films



(a)



(b)

Render the Scene from a new viewpoint

A depth map is used to decide how much to shift pixels left or right.

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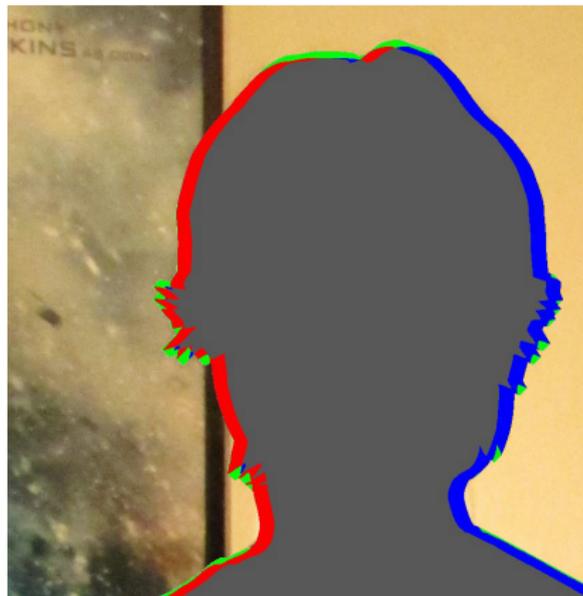
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Figure

Incomplete Disocclusion

In classical inpainting, our goal is to fill the hole in an image *in its entirety*. But in the present case, we only have to fill *part* of the hole.



Figure

Classical Inpainting vs. Present Case

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- Shell based approaches are natural, as they can be “stopped early”.
- GPU implementation is extremely fast.

ErodeFill and Compass Fill

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ErodeFill and Compass Fill

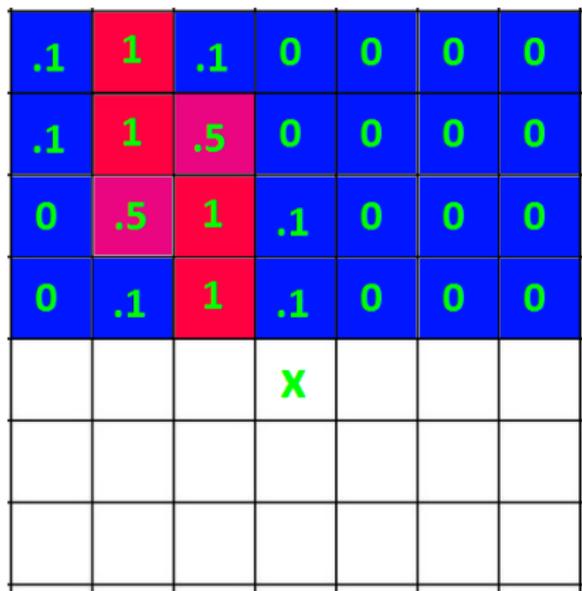
- When I arrived at Gener8, the company was using two inpainting schemes, both special cases of the approach described earlier.
- Both suffered from the “kinking” phenomena we saw earlier.
- My Job: “Can you make that kinking go away?”

Coherence Transport

Idea: Instead of uniform weights, assign higher weights to pixels on edges.

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Coherence Transport

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For each pixel x due to be filled, estimate local edge magnitude $\mu(x) \geq 0$ and direction $\mathbf{g}(x) \in S^1$.

Adapt weights accordingly:

$$w(x, y) = \frac{\exp\left(-\frac{\mu(x)^2}{2\epsilon^2} (\mathbf{g}^\perp(x) \cdot (y - x))^2\right)}{\|x - y\|}.$$

Coherence Transport

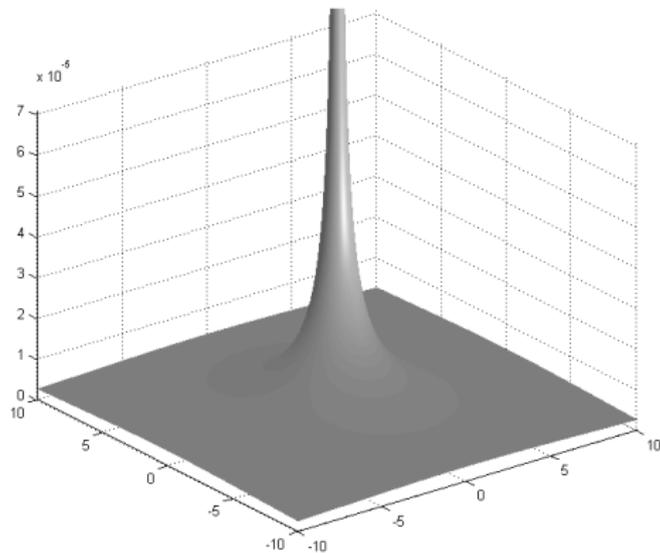


Figure: $\mu = 0$

Coherence Transport

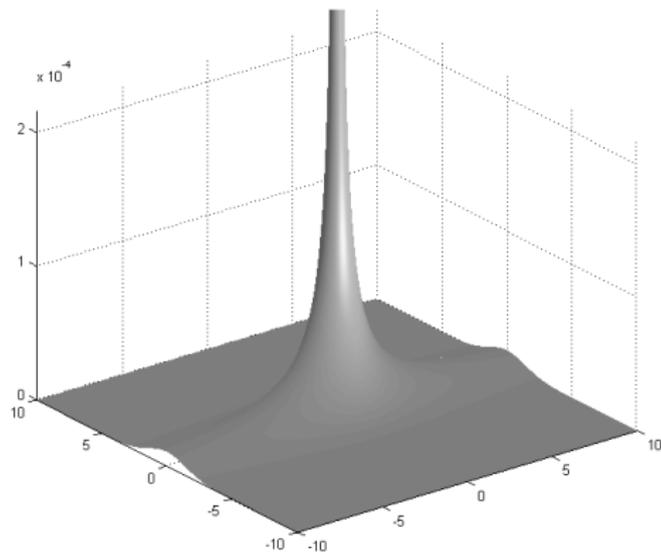


Figure: $\mu = 1$

Coherence Transport

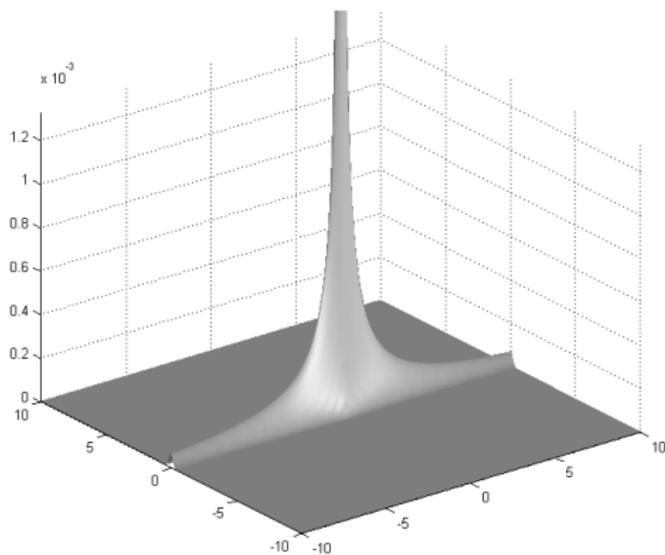
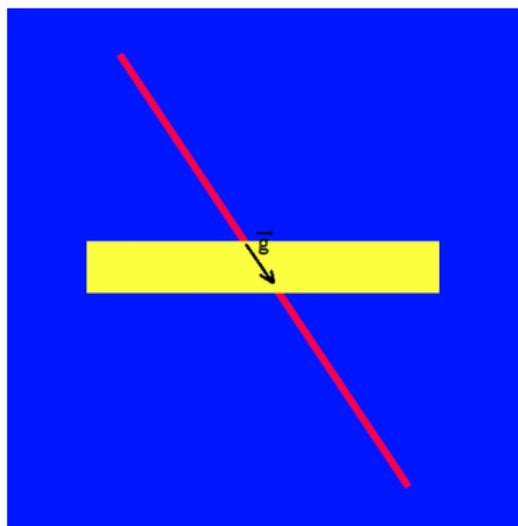


Figure: $\mu = 10$

Is the kinking fixed?

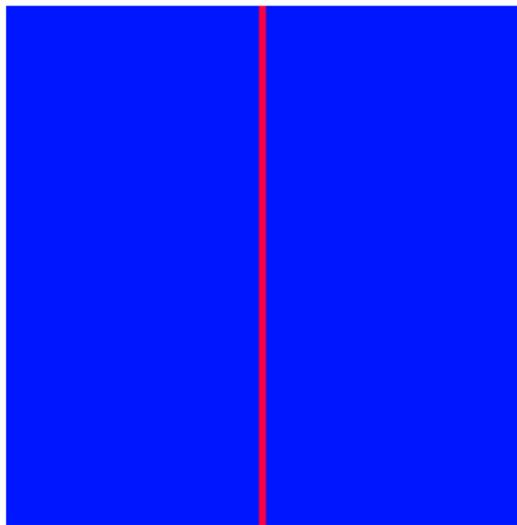
Test the method on a toy problem, feeding in the correct g by hand and set $\mu \gg 1$.



Figure

Actually, no

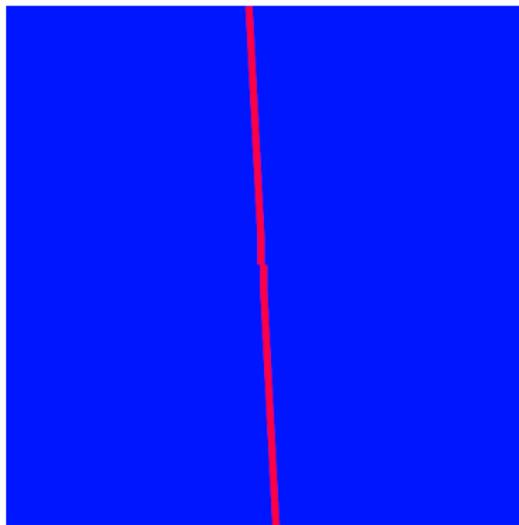
$$\epsilon = 3\sigma, \mu = 50.$$



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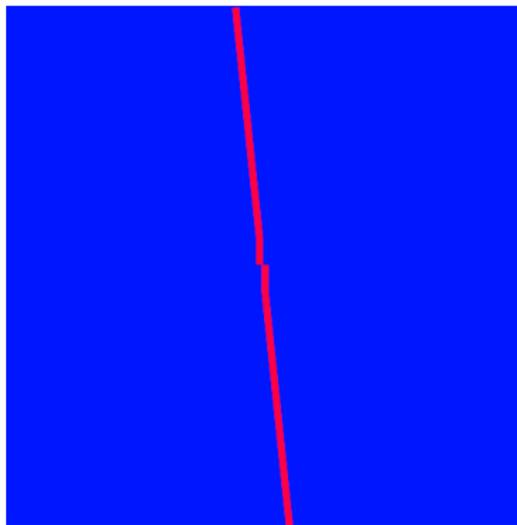
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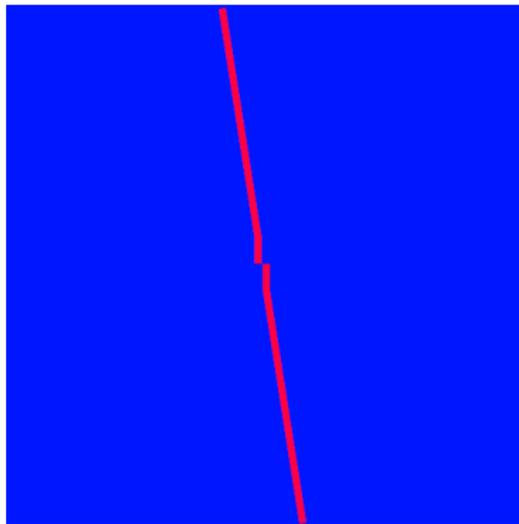
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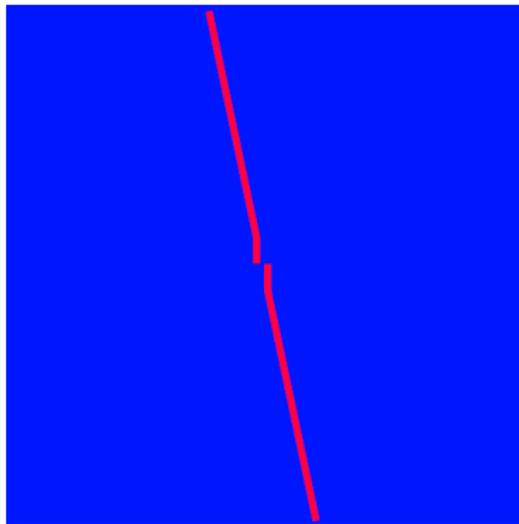
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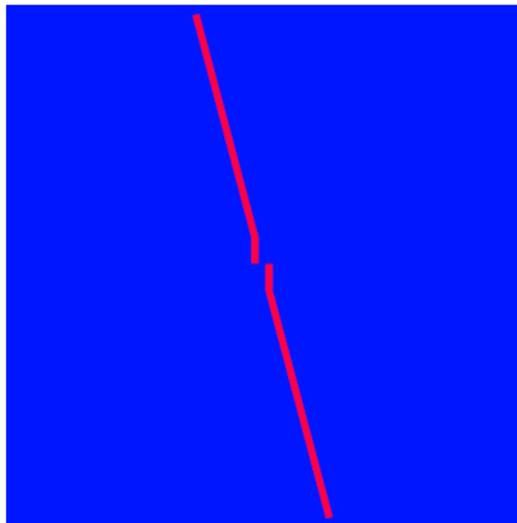
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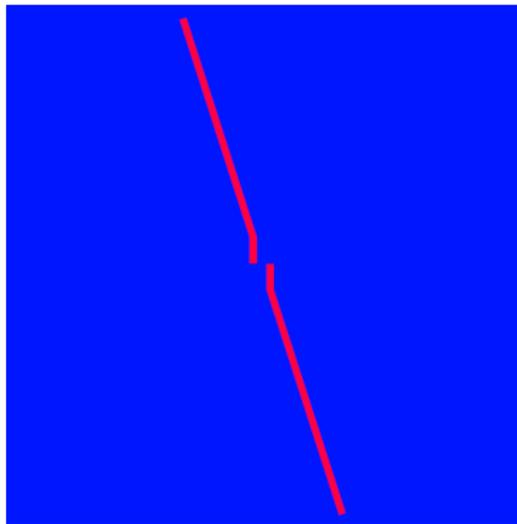
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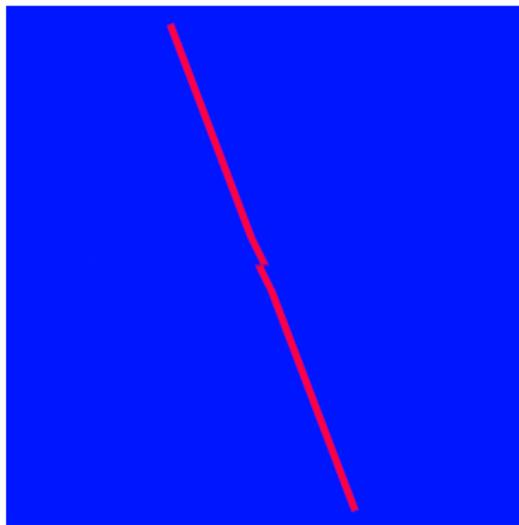
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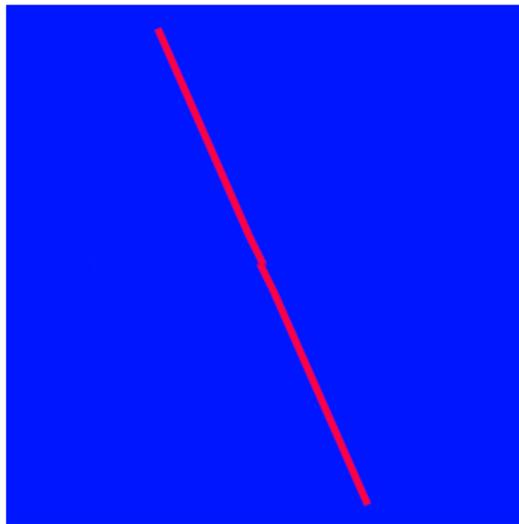
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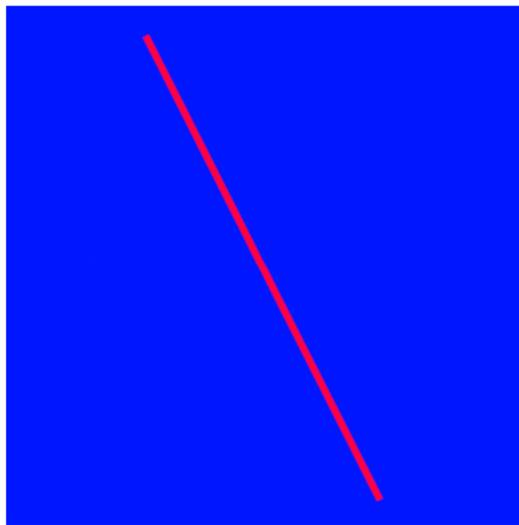
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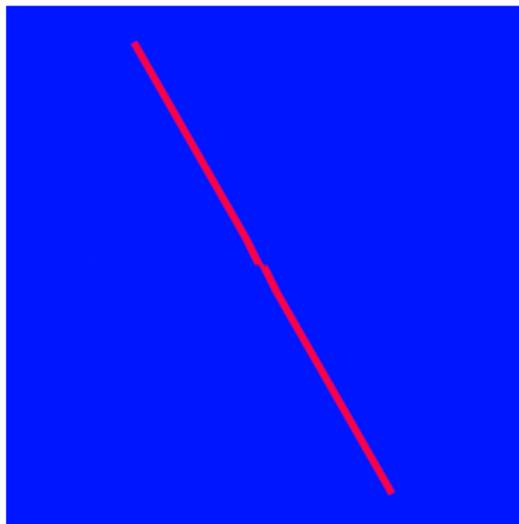
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Intuitive Fix

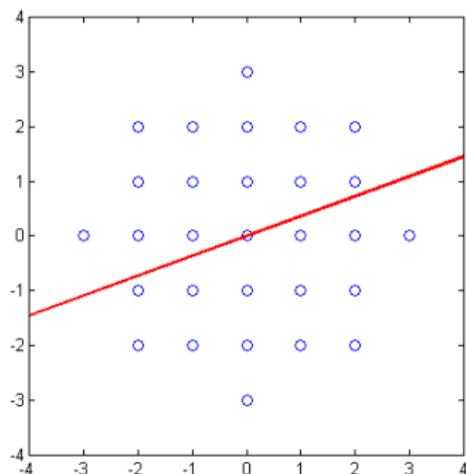
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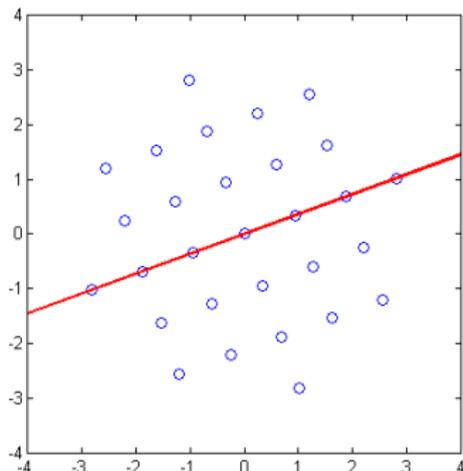
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Sum over a rotated ball of “ghost pixels” instead.

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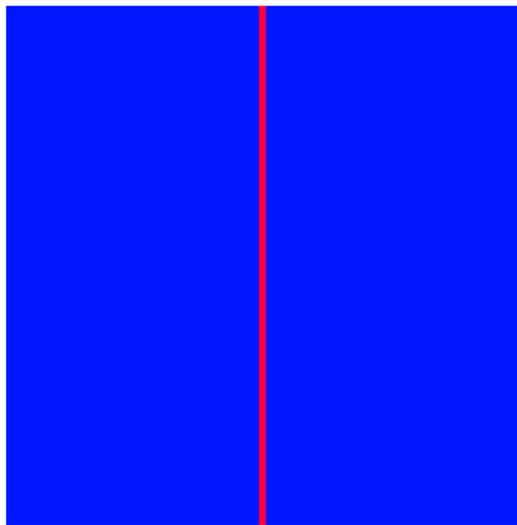


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Define “ghost pixels” using bilinear interpolation.

Lines do not bend

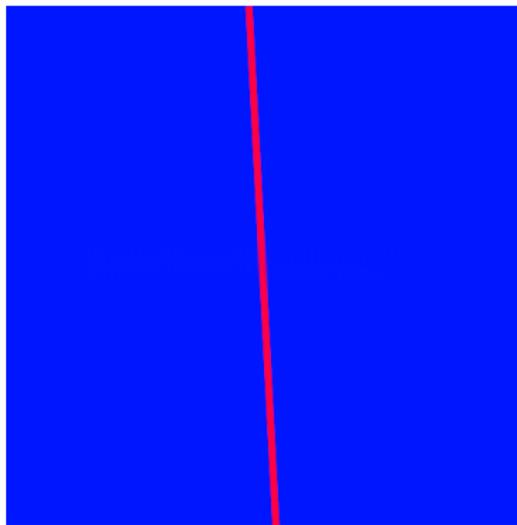
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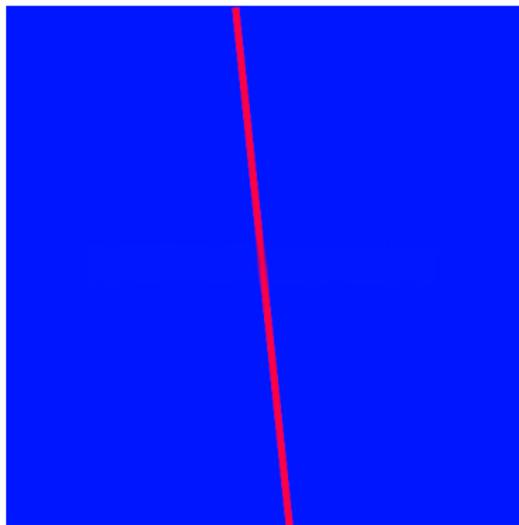
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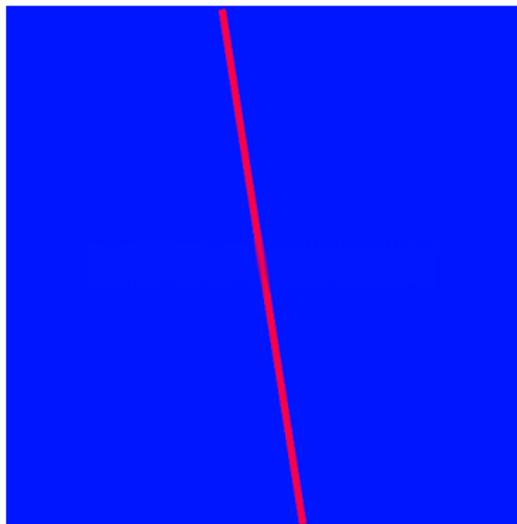
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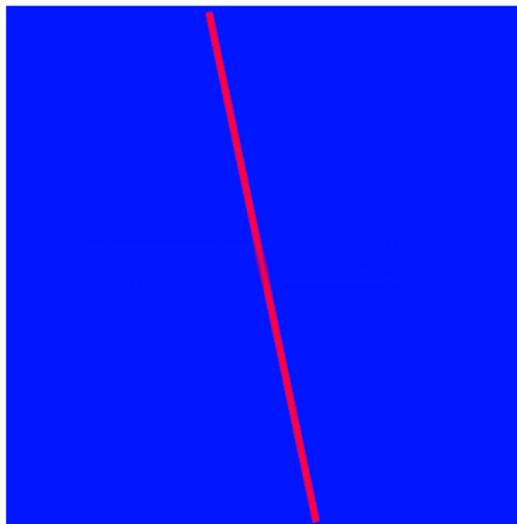
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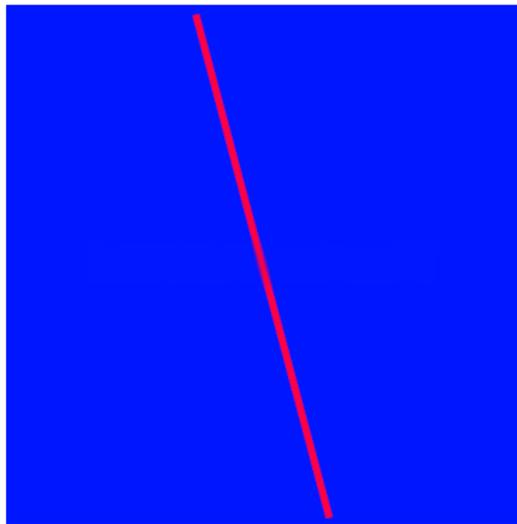
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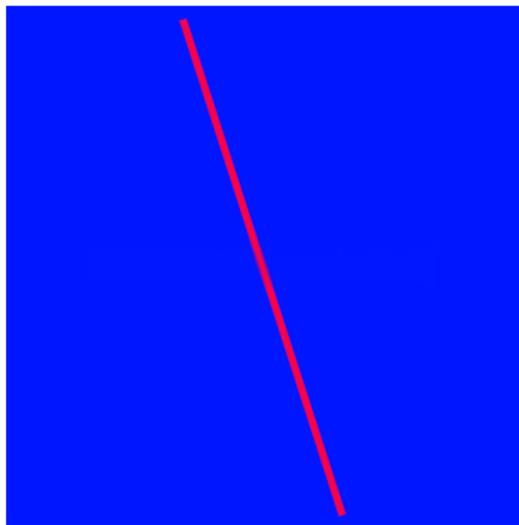
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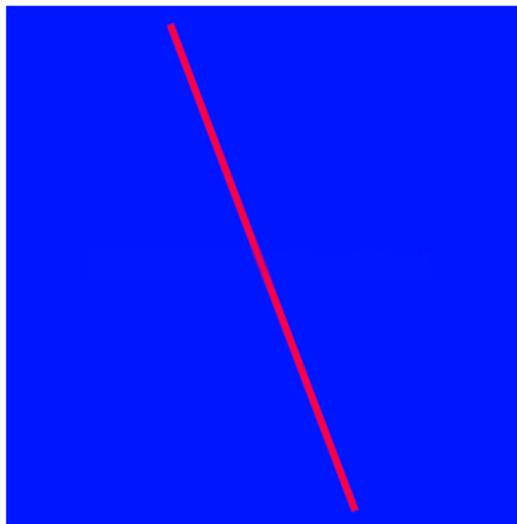
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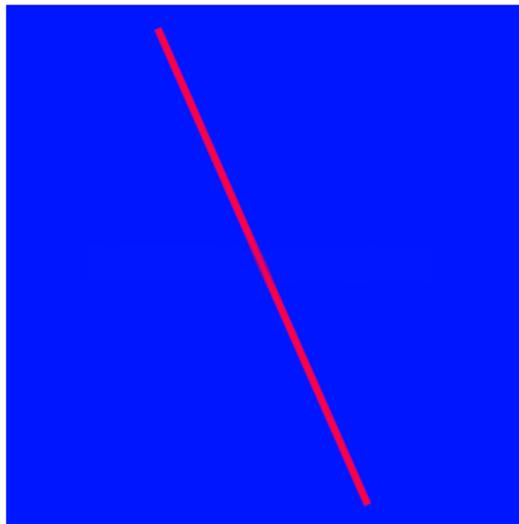
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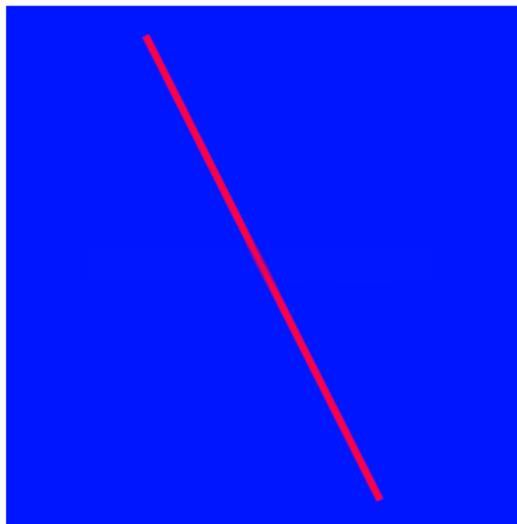
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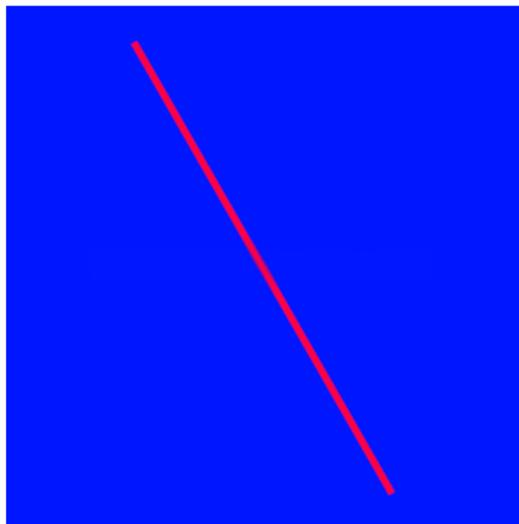
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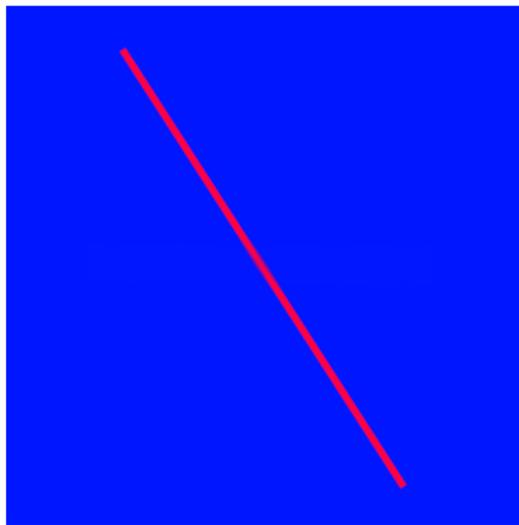
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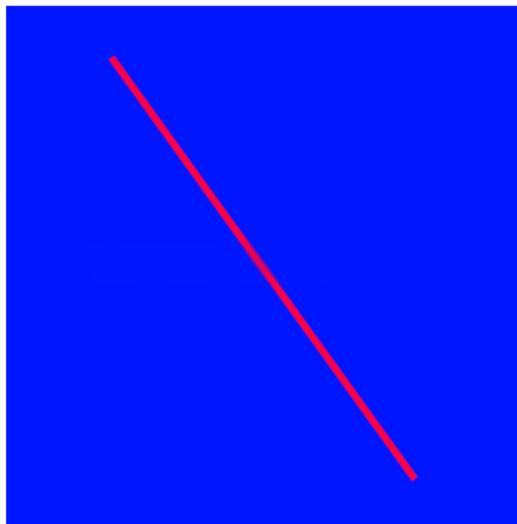
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Figure

Maths

What's going on?

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- Why did ghost pixels fix it?

Assumptions

From now on, assuming inpainting domain $O = (0, 1]^2$, with periodic boundary conditions at $x = 0$ and $x = 1$.

Boundary data is on the strip $(0, 1] \times (-\delta, 0]$.

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Also some restrictions on the weights.

März justifies Coherence Transport by arguing that in the *double limit* $h \rightarrow 0$ and then $\epsilon \rightarrow 0$, the method behaves like the transport PDE

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where,

- $\mathbf{g}_\mu^* := \int_{y \in B_1^-} w_{\mu,1}(0, y) y dy \rightarrow \mathbf{g}$ as $\mu \rightarrow \infty$,
- $B_1^- := \{(y_1, y_2) \in \mathbb{R}^2 : y_1^2 + y_2^2 \leq 1 \text{ and } y_2 \leq 0\}$,

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Call the continuum solution $u_{\text{m\u00e4rz}}$, discrete solution u_h .

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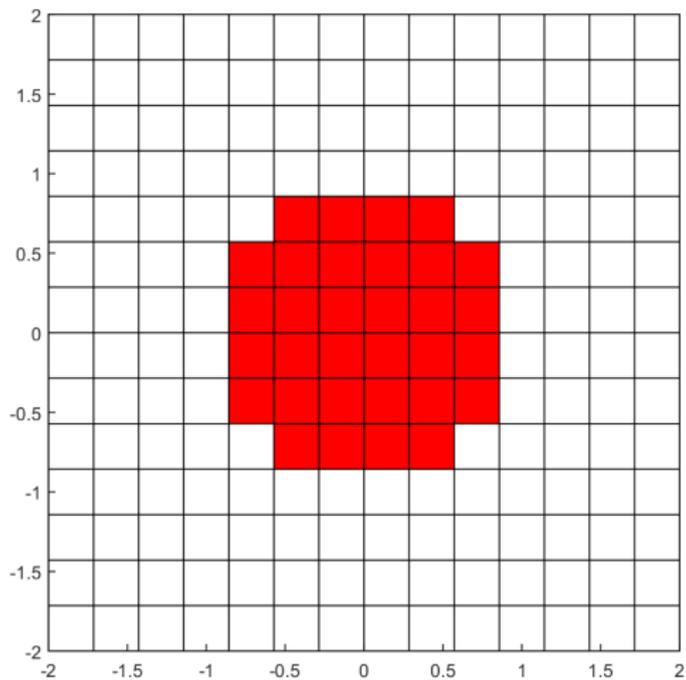
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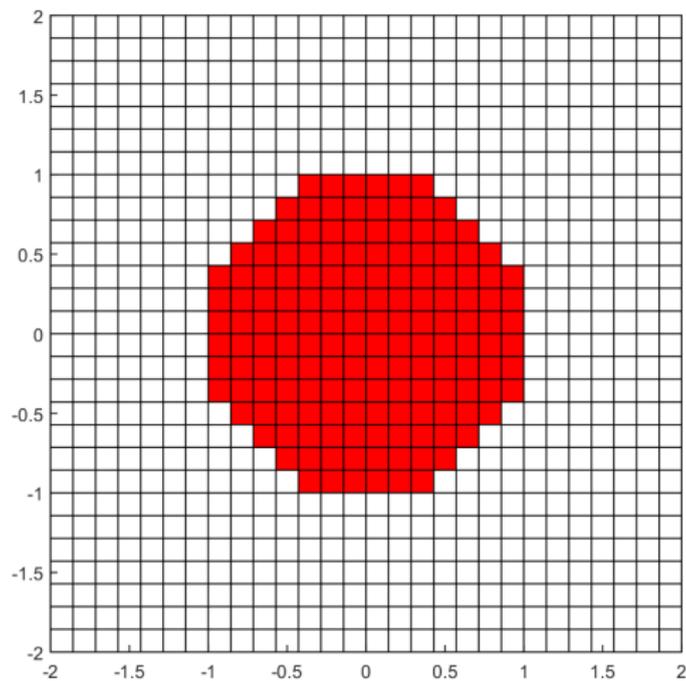
First, let's take $h \rightarrow 0$...



Figure

Closer Look

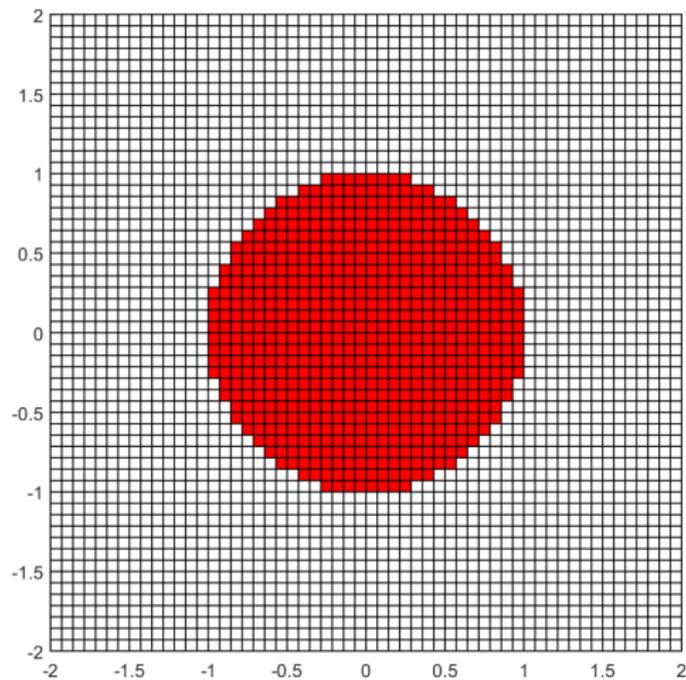
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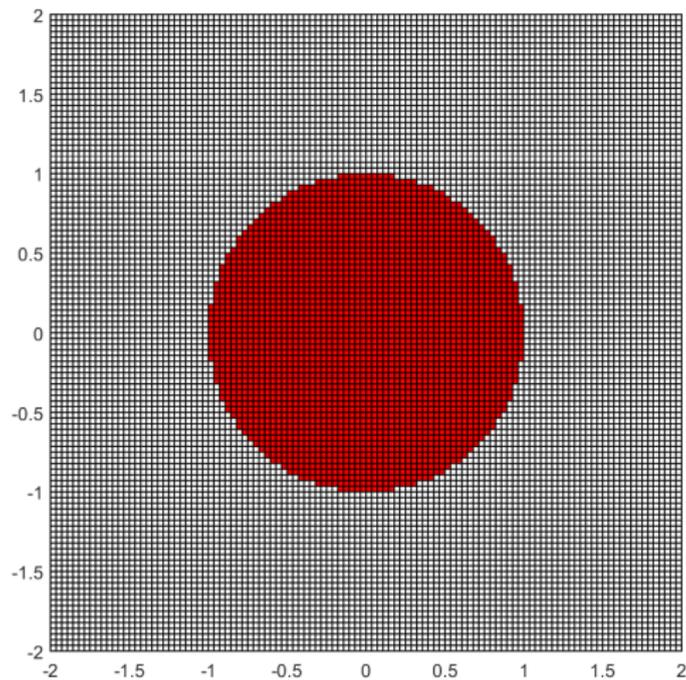
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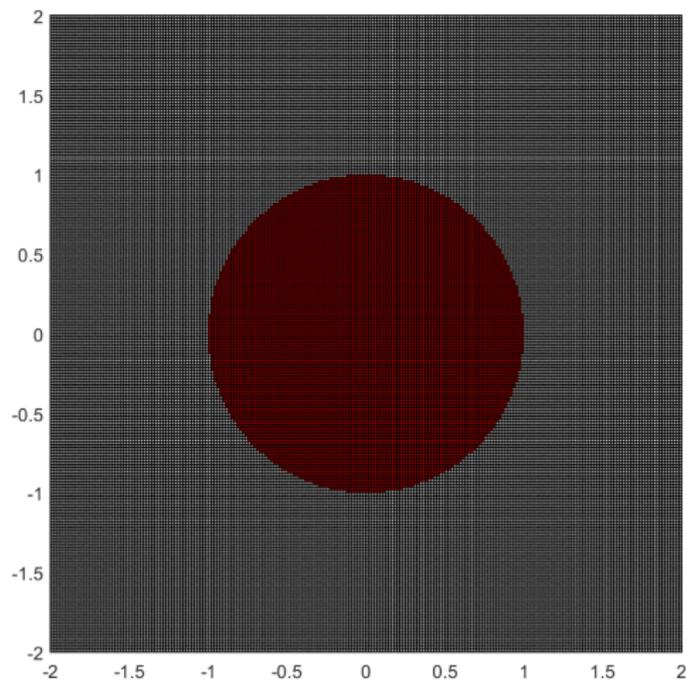
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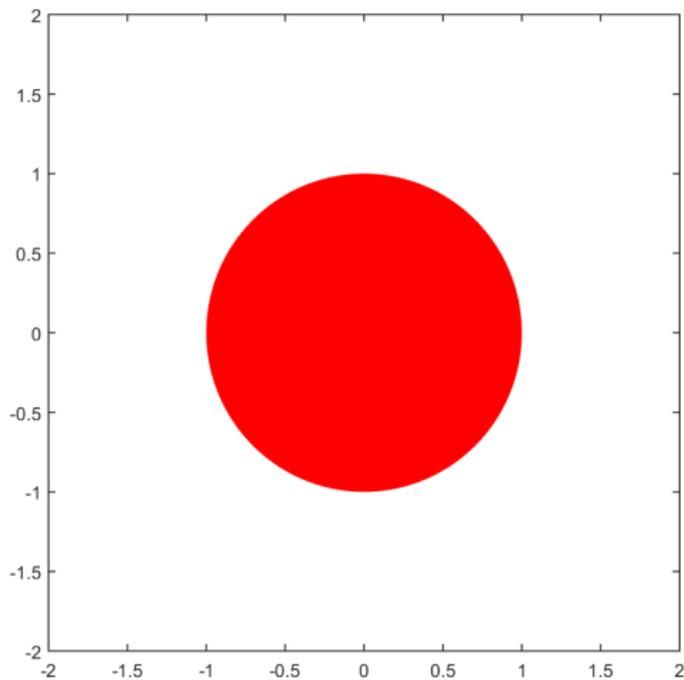
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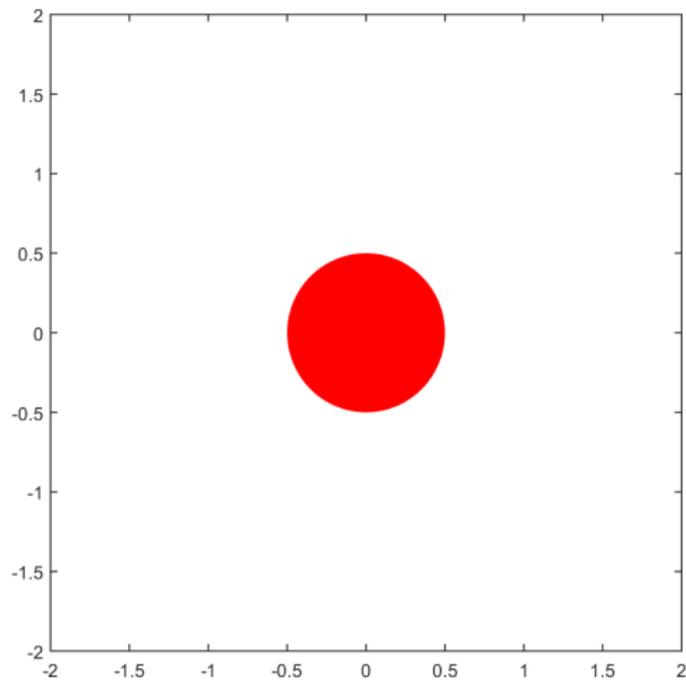
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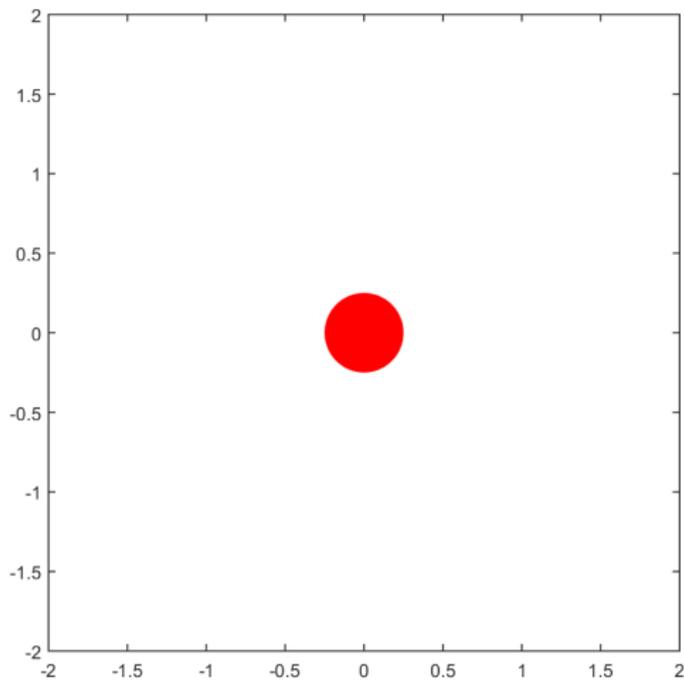
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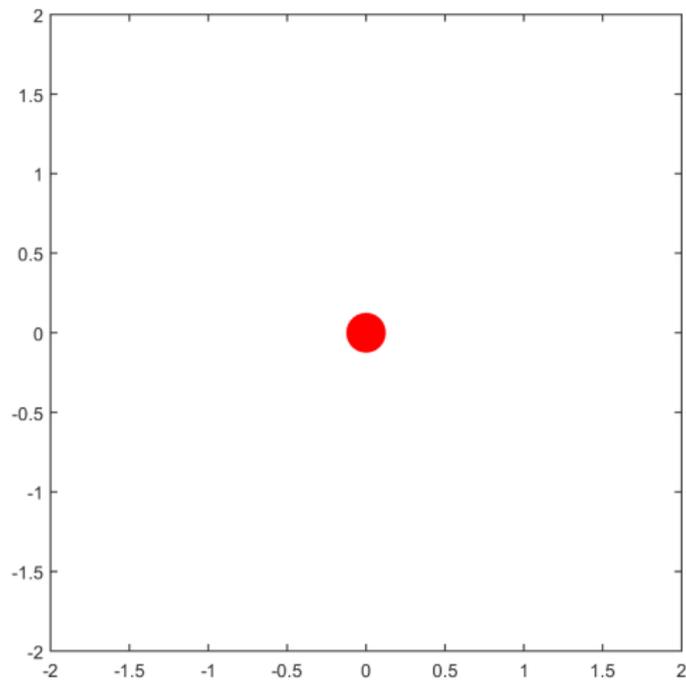
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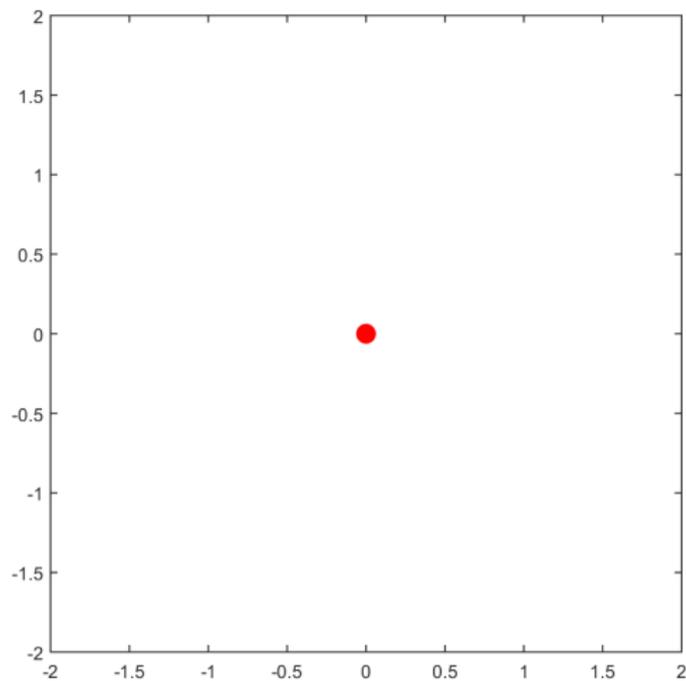
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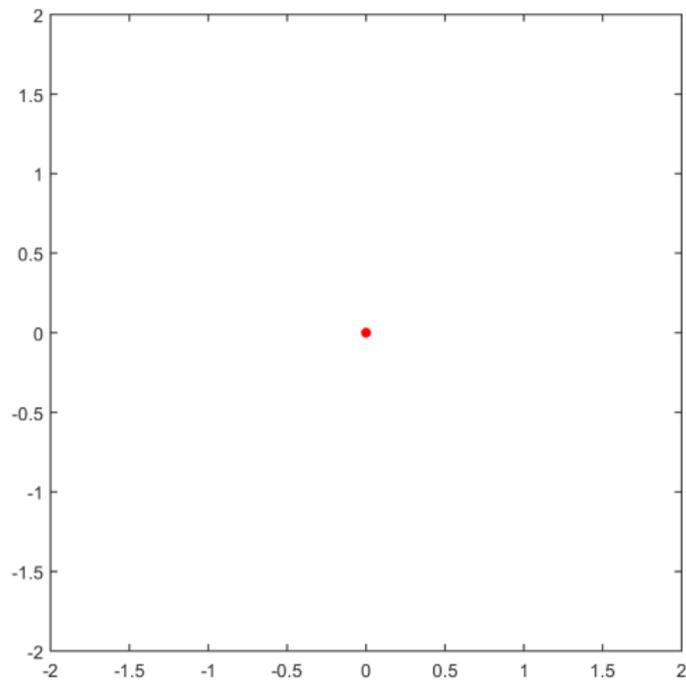
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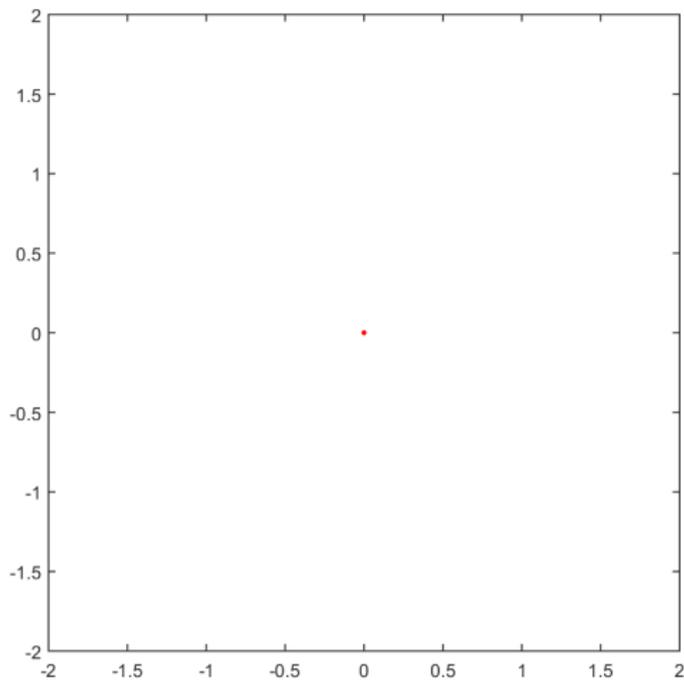
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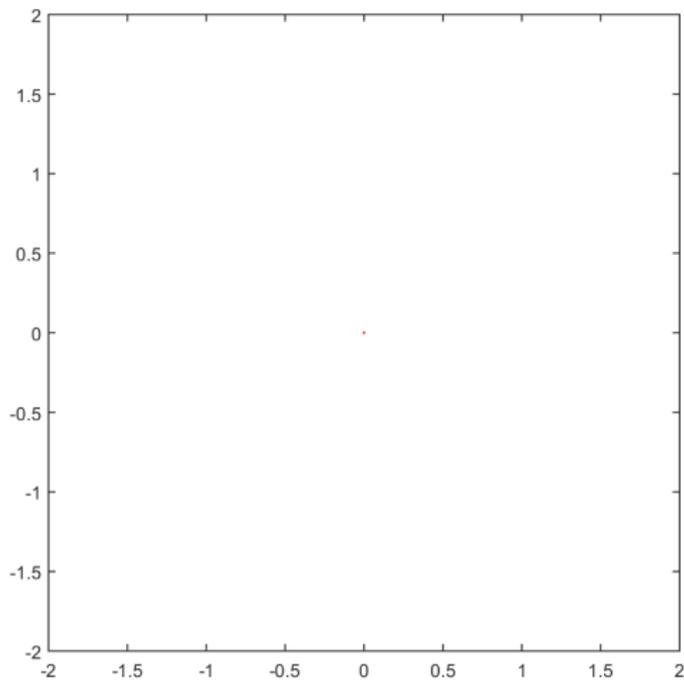
Next, take $\epsilon \rightarrow 0\dots$



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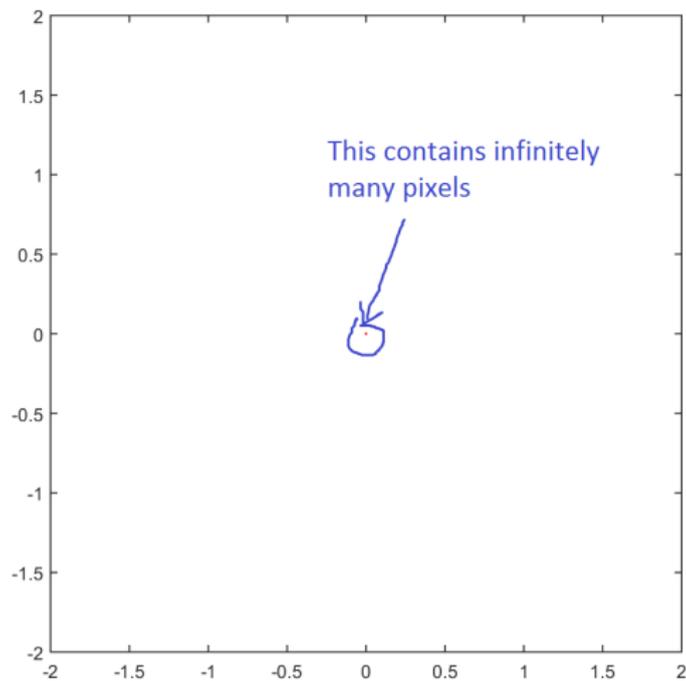
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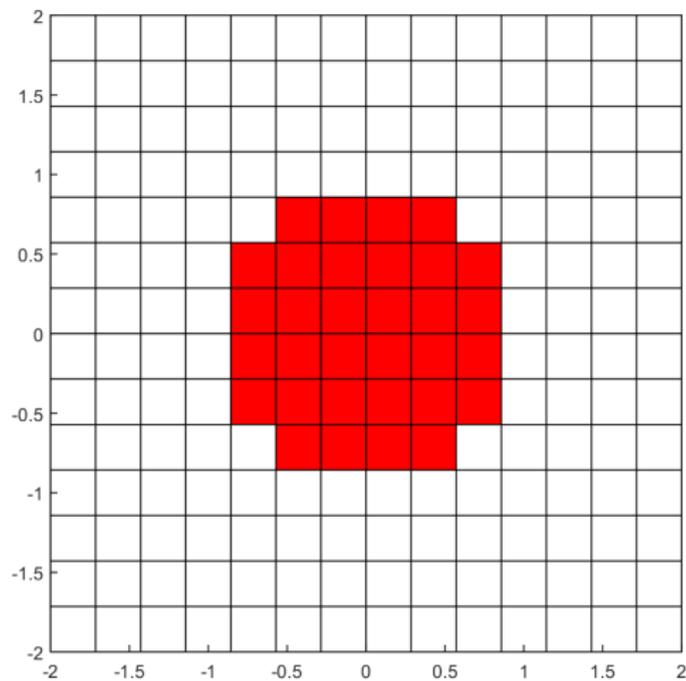
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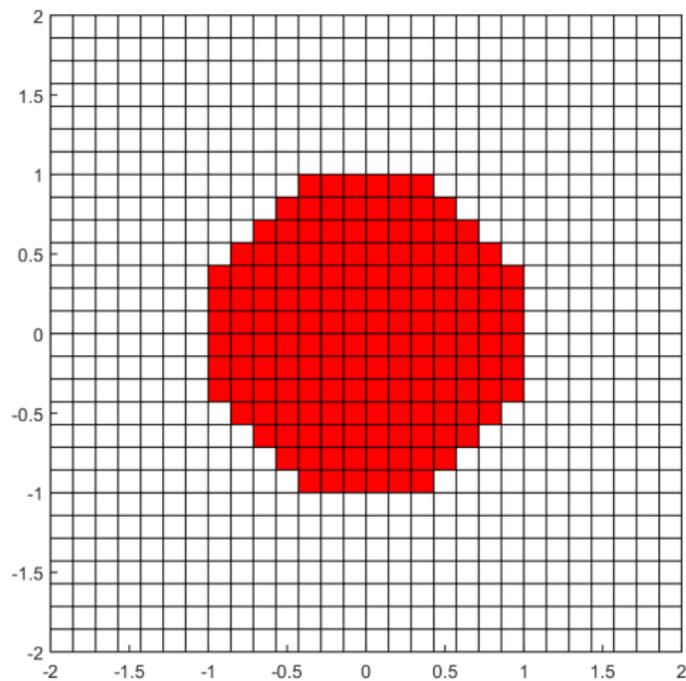
First, let's make h "pretty small".



Figure

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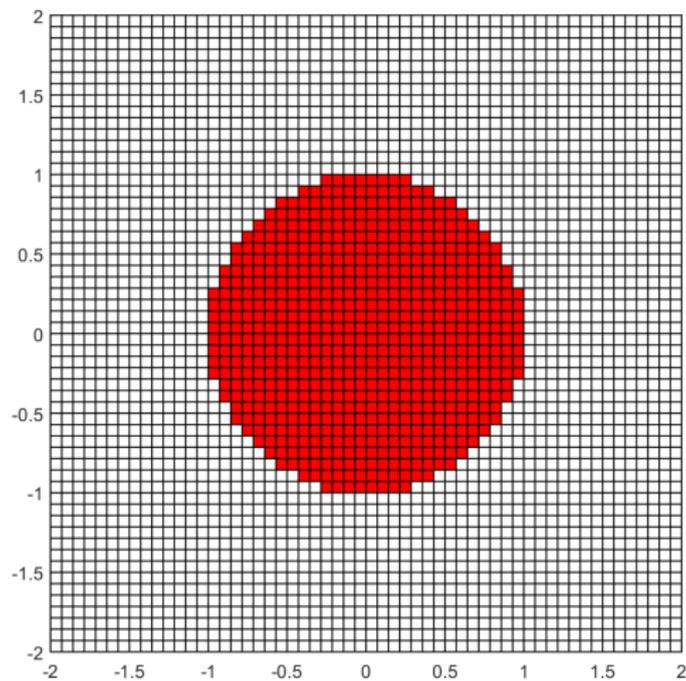
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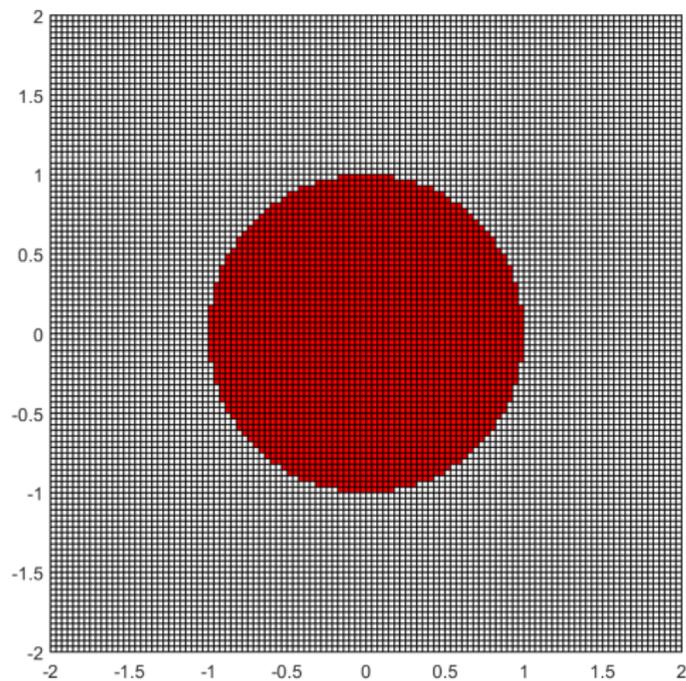
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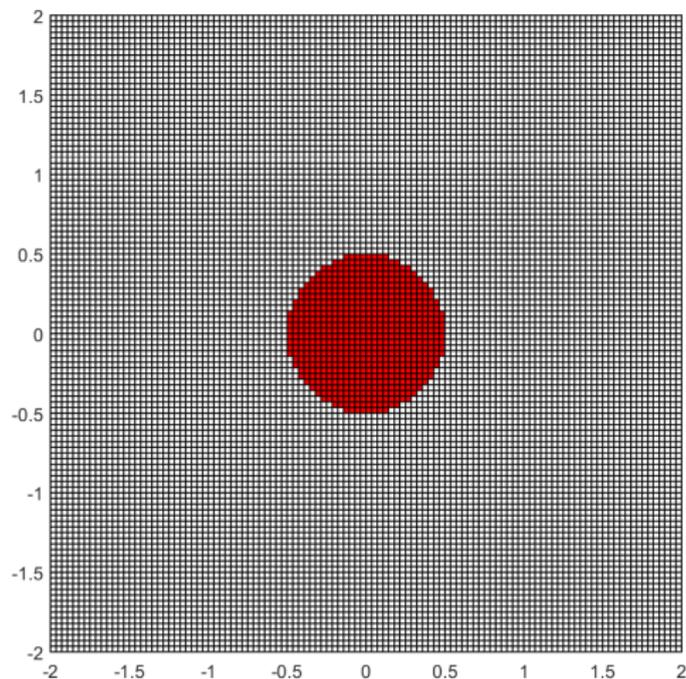
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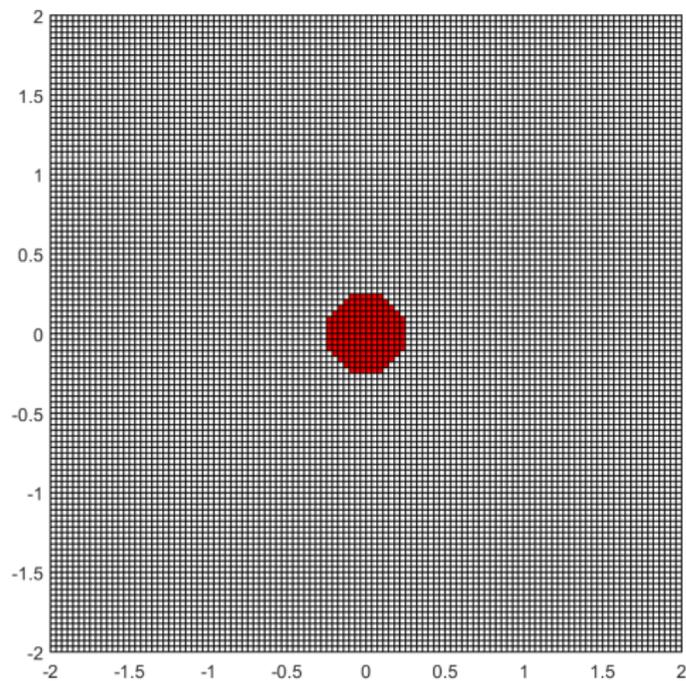
Now let's make ϵ "pretty small" as well.



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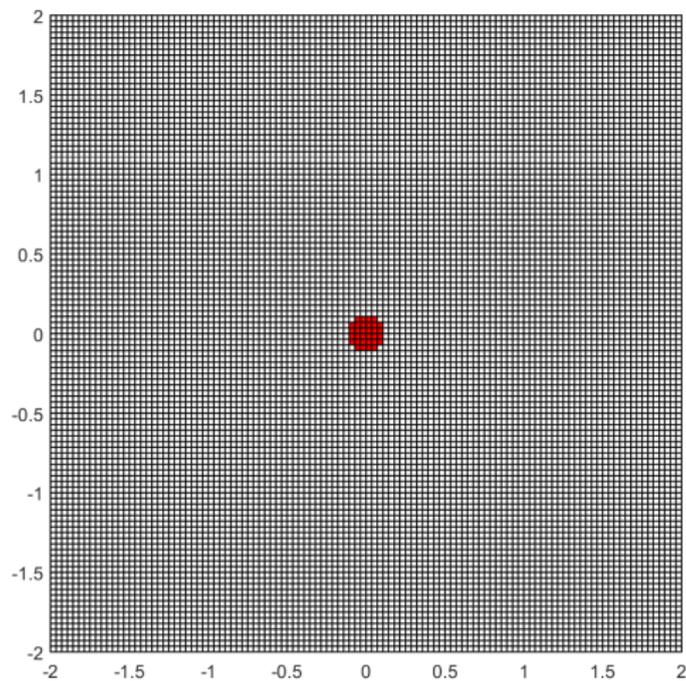
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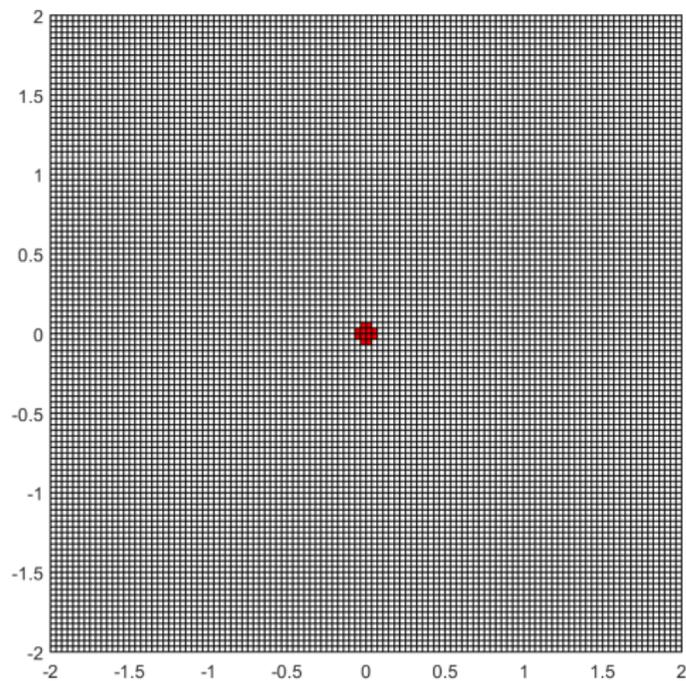
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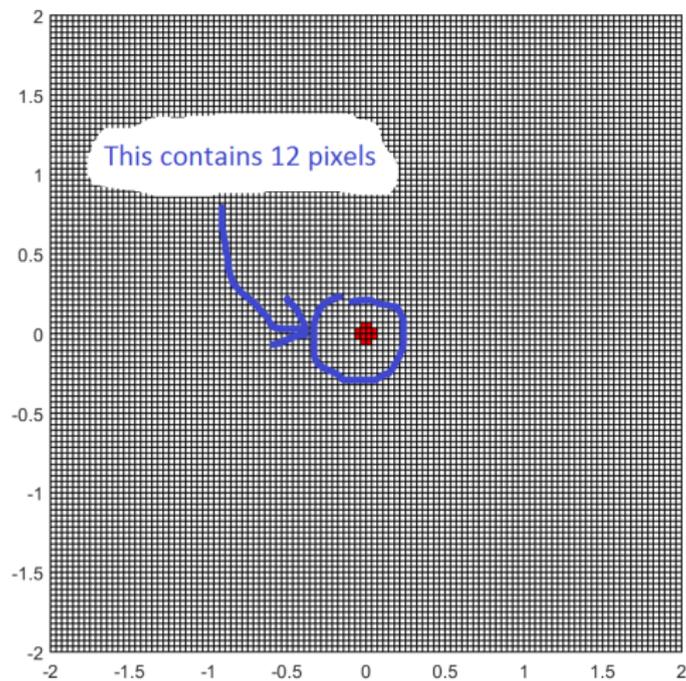
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Figure

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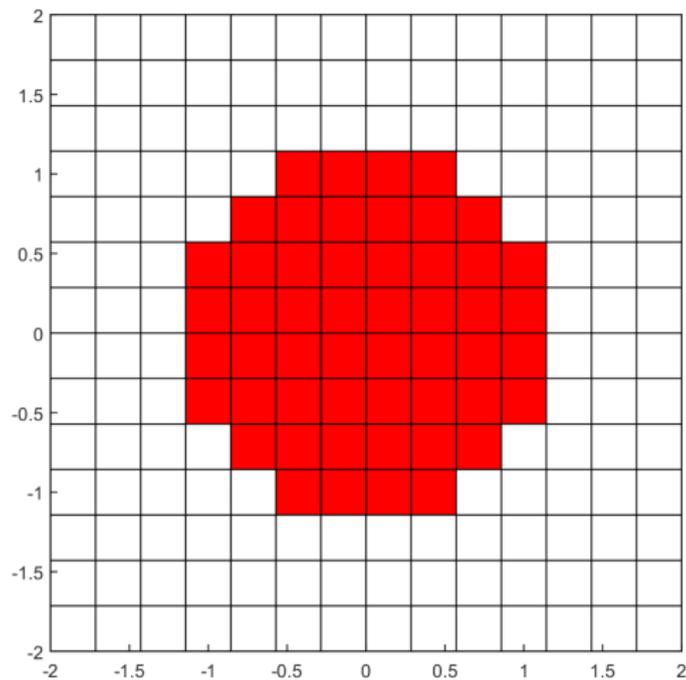
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Makes more sense to study the limit $h \rightarrow 0$ with r fixed.

New Limit

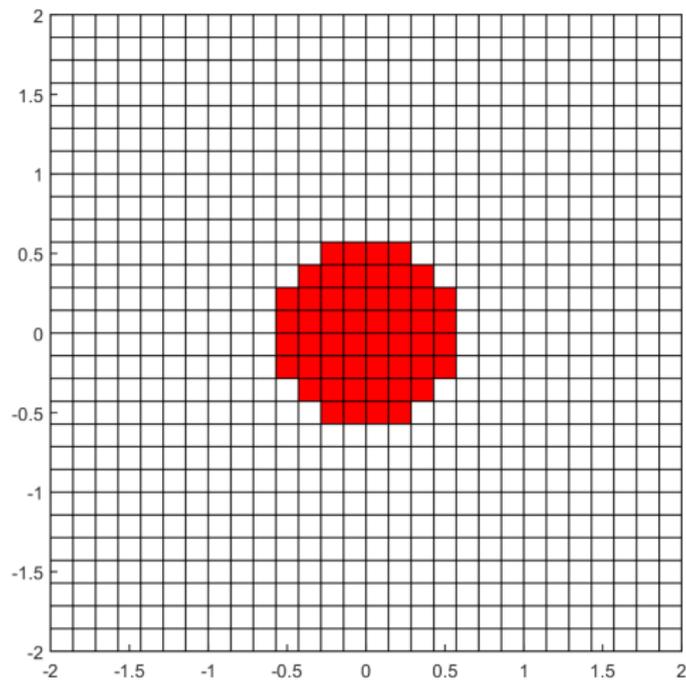
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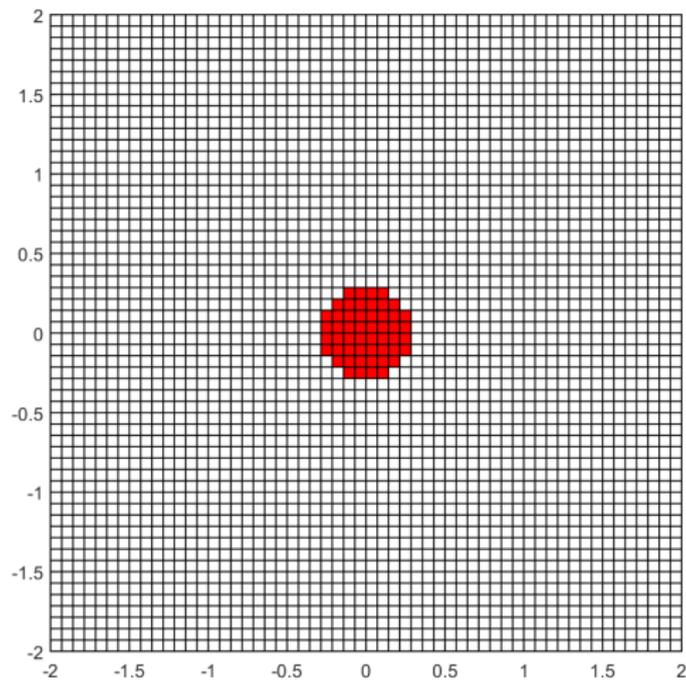
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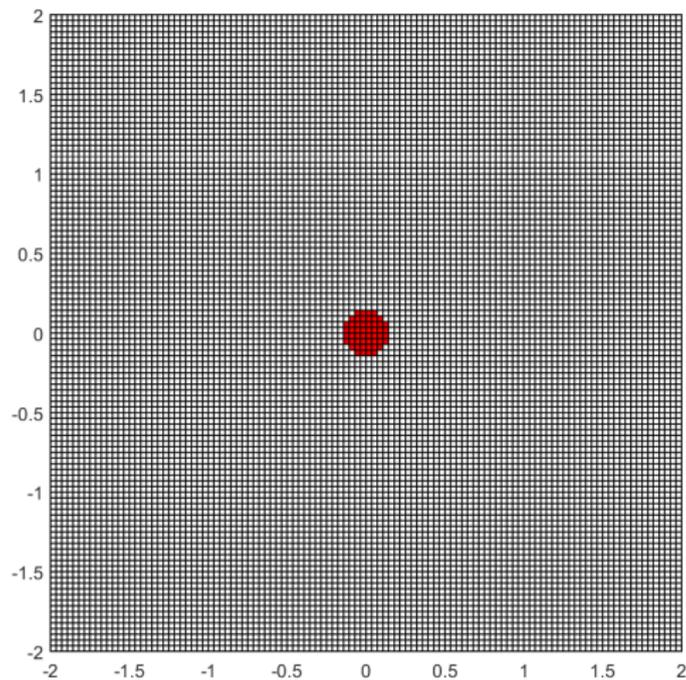
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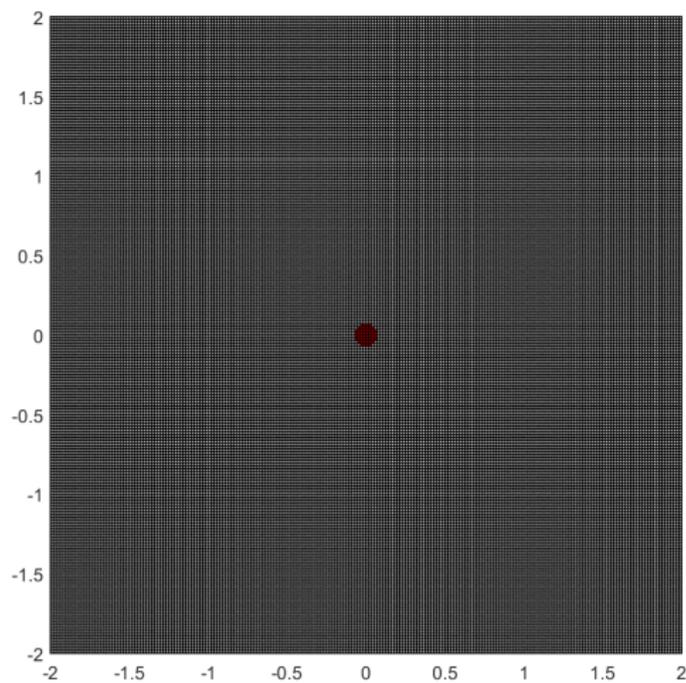
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- true for all $1 \leq p \leq \infty$ if there are no jumps.

Which Limit is Closer?

But for *fixed* h , one finds:

$$\|u_h - u_r\|_p \lesssim C_r \|u_h - u_{\text{märz}}\|_p^2.$$

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$$\mathbf{g}_{\mu}^* := \int_{y \in B_1^-} w_{\mu,1}(0, y) y dy$$

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Explanation of Clamping

In this case we have

$$\mathbf{g}_{r,\mu}^* := \sum_{\mathbf{j} \in b_r^-} e^{-\frac{\mu^2}{2r^2} (\mathbf{j} \cdot \mathbf{g}^\perp)^2} \frac{\mathbf{j}}{\|\mathbf{j}\|},$$

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Then, rescaling by $e^{\frac{\mu^2}{2r^2} (\mathbf{j}^* \cdot \mathbf{g}^\perp)^2}$ we have

$$\mathbf{g}_{r,\mu}^* = \frac{\mathbf{j}^*}{\|\mathbf{j}^*\|} + \sum_{\mathbf{j} \in b_r^- \setminus \{\mathbf{j}^*\}} e^{-\frac{\mu^2}{2r^2} \{(\mathbf{j} \cdot \mathbf{g}^\perp)^2 - (\mathbf{j}^* \cdot \mathbf{g}^\perp)^2\}} \frac{\mathbf{j}}{\|\mathbf{j}\|}$$

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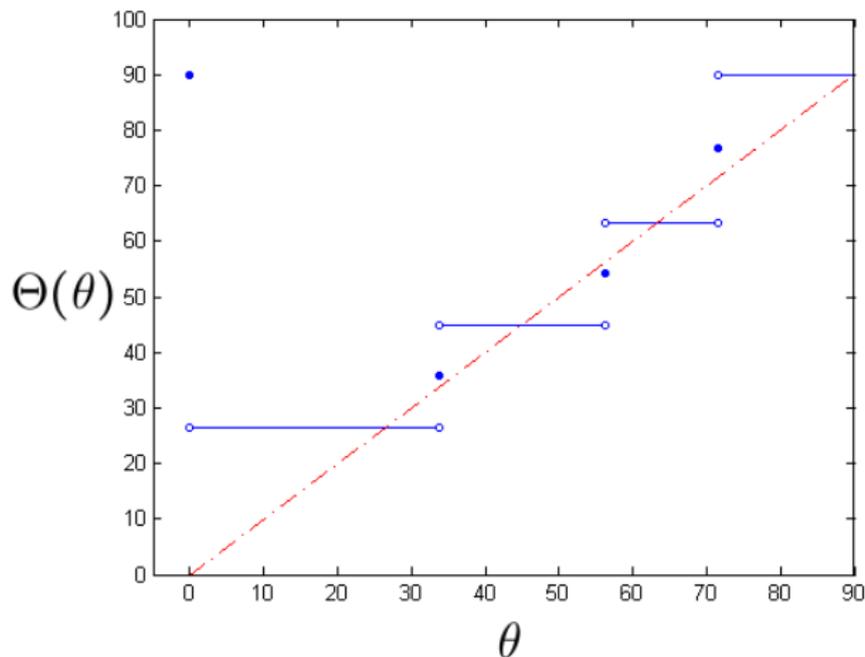
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Comparison with real Life

Define $\theta = \angle \mathbf{g}$, $\theta_r^* = \angle \mathbf{g}_r^*$, consider $\theta_r^* = \Theta(\theta)$.

Comparison with real Life

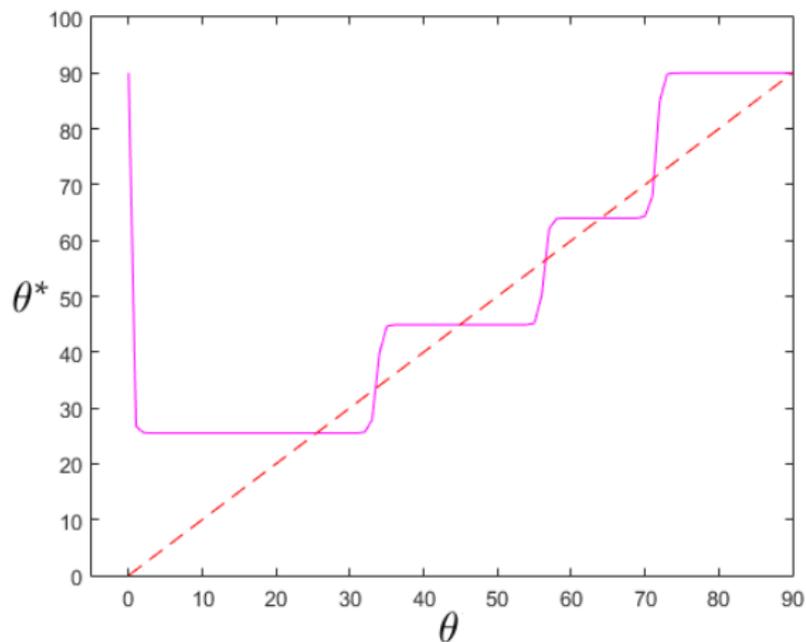
Theoretical Curve ($r = 3$)



Figure

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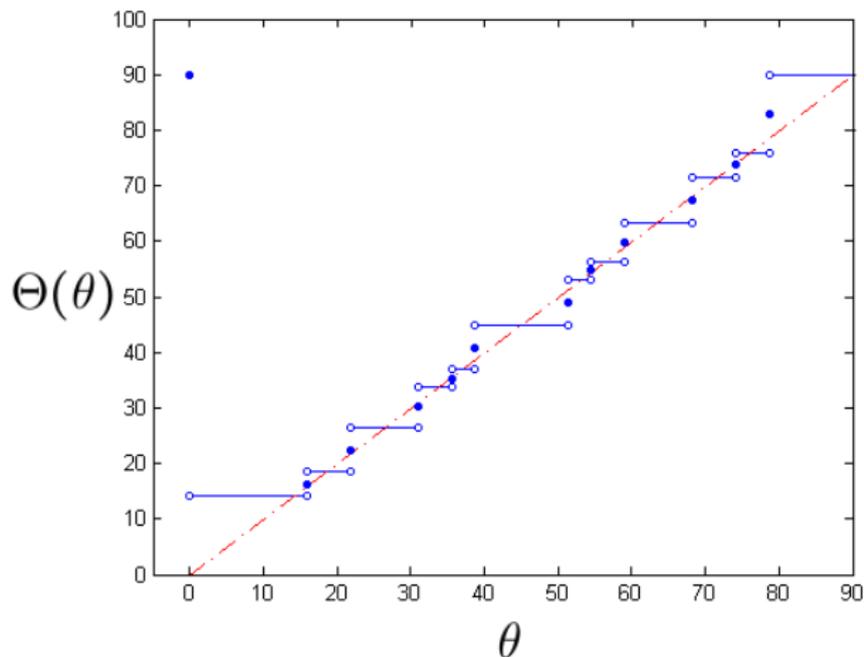
Real Curve ($r = 3, \mu = 40$)



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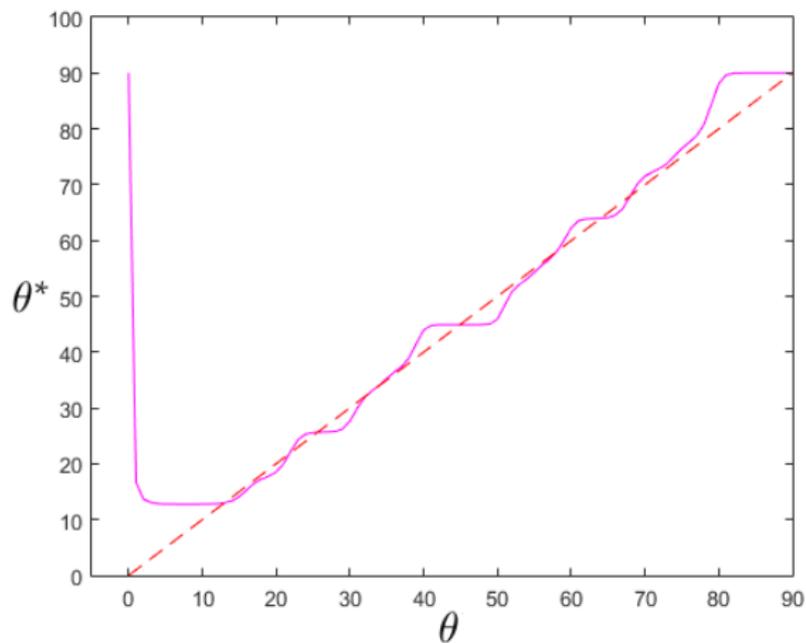
Theoretical Curve ($r = 5$)



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Comparison with real Life

Real Curve ($r = 5, \mu = 40$)



Figure

Explaining our Fix

- This explains earlier kinking.

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- But why were we able to make it go away?

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$$\mathbf{g}_r^* = \sum_{\mathbf{j} \in b_r^-} w(0, \mathbf{j}) \mathbf{j}.$$

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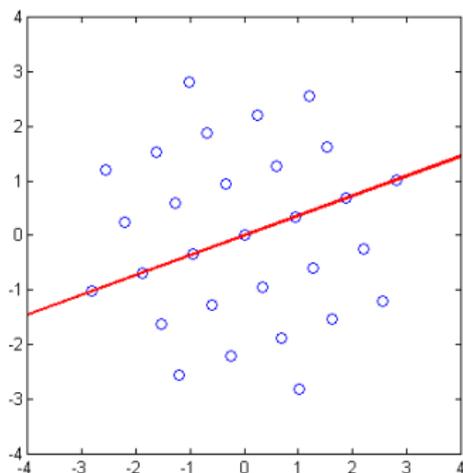
Earlier, we proposed a continuum limit u_r with transport direction

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But actually this assumed that we use no ghost pixels.

Back to the Fix

Now assume we sum over a rotated ball $\tilde{B}_{\epsilon,h}(\mathbf{x})$.



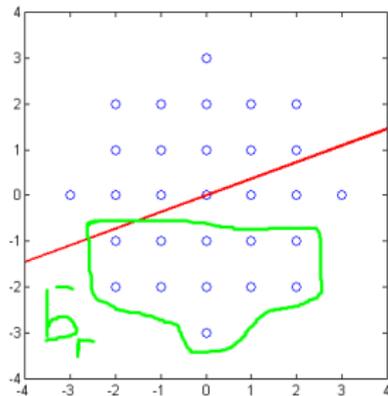
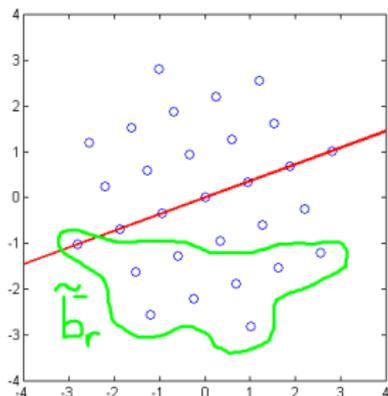
Figure

Back to the Fix

We get a remarkably similar formula:

$$\mathbf{g}_r^* = \sum_{\mathbf{j} \in \tilde{b}_r^-} w(0, \mathbf{j}) \mathbf{j} \quad \text{vs.}$$

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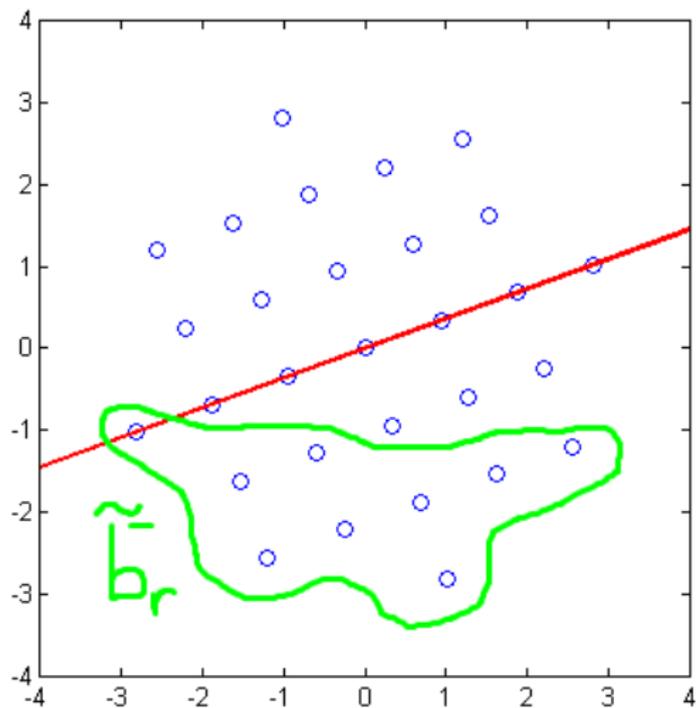
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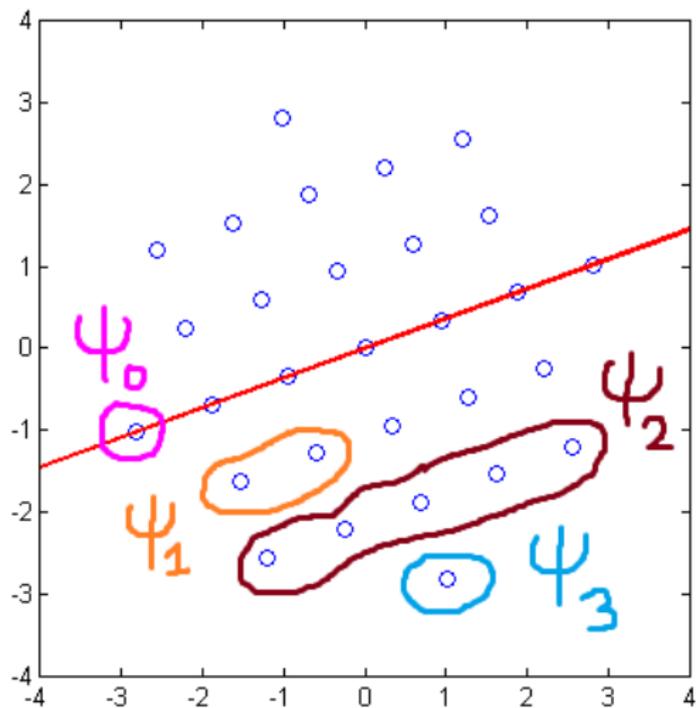
- This simple formula is a consequence of our choice to define ghost pixels via bilinear interpolation.
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- Now let's see why everything is fixed.

Back to the Fix



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$$\mathbf{g}_{r,\mu}^* = \sum_{\mathbf{j} \in \Psi_0} \frac{\mathbf{j}}{\|\mathbf{j}\|} + \sum_{k=1}^r \sum_{\mathbf{j} \in \Psi_k} e^{-\frac{\mu^2}{2r^2} k^2 \|\mathbf{g}\|^2} \frac{\mathbf{j}}{\|\mathbf{j}\|}$$

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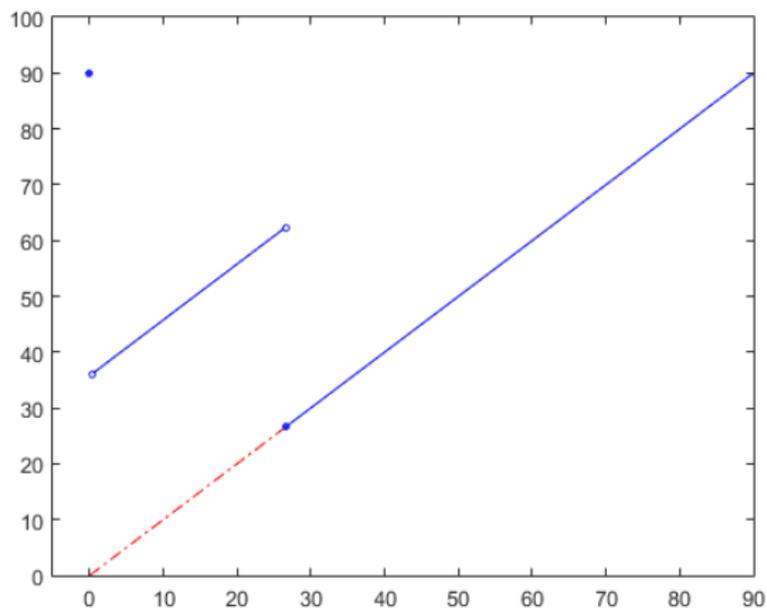
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Can show $\Psi_0 \neq \emptyset$ if $\theta > \theta_c(r)$.

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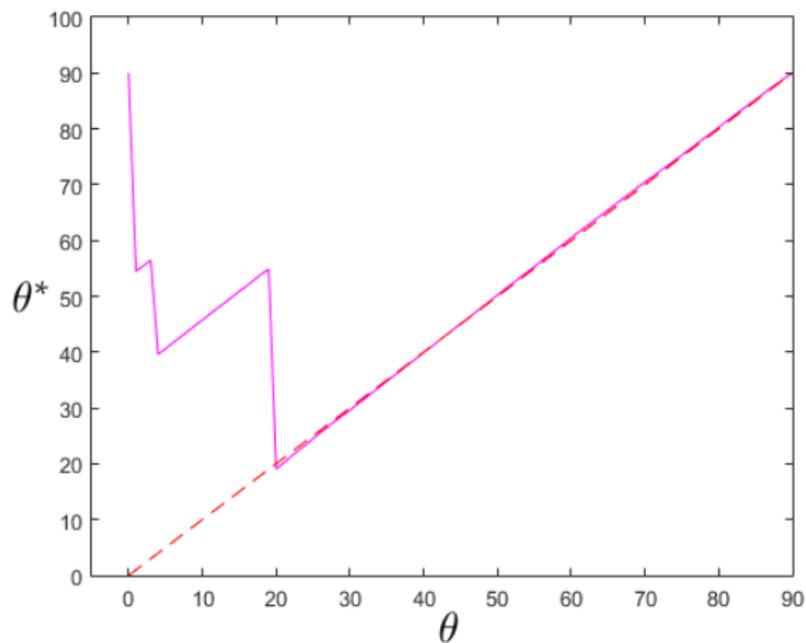
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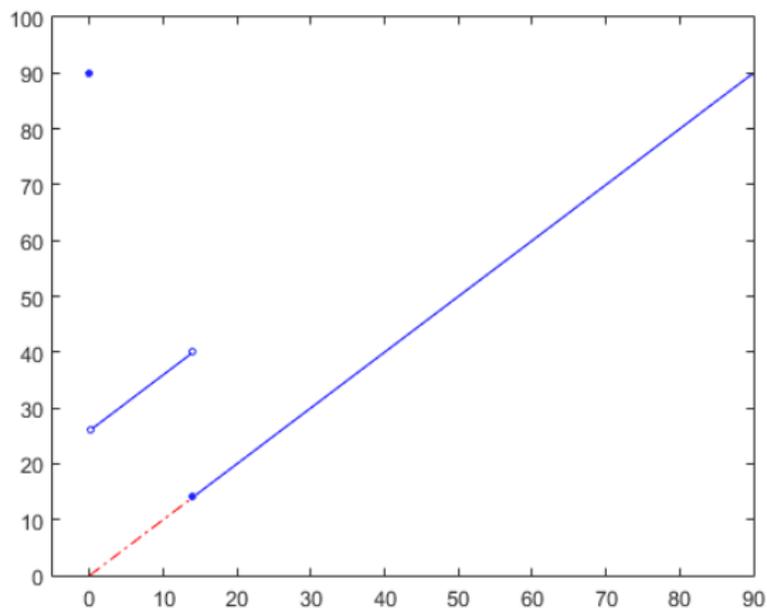
Real Curve ($r = 3, \mu = 40$)



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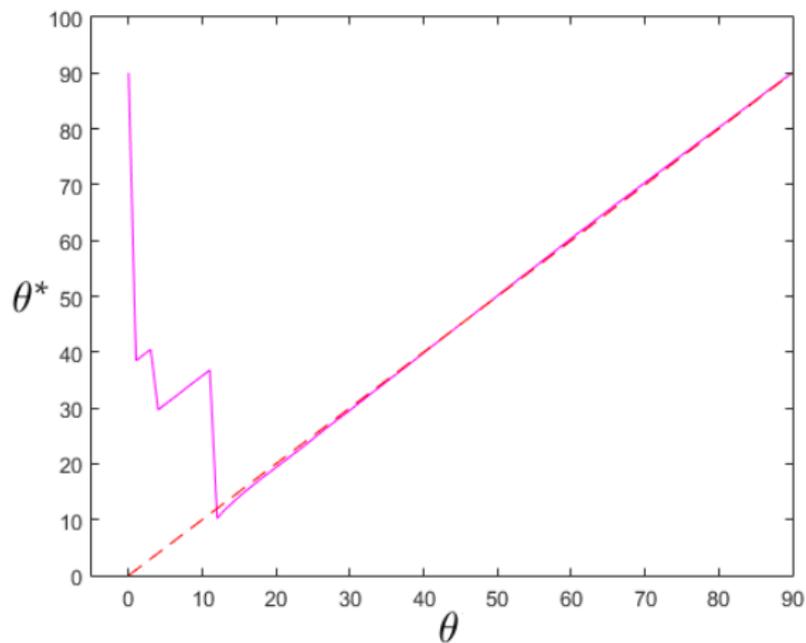
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Figure

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Convergence

- For *smooth* boundary data, proving convergence is routine.
- However, *nonsmooth* boundary data (e.g. images) is much more challenging.

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- In this case, the fact that our weights are non-negative and sum to one means they can be interpreted as a probability density.
- This opens the door to a probabilistic line of attack based on martingales.
- Enables us to prove convergence even for data with jump discontinuities.

Better weights



Figure

Better weights



Figure

Better weights



Figure

Better weights



Figure

Better weights



Figure

Better weights



Figure

Better weights



Figure

Better weights



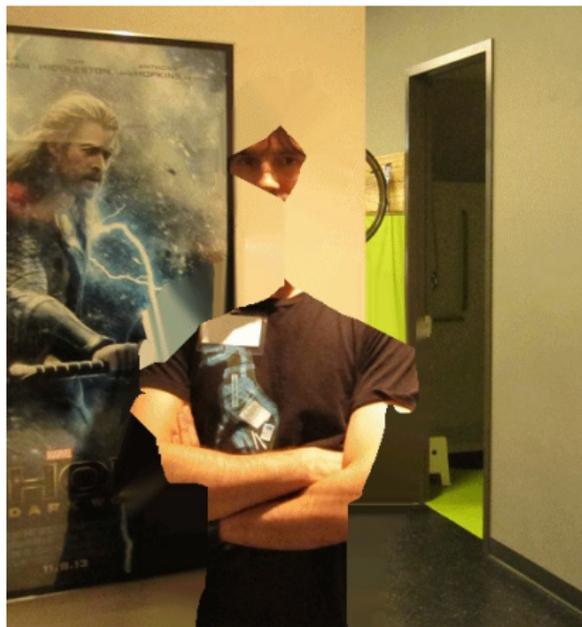
Figure

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Figure

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Figure

The End