Nonlinear Localization and Energy Harvesting in Granular Media

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References

Dark Breathers in Granular Crystals
C. Chong, P.G. Kevrekidis, G. Theocharis, and C. Daraio
PRE **87** (2013) 042202

Damped-Driven Granular Crystals: An Ideal Playground for Dark Breathers and Multibreathers
C. Chong, F. Li, J. Yang, M.O. Williams, I.G. Kevrekidis, P.G. Kevrekidis, and C. Daraio
PRE **89** (2014) 032924

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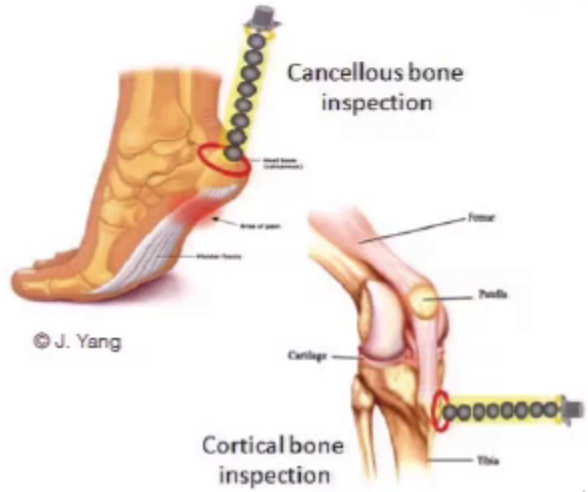
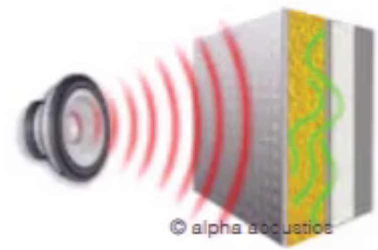
Outline

- Hamiltonian System: Breathers and the NLS
- Vibration Energy Harvesting
- Damped-Driven Dynamics
- Revisiting NLS Prediction

Applications



Non-destructive Evaluation



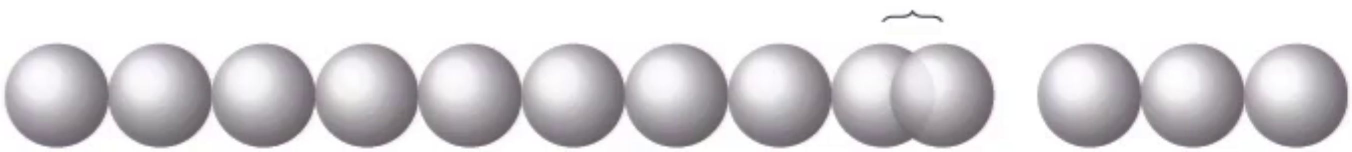
Sound Proofing

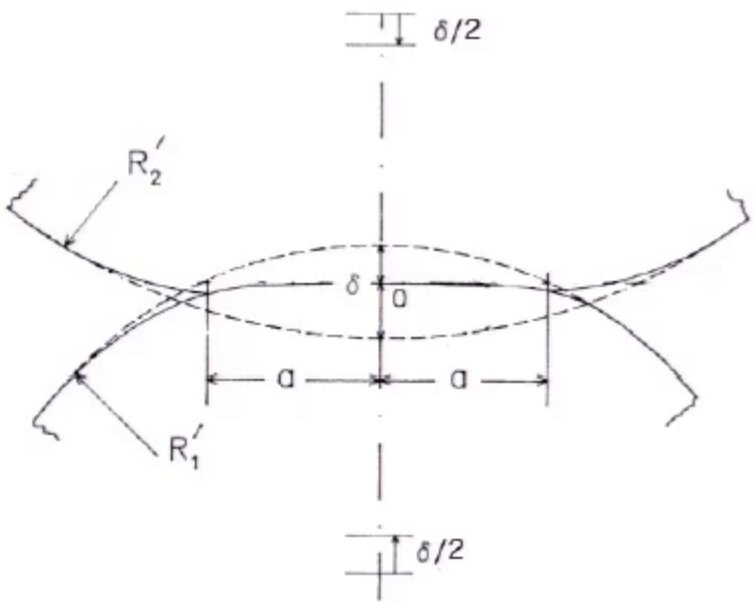
Vibration Energy Harvesting

Impact Absorption

Modeling the Granular Chain

$$\delta_n = u_n - u_{n+1}$$





$$F(\delta) = \gamma\,[\,\delta\,]_+^{3/2}$$
$$[s]_+ = \max(0, s)$$

Source: Johnson, K. L. (1985) *Contact Mechanics* (Cambridge University Press, Cambridge)

Modeling the Granular Chain

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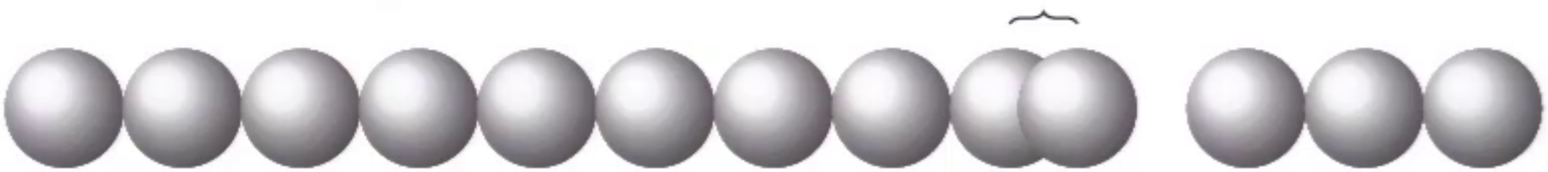
$$[s]_+ = \max(0, s)$$

$$\delta_n = u_n - u_{n+1}$$



Assuming a static load induced by the force $F_0 = \gamma\delta_0^{3/2}$ we have the following model

$$M\partial_t^2 u_n = \gamma[\delta_0 + u_{n-1} - u_n]_+^{3/2} - \gamma[\delta_0 + u_n - u_{n+1}]_+^{3/2}$$

Strain Variable

We will work in terms of the strain variable $y_n = u_n - u_{n+1}$.

$$\ddot{y}_n = F(y_{n+1}) - 2F(y_n) + F(y_{n-1}).$$

A displacement variable solution can be recovered by using the relation,

$$u_{n+1} = u_1 - \sum_{i=1}^{n} y_i$$

where an arbitrary choice for the first node is made.

Weakly Nonlinear Regime

Considering small amplitude solutions $\frac{|y_n|}{\delta_0} \ll 1$, implies we can Taylor expand $F(x)$:

$$F(x) \approx \gamma \delta_0^{3/2} + \frac{3}{2}\gamma \delta_0^{1/2} x + \gamma \frac{3}{8}\delta_0^{-1/2} x^2 - \gamma \frac{3}{48}\delta_0^{-3/2} x^3$$

such that the linear problem becomes:

$$M\ddot{y}_n = \frac{3}{2}\gamma \delta_0^{1/2}(y_{n-1} - 2y_n + y_{n+1})$$

Spectral Situation

The linear problem is solved by,

$$y_n(t) = e^{i(kn+\omega t)}$$

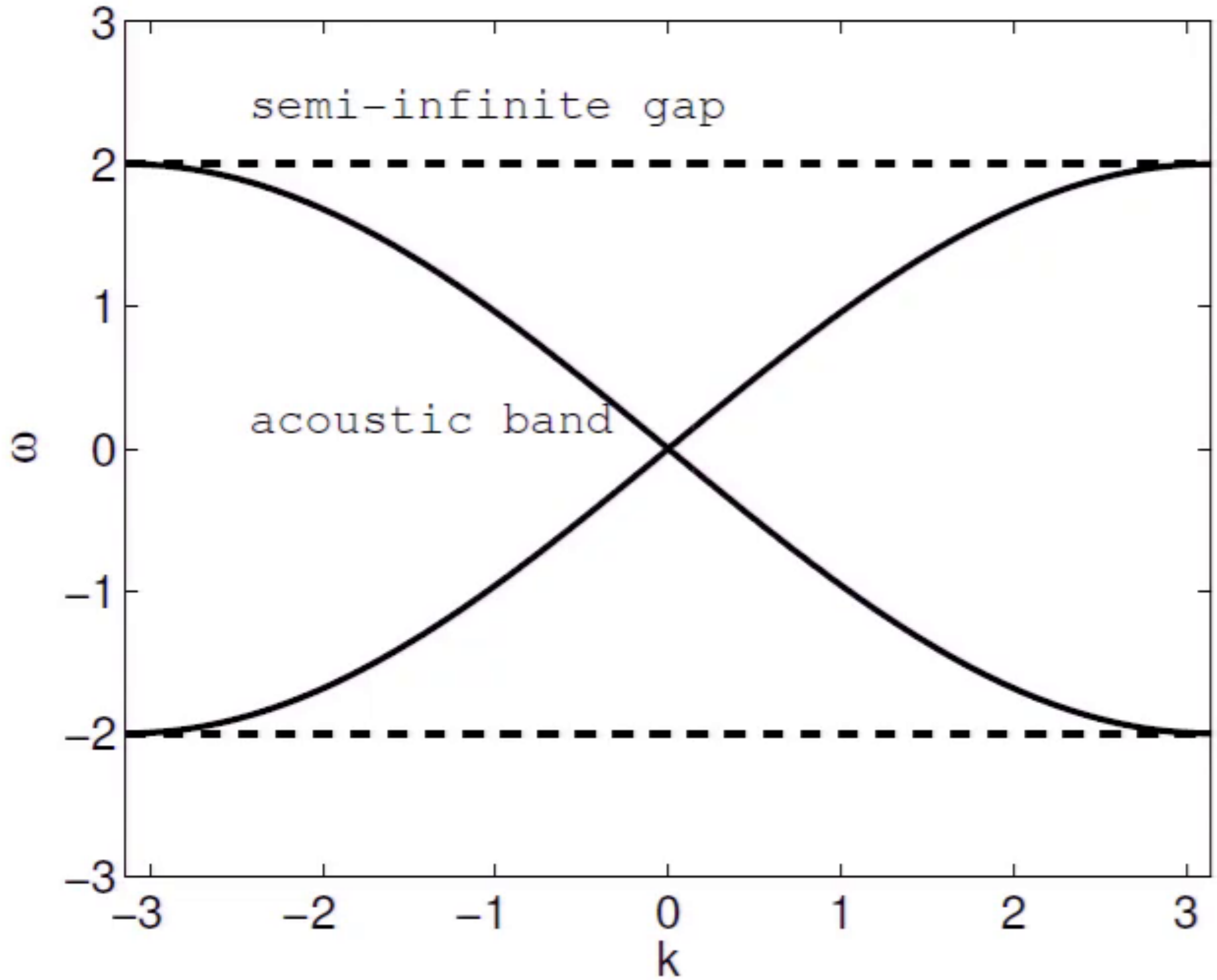
for all $k \in \mathbb{R}$ where ω and k are related through the dispersion relation,

$$\omega(k)^2 = 2K_2/M(1 - \cos k)$$

such that linear spectrum is given by

$$\sigma = \left[-2\sqrt{K_2/M}, 2\sqrt{K_2/M}\right]$$

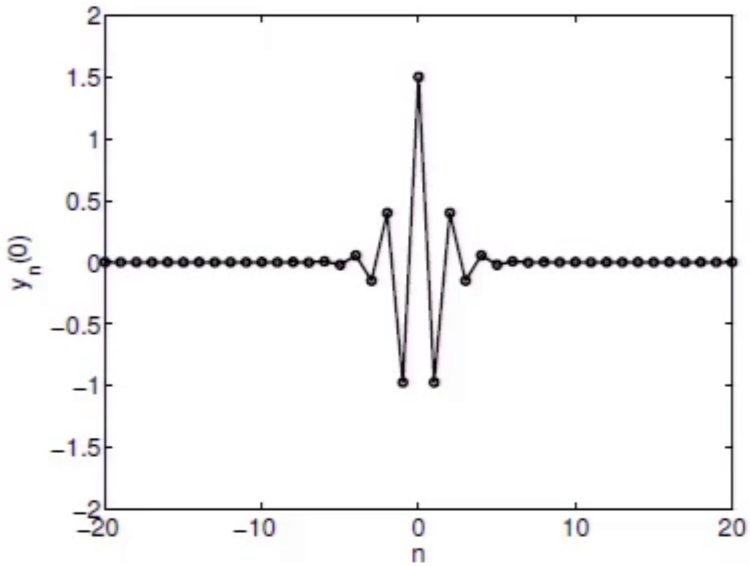
Spectral Situation



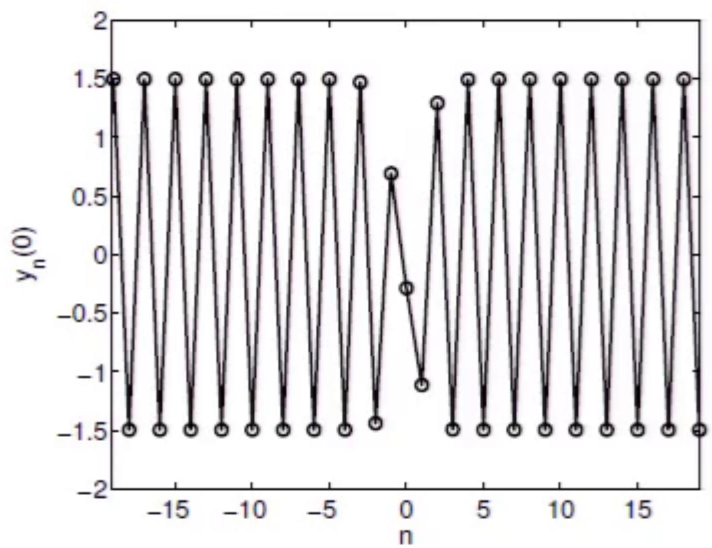
Discrete Breathers

Breathers are time period solutions that are localized in space, e.g.:

$$y_n = (-1)^n a_n \cos(\omega t) + h.o.t.$$



Bright: $a_n = \alpha \text{sech}(\beta n)$



Dark: $a_n = \alpha \tanh(\beta n)$

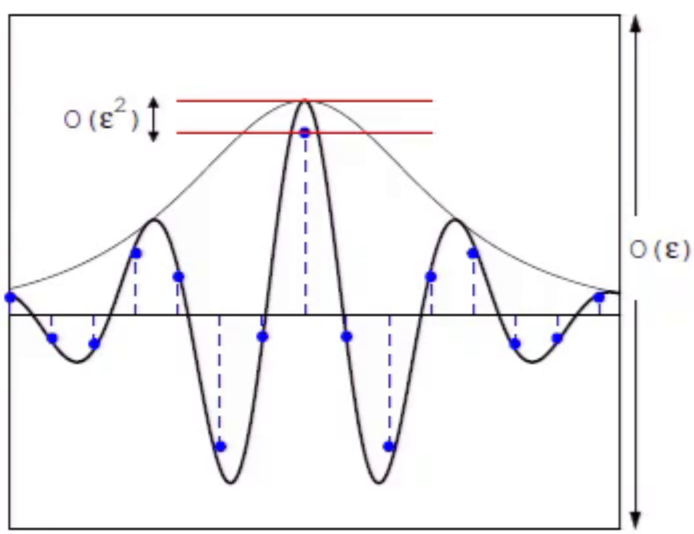
Derivation of NLS Equation

The multiple scale ansatz

$$y_n(t) \approx \psi_n(t) := \epsilon A(X,T)e^{i(k_0 n+\omega_0 t)} + \text{c.c.}, \quad X = \epsilon(n+ct), \quad T = \epsilon^2 t$$

where $\epsilon \ll 1$ is a small parameter, is the canonical ansatz to derive the NLS equation from wave equations e.g. in optics and water waves.

** $A(X,T)$ is *continuous* in both its arguments **



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$\mathcal{O}(\varepsilon^{-1}\varepsilon)$: the dispersion relation $\omega_0 = \omega(k_0)$

$\mathcal{O}(\varepsilon^{-1}\varepsilon^2)$:the group velocity relation $c = \omega'(k_0)$

$\mathcal{O}(\varepsilon^{-1}\varepsilon^3)$:and the nonlinear Schrödinger equation (NLS),

$$i\partial_T A(X,T) + \nu_2 \partial_X^2 A(X,T) + \nu_3 A(X,T)|A(X,T)|^2 = 0$$

The NLS equation

The NLS equation,

$$i\partial_T A(X,T) + \nu_2 \partial_X^2 A(X,T) + \nu_3 A(X,T)|A(X,T)|^2 = 0$$

has steady-states which have the form,

$$A(X,T) = \tilde{A}(X)e^{i\kappa T}$$

where $\kappa \in \mathbb{R}$ and $\tilde{A}(X)$ satisfies the Duffing equation,

$$\partial_X^2 \tilde{A} = \frac{1}{\nu_2}\left(\kappa \tilde{A} - \nu_3 \tilde{A}^3\right)$$

Which has the exact solutions,

$$\tilde{A}(X) = c_1 \tanh\left(c_2 X\right) \qquad \text{if } \nu_3 < 0, \kappa < 0$$

$$\tilde{A}(X) = c_3 \mathrm{sech}\left(c_4 X\right) \qquad \text{if } \nu_3 > 0, \kappa > 0$$

Non-Existence of (bright) breathers

Since we seek standing wave solutions, we chose the wavenumber to be at the edge of the phonon band $k_0 = \pi$, such that the group velocity vanishes, $\omega_0 = 2\sqrt{K_2/M}$, and

$$\nu_3|_{k_0=\pi} = 3K_2K_4 - 4K_3^2 = B < 0.$$

In the case of the granular chain, NLS is defocusing

Thus, the only possible (NLS) steady-state is,

$$\tilde{A}(X) = c_1 \tanh(c_2 X) \qquad \text{if } \nu_3 < 0, \kappa < 0$$

- Bright breathers only if $B > 0$ (modulation instability) [Flach 1997]
- Bright breathers only if $B > 0$ (center manifold) [James 2003]

NLS approximation

Returning to our ansatz we have the following approximation,

$$
\begin{aligned}
y_n(t) &\approx \epsilon \tilde{A} e^{i\kappa T} e^{i(k_0 n + \omega_0 t)} + \text{c.c.} \\
&= 2\epsilon(-1)^n \sqrt{\frac{\kappa}{\nu_3}} \tanh\left(\sqrt{\frac{-\kappa}{2\nu_2}}\epsilon n\right) \cos(\omega_b t) \qquad (\text{for } k_0 = \pi)
\end{aligned}
$$

where $\omega_b = \omega_0 + \kappa\epsilon^2$ is the frequency of the breather, where $\kappa < 0$ is a fixed but arbitrary parameter.

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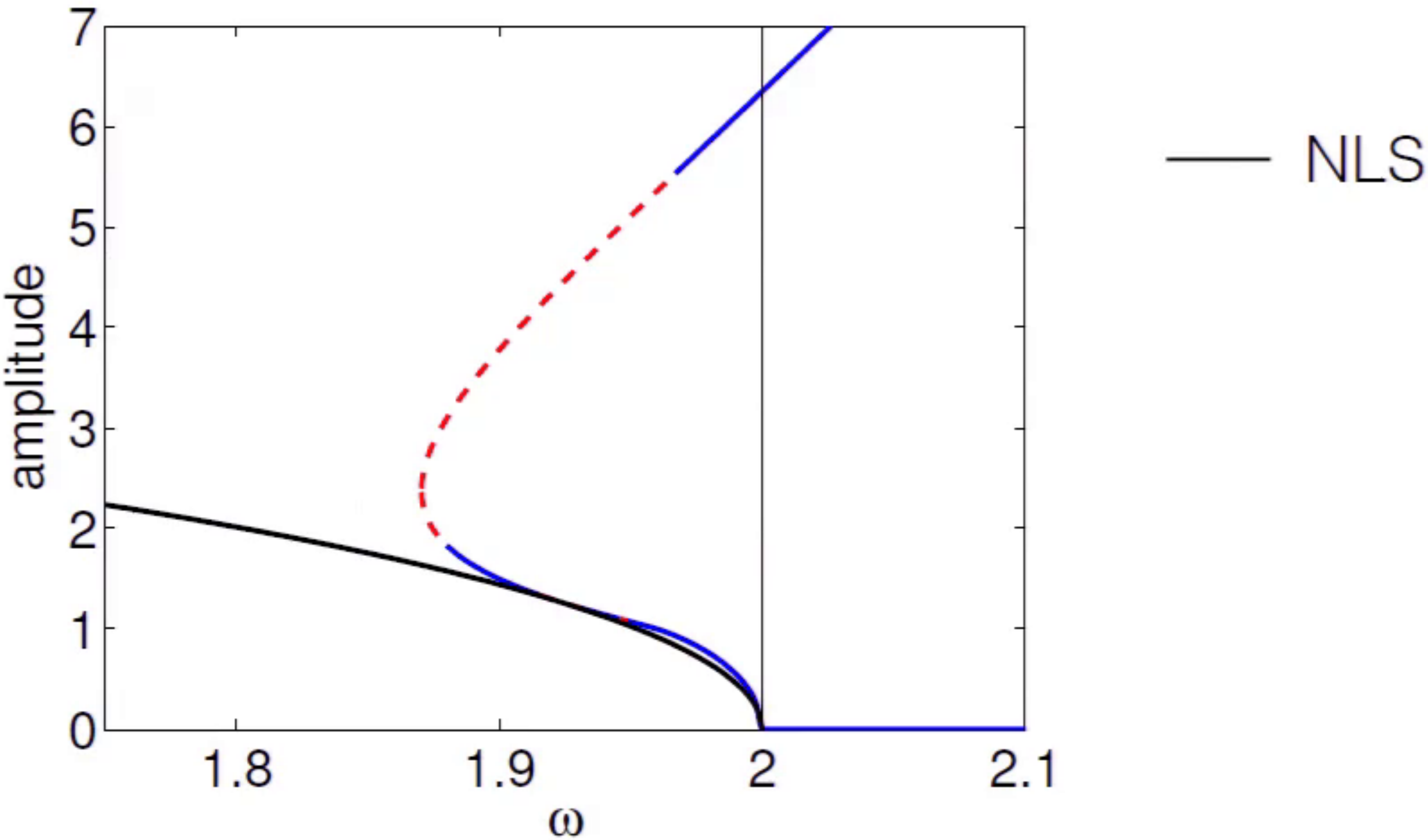
Time-period solutions and their spectral stability and bifurcations

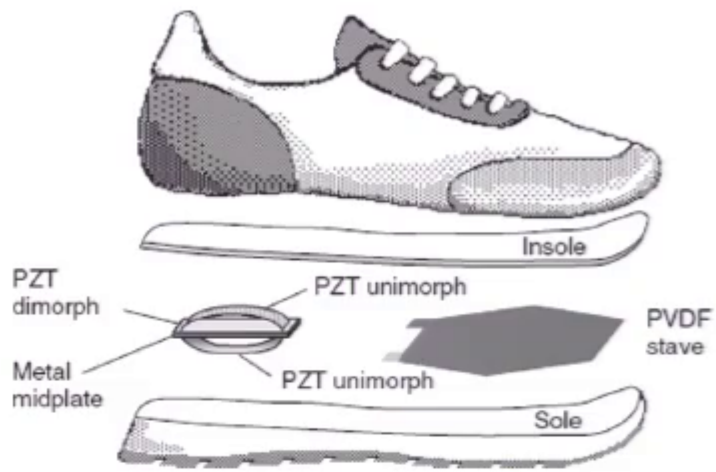
- We numerically compute time-periodic solutions by finding roots of the map $F(x) = x(0) - x(T)$ via Newton iterations.
- Solutions are continued via pseudo-arclength continuation
- We linearize about these time periodic solutions and compute the Floquet multipliers of the resulting Hills equation

Comparison of NLS Prediction to Numerics



Vibration Energy Harvesting

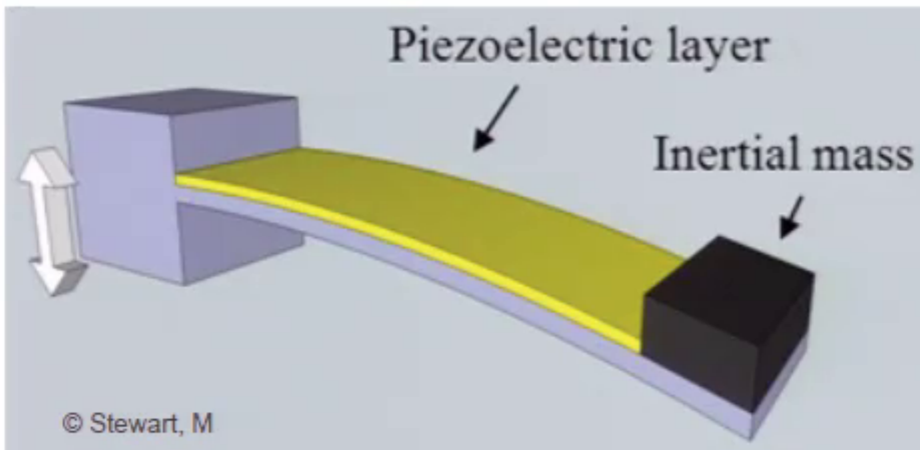
- Goal: Enable self-powered electronic devices by scavenging ambient energy (vibrations)
- Energy Conversion: Electromagnetic Induction, Electrostatic Transduction, **Piezoelectric Transduction**
- Ambient vibration:
  white noise <—> **harmonic**



Energy scavenging with shoe mounted piezoelectrics
N. A. Shenck, J. A. Paradiso
IEEE Micro, vol. 21, p. 30–42, 2001-05/06

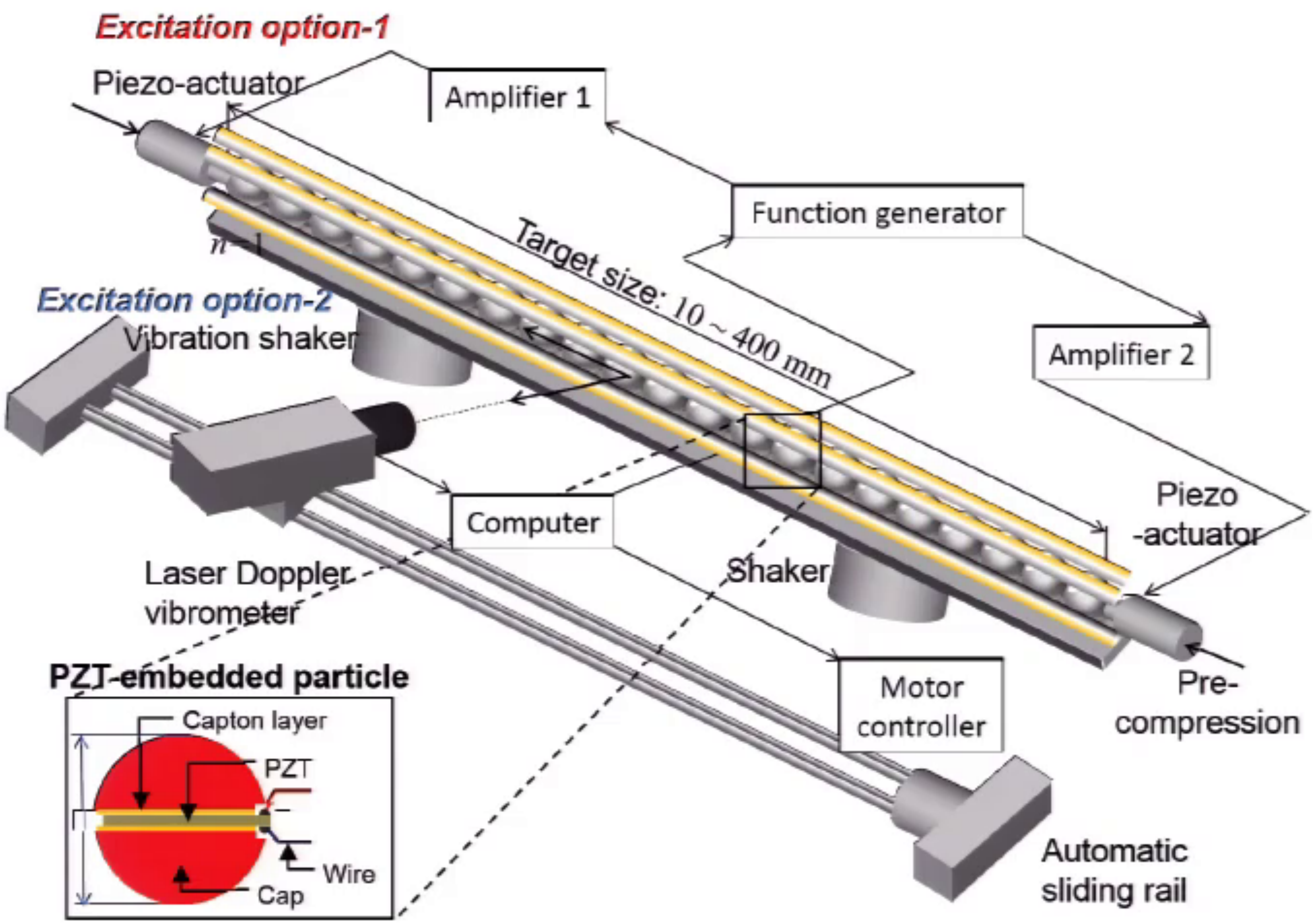
Linear Harvester

- Conventional devices are linear and well studied (Erturk 2013)
- Inefficient if input frequency varies or is not known a priori

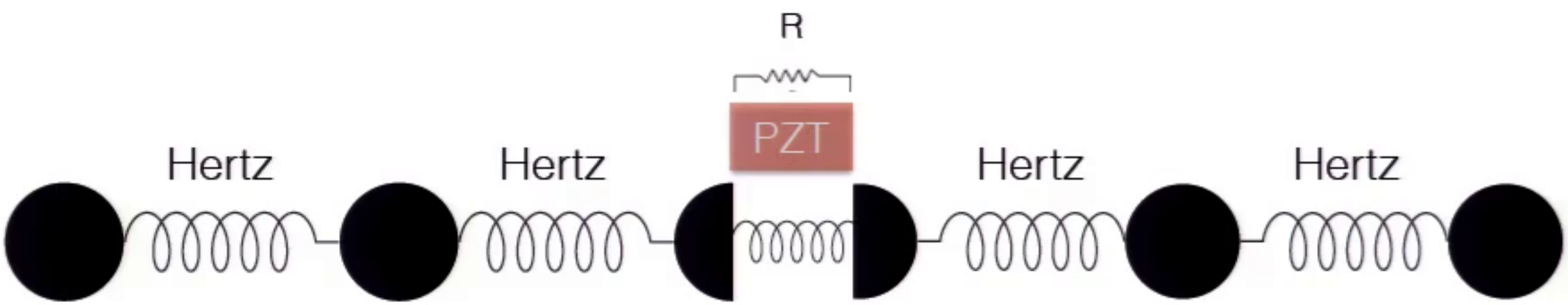


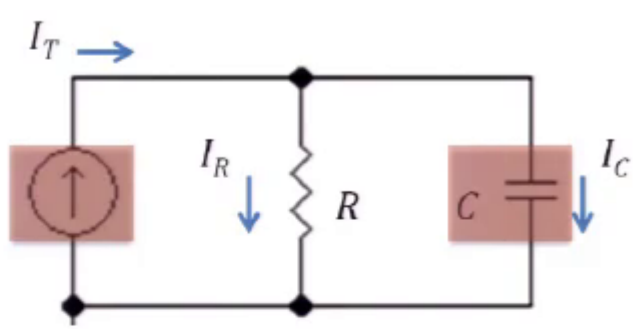
- We consider a spatially extended and nonlinear system and test its efficiency as an energy harvester

Experimental Set-up for Energy Harvesting



Electromechanical Model





Electromechanical Model

Constitutive relation

$$D = \epsilon^T E + d_{33} T$$
$$S = d_{33} E + s^E T$$

+

Granular Model

$$M\partial_t^2 u_n = \gamma[\delta_0 + u_{n-1} - u_n]_+^{3/2} - \gamma[\delta_0 + u_n - u_{n+1}]_+^{3/2}$$

$$\begin{aligned}
M\ddot{u}_n &= \gamma[\delta_0 + u_{n-1} - u_n]_+^{3/2} - \gamma[\delta_0 + u_n - u_{n+1}]_+^{3/2} - \frac{M}{\tau}\dot{u}_n, \quad n \notin \{m, m+1\} \\
\frac{M}{2}\ddot{u}_m &= \gamma[\delta_0 + u_{m-1} - u_m]_+^{3/2} - K_a(\delta_1 + u_m - u_{m+1}) - d_{33}K_a V - \frac{M}{2\tau}\dot{u}_m, \\
\frac{M}{2}\ddot{u}_{m+1} &= K_a(\delta_1 + u_m - u_{m+1}) - \gamma[\delta_0 + u_{m+1} - u_{m+2}]_+^{3/2} + d_{33}K_a V - \frac{M}{2\tau}\dot{u}_{m+1}, \\
RC(1-k^2)\dot{V} &= d_{33}K_a R(\dot{u}_m - \dot{u}_{m+1}) - V
\end{aligned}$$

$$k^2 = \frac{d_{33}^2}{S^E \epsilon^T}, \qquad K_a = \frac{A}{S^E d} \qquad C = \frac{\epsilon^T A}{d}$$

<u>Boundary Actuation</u>

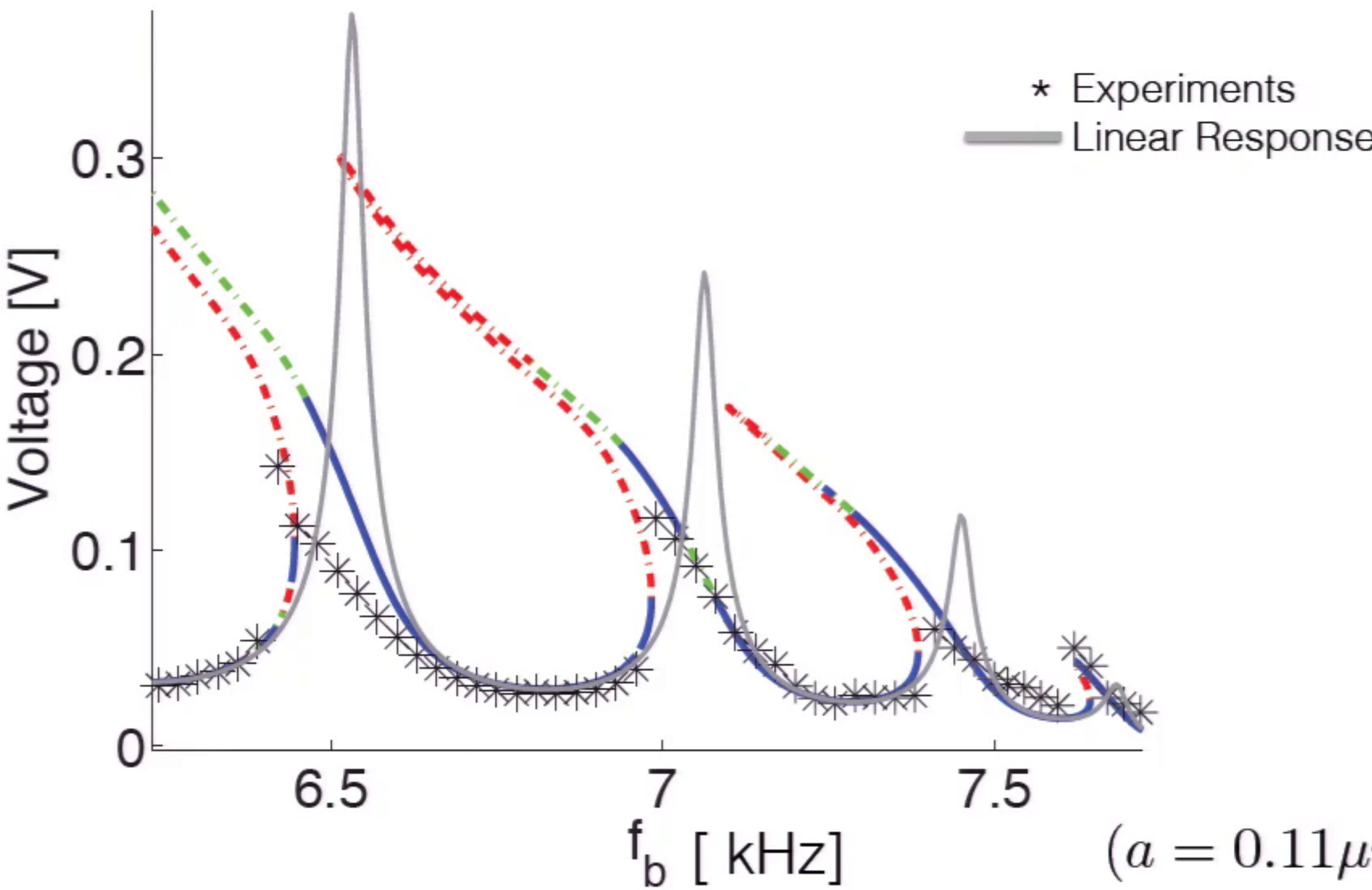
$$u_0 = a\cos(2\pi f_b t)$$
$$u_{N+1} = b\cos(2\pi f_b t)$$

Experimental Values and Set-up

Mechanical Parameters			Electrical Parameters		
Bead Mass	M	28.2 [g]	Piezo Constant	d_{33}	$360 \cdot 10^{-12}$ m/V
Bead Young's Modulus	E	200 [Gpa]	Piezo Permittivity	ϵ^T	$14166 \cdot 10^{-12}$ F/m
Bead Radius	r	9.53 [mm]	Piezo Compliance	S^E	72 Gpa^{-1}
Bead Poisson Ration	ν	0.3	Piezo Disc Area	A	28.353 mm
Damping Coefficient	τ	5 [ms]	Piezo Thickness	d	0.3 mm
			Resistance	R	$3\,K\Omega$

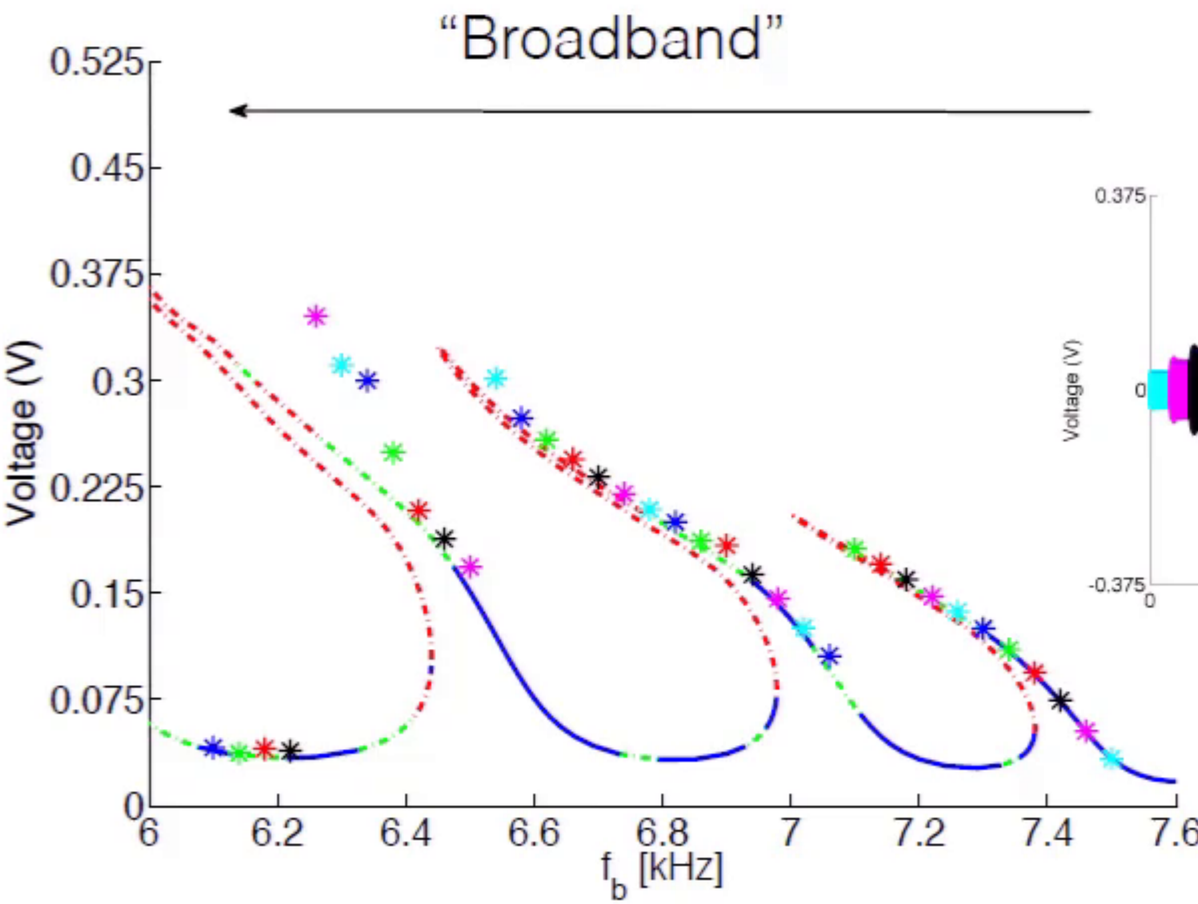
TABLE I: Theoretical values of the electromechanical parameters.

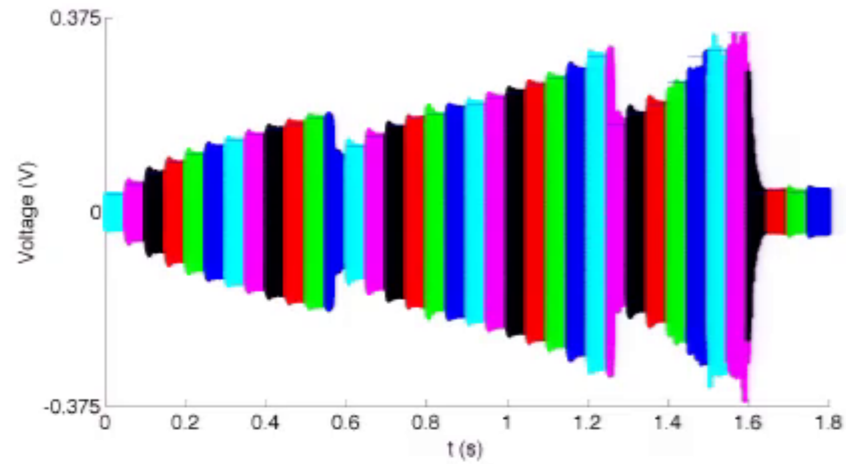
Theory vs Experiments



$(a = 0.11\mu m)$

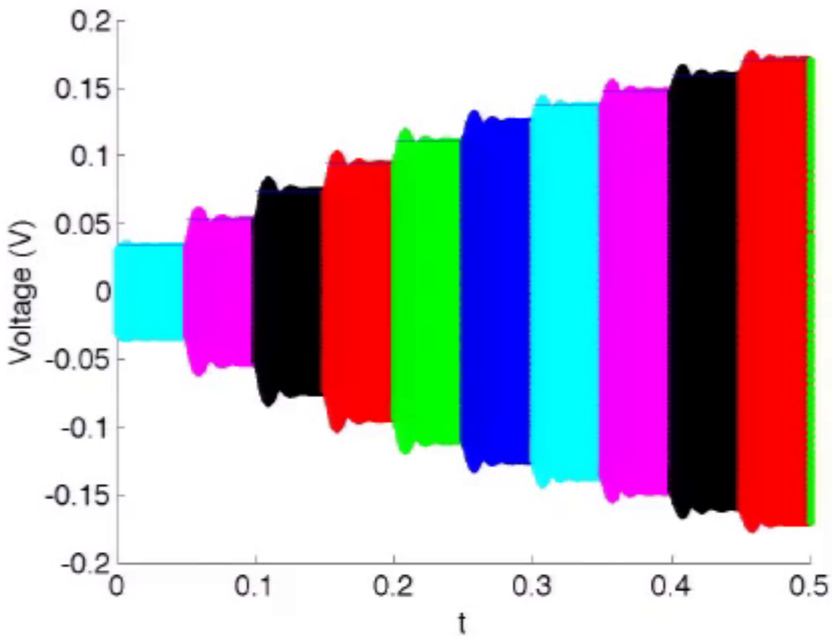
High to Low Frequency Sweep



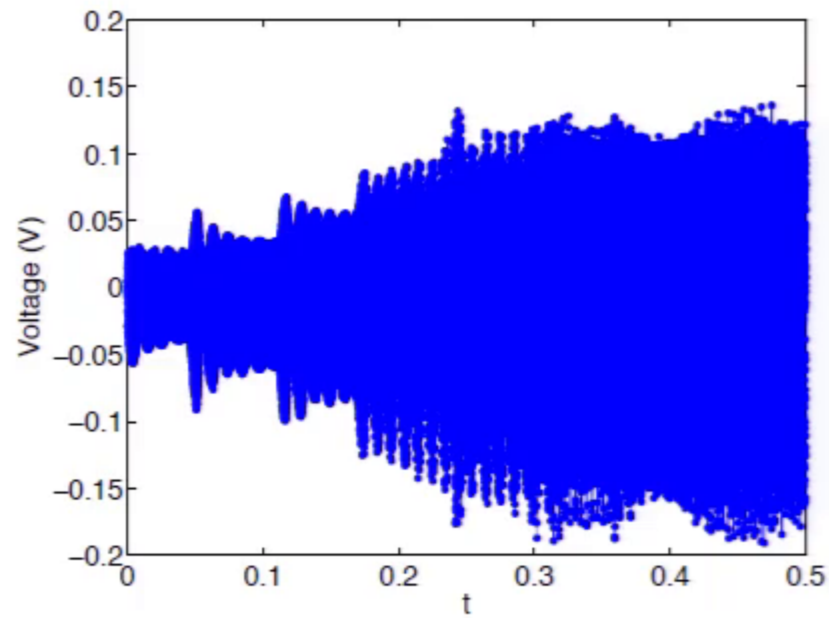


decrease frequency 40 Hz
every .05 seconds

Numerics vs Experiment: High to Low Frequency Sweep

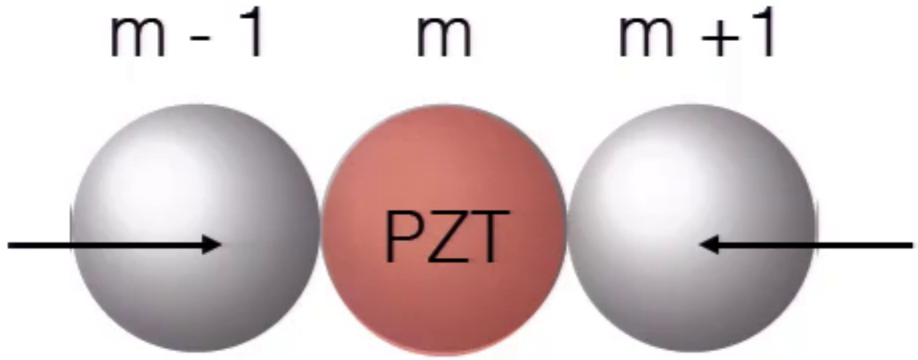


Theory



Experiment

Mechanical System



The amount the PZT bead is squeezed is related to voltage production:

$$s_m(t) = u_{m-1} - u_{m+1}$$

In terms of the strain $y_m = u_{m-1} - u_m$ we have,

$$s_m(t) = y_m + y_{m+1}$$

NLS Prediction

Recall, the NLS approximation:

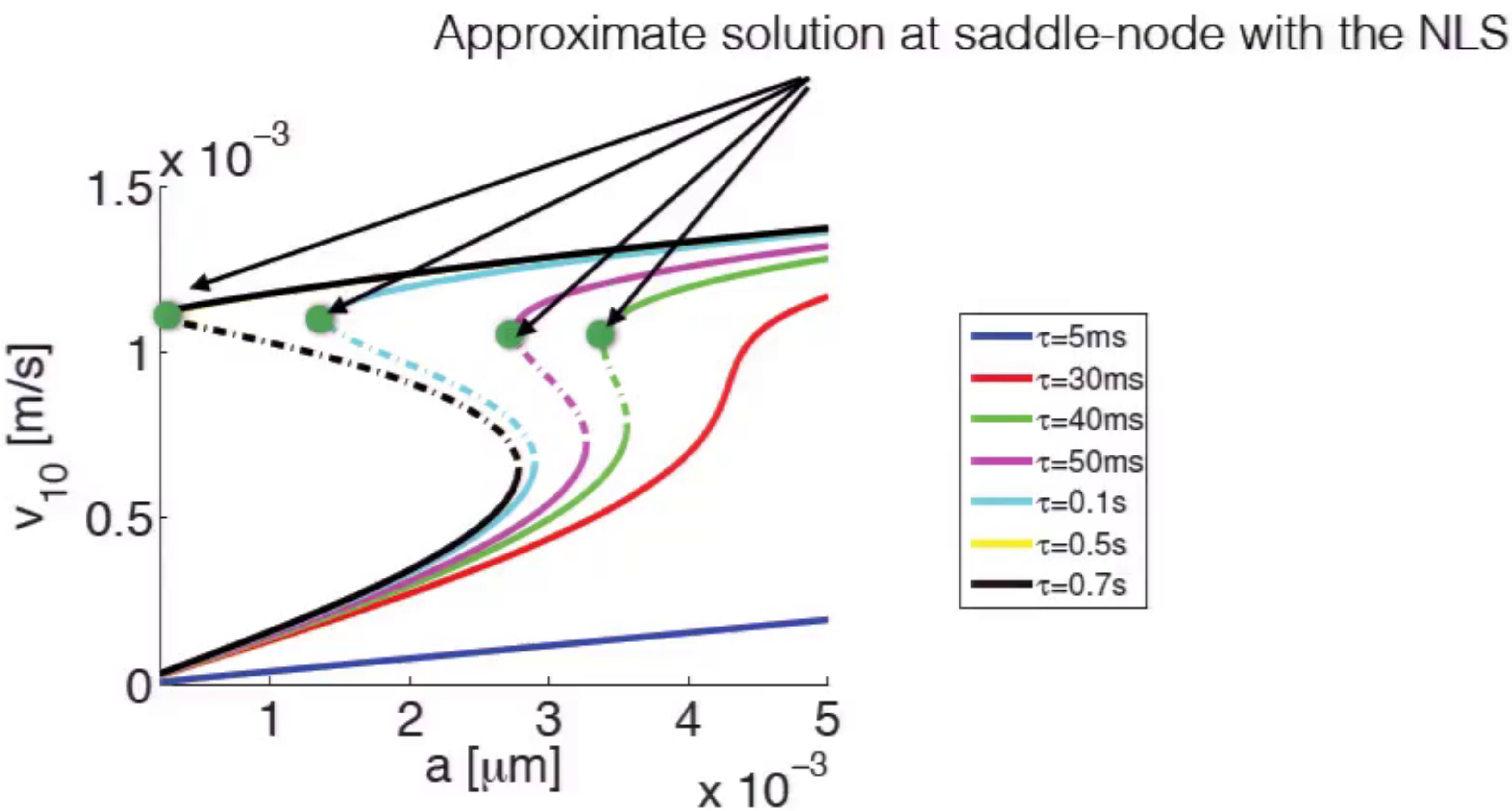
$$y_n(t) = 2\varepsilon(-1)^n\sqrt{\frac{\kappa}{\nu_3}}\tanh\left(\sqrt{\frac{-\kappa}{2\nu_2}}\varepsilon(n-x_0)\right)\cos(\omega_b t)$$

Where $\omega_b = \omega_0 + \kappa\epsilon^2$ is the frequency and $\kappa < 0$ is a fixed but arbitrary parameter. In terms of the "squeeze", we have

$$s_n(t) = 2\epsilon(-1)^{n+1}\sqrt{\tfrac{\kappa}{\nu_3}}\left(\tanh\left(\sqrt{\tfrac{-\kappa}{2\nu_2}}\epsilon(n+1)\right) - \tanh\left(\sqrt{\tfrac{-\kappa}{2\nu_2}}\epsilon n\right)\right)\cos(\omega_b t)$$

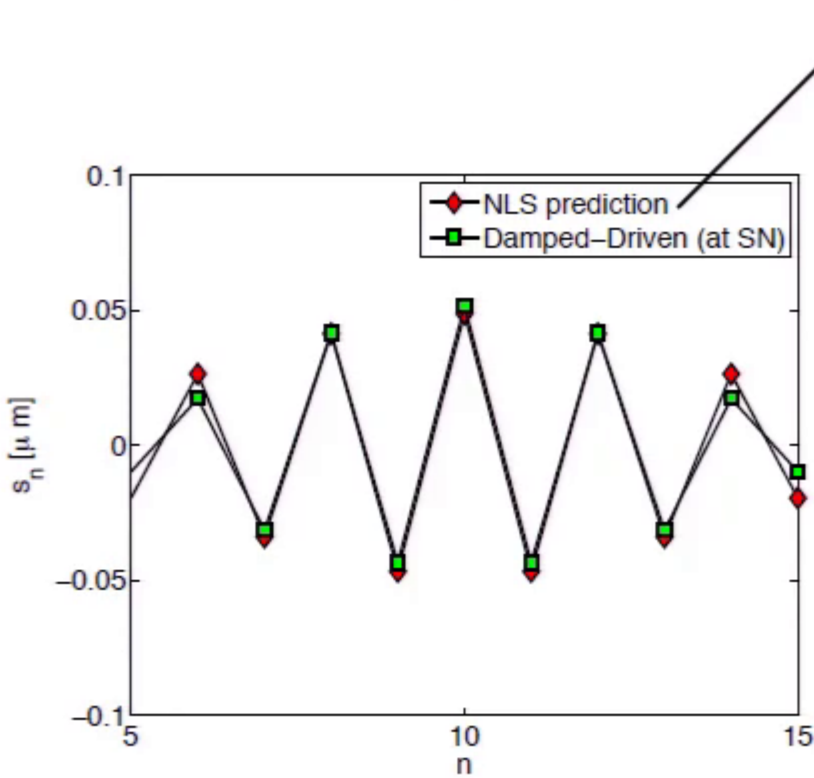
$$\approx 2\epsilon^2(-1)^{n+1}\sqrt{\tfrac{\kappa}{\nu_3}}\left(1 - \tanh^2\left(\sqrt{\tfrac{-\kappa}{2\nu_2}}\epsilon n\right)\right)\cos(\omega_b t)$$

Continuation in Amplitude and the NLS approximation



NLS Prediction

$$s_n(t) = 2\epsilon(-1)^{n+1}\sqrt{\frac{\kappa}{\nu_3}}\left(\tanh\left(\sqrt{\frac{-\kappa}{2\nu_2}}\epsilon(n+1.5)\right) - \tanh\left(\sqrt{\frac{-\kappa}{2\nu_2}}\epsilon(n+.5)\right)\right)\cos(\omega_b t)$$



We can back calculate the voltage: