# Dissipative Time Reversal Technique for Photoacoustic Tomography in a Cavity

#### Linh Nguyen (University of Idaho) Leonid Kunyansky (University of Arizona)

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## Outline



### Free Space Photoacoustic Tomography

- Time Reversal Technique
- Photoacoustic Tomography in A Cavity
  - Mathematical Model and Difficulty
  - Current Time Reversal Techniques
  - Dissipative Time reversal

# Photoacoustic Tomography

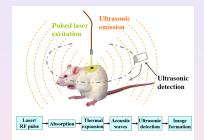


Figure: Photoacoustic tomography - Source: Wikipedia

- Ultrasonic pressure measured on observation surface S.
- One needs to find the initial pressure f(x).

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Wave propagation:

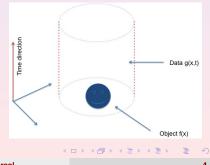
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Data: 
$$g = u|_{S \times (0,T)}$$
.

Problem: Find f from g

Assumptions:

- S is a closed surface.
- T can be either a fixed finite number or infinity.



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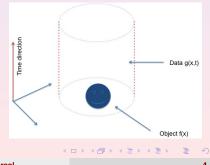
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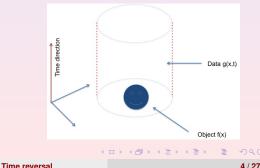
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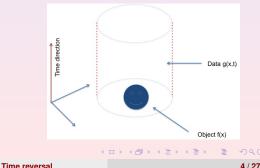
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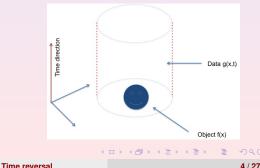
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## Techniques

- Exact Inversion formulas: Norton & Linzer ('79-'81), Xu & Wang ('02, '04), Finch, Patch & Rakesh ('04), Finch, Haltmeier & Rakesh ('07), Kunyansky ('07, '11,'15), N. ('09), Palamodov ('11 & 14), Natterer ('12), Salman ('12), Haltmeier ('11 & '13) ...
- Series solutions: Agranovsky & Kuchment ('07), Kunyansky ('07) ...
- Time-reversal: Finch, Patch & Rakesh ('04), Hristova, Kuchment & N. ('08), Hristova ('09), Stefanov & Uhlmann ('09 & '11), J. Qian, Stefanov, G. Uhlmann & H. Zhao ('11)
- Iterative method: Anastasio, Su, & Oraevsky ('12), Belhachmi, Glatz, & Scherzer ('15) ...

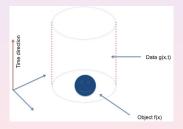
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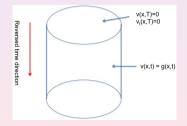
## Näive Time Reversal

Consider the time reversed problem

$$\begin{cases} v_{tt}(x,t) - c^{2}(x) \Delta v(x,t) = 0, \ (x,t) \in \Omega \times (0,T), \\ v(x,T) = 0, \quad v_{t}(x,T) = 0, \quad x \in \Omega, \\ v(x,t) = g(x,t), \quad (x,t) \in S \times (0,T). \end{cases}$$



#### Figure: Forward model



#### Figure: Time reversed

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# Näive Time Reversal

Denote: Ag = v(., 0). Then,

$$Ag o f$$
 in  $H^1(\Omega)$ ,  $T \to \infty$ .

Idea:

• w = u - v solves wave equation with zero Dirichlet BC.

• Conservation of energy:

$$E_{\Omega}(w,0) = E_{\Omega}(w,T) = E_{\Omega}(u,T).$$

• Then,

$$\|\mathcal{A}g - f\|_{H^1(\Omega)} = \|w(.,0)\|_{H^1(\Omega)} \leq C E_{\Omega}(w,0) \rightarrow 0.$$

A technical issue: The initial condition and boundary condition may not be compatible.

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# **Fixing The Problem**

• Fix 1: Hristova, Kuchment, & N. ('08), Hristova ('09).

$$\begin{cases} v_{tt}(x,t) - c^{2}(x) \Delta v(x,t) = 0, \ (x,t) \in \Omega \times (0,T), \\ v(x,T) = 0, \quad v_{t}(x,T) = 0, \quad x \in \Omega, \\ v(x,t) = \chi(t)g(x,t), \quad (x,t) \in S \times (0,T). \end{cases}$$

Here,  $\chi \equiv 1$  for  $t \in [0, T - \epsilon]$  and  $\chi(T) = 0$ .

• Fix 2: Stefanov& Uhlmann ('09)

$$\begin{cases} v_{tt}(x,t) - c^2(x) \Delta v(x,t) = 0, \quad (x,t) \in \Omega \times (0,T], \\ v(x,T) = \phi(x), \quad v_t(x,T) = 0, \quad x \in \Omega, \\ v(x,t) = g(x,t), \quad (x,t) \in S \times (0,T). \end{cases}$$

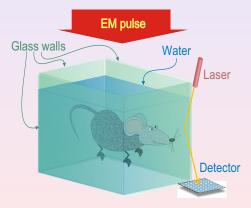
Here,  $\phi$  harmonic extension of g(., T) on  $\Omega$ . Stefanov & Uhlmann obtained a Neumann series solution.

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Time reversal

## Photoacoustic Tomography in A Cavity

Photoacoustic setup at University College London:



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#### Mathematical model:

$$\begin{cases} u_{tt}(x,t) - c^2(x)\Delta u(x,t) = 0, \ (x,t) \in \Omega \times (0,T), \\ u(x,0) = f(x), \ u_t(x,0) = 0, \quad x \in \Omega, \\ \partial_{\nu}u(x,t) = 0, \ (x,t) \in \partial\Omega \times (0,T). \end{cases}$$

Data:  $g = u|_{\Gamma \times (0,T)}$ , where  $\Gamma \subset \partial \Omega$ . Problem: Find *f* given *g*. Equivalently, invert  $\Lambda : f \to g$ .

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## **Historical Remarks**

- Contributions: Kunyansky, Holman, & Cox ('13): series expansion; Holman & Kunyansky ('15): gradual time reversal; Stefanov & Yang ('15): averaged time reversal; Acosta & Montalto ('15): conjugate gradient method & time reversal for partially reflecting boundary.
- Recently: Stefanov & Yang ('16): Landweber's iteration, Chernova & Okasen ('16): Time reversal.
- Difficulty: The energy  $E_{\Omega}(u, t)$  is conserved for all time *t*.

Gradual time reversal - Holman & Kunyansky ('15):

$$\begin{cases} v_{tt}(x,t) - c^2(x) \Delta v(x,t) = 0, \ (x,t) \in \Omega \times (0,T), \\ v(x,T) = 0, \quad v_t(x,T) = 0, \quad x \in \Omega, \\ v(x,t) = \chi(\epsilon t) g(x,t), \quad (x,t) \in \partial\Omega \times (0,T). \end{cases}$$

Here,  $\epsilon = \frac{1}{T}$  and  $\chi \in C_0^{\infty}[0, 1)$ . Define:

$$Ag = v(., 0).$$

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#### Then,

$$Ag = v(.,0) \rightarrow f$$
 weakly  $H^1(\Omega)$  as  $\epsilon \rightarrow 0$ ,

under assumption on the eigenvalues of Laplacian with Dirichlet and Neumann boundary condition.

#### Features:

Adaptable to partial data problem (i.e., Γ ≠ ∂Ω).
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- The rate of convergence is unclear.

Rough idea - Stefanov & Yang ('15):

- Time reversal at all time *τ* ∈ (0, *T*] with Dirichlet boundary condition. Then average the result.
- Neumann series solution

$$f=\sum_{m=0}^{\infty}K^{m}Ag=f_{0}+f_{1}+\ldots,$$

where  $K = I - A\Lambda$ .

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#### Features:

- Converges exponentially as the  $N \to \infty$  as long as  $T > \frac{1}{2}T(\Omega)$ .
- **2** Not clear that  $f_0 \to f$  as  $T \to \infty$ .
- Theory not available for partial data problem (but some convincing numerical examples).

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## **New Boundary Condition**

Consider the time reversed problem - N. & Kunyansky('16):

$$\begin{cases} v_{tt}(x,t) - c^2(x) \Delta v(x,t) = 0, \ (x,t) \in \Omega \times (0,T), \\ v(x,T) = 0, \quad v_t(x,T) = 0, \quad x \in \Omega, \\ \partial_n v(x,t) - \lambda v_t(x,t) = -\lambda g_t(x,t), (x,t) \in \partial\Omega \times [0,T]. \end{cases}$$

Define: Ag = v(., 0).

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### Results

#### Assume that $T(\Omega) < \infty$ . Then, If $T \to \infty$ , then

 $Ag \rightarrow f$  exponentially.

**2** If  $T > T(\Omega)$ , one has the Neumann series formula:

$$f=\sum_{m=0}^{\infty}K^{m}Ag=f_{0}+f_{1}+\ldots,$$

where  $K = I - A\Lambda$ .

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## **Quick Proof**

Let w = u - v. Then,

$$\begin{cases} w_{tt}(x,t) - c^2(x) \Delta w(x,t) = 0, \ (x,t) \in \Omega \times (0,T), \\ w(x,T) = 0, \quad w_t(x,T) = 0, \quad x \in \Omega, \\ \partial_n w(x,t) - \lambda w_t(x,t) = 0, (x,t) \in \partial\Omega \times (0,T). \end{cases}$$

We have

$$\mathsf{E}_{\Omega}(w,T) - \mathsf{E}_{\Omega}(w,0) = \iint\limits_{\partial\Omega imes(0,T)} \lambda |w_t(x,t)|^2 \, dx \, dt.$$

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## New Idea: Quick Proof

Key point: the energy of *w* leaks out on the boundary (going backward in time). Therefore Bardos, Lebeau & Rauch('92):

• As  $T \to \infty$ ,

 $E_{\Omega}(w,0) 
ightarrow 0$  exponentially.

That is,  $\nu(.,0) \rightarrow f$  exponentially.

• If  $T > T(\Omega)$ , then

$$E_{\Omega}(w,0) \leq (1-\epsilon) E_{\Omega}(w,T).$$

That is,

$$\|f - A \wedge f\| \leq (1 - \epsilon) \|f\|_{H^1(\Omega)}.$$

 $\implies$  Neumann series formula.

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## Partial Data Problem

Consider  $\Gamma \neq \partial \Omega$ . Assumption (GCC): All the generalized geodesics rays hit  $\Gamma$  in a positive time  $T_{\Gamma} < T$ . Then

• Neumann series formula:

$$f=\sum_{m=0}^{\infty}K^{m}Ag=f_{0}+f_{1}+\ldots,$$

where  $K = I - A \Lambda$ .

• Moreover,

 $f_0 = Ag \rightarrow f$  exponentially as  $T \rightarrow \infty$ .

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# Numerical Experiments (0/4)



Figure: Phantom

Numerical setup:

- Domain  $\Omega$  is a a square of size 2  $\times$  2.
- Sound speed c = 1,  $T(\Omega) = 2\sqrt{2}$ .
- Experiment for full data:  $T > \sqrt{2}$ .
- Experiment for partial data: left and lower edges,  $T > T(\Omega) = 2\sqrt{2}$ .

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# Numerical Experiments (1/4)

Reconstruction with large time: T = 5.

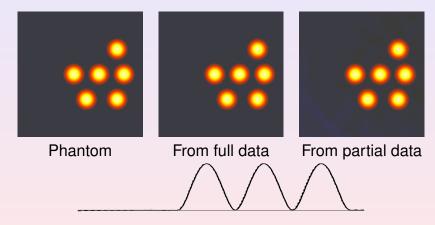
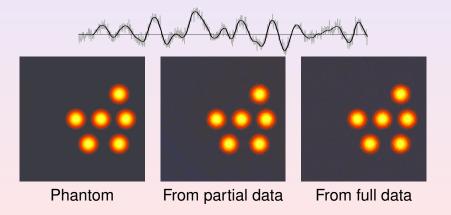


Figure: Gray: phantom, dashed: partial data, black: full data

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# Numerical Experiments (2/4)

Reconstruction with large time: T = 5. Noisy data: 50% noise (in  $L^2$  norm).



## Numerical Experiments (2/4)



Figure: Central line profile

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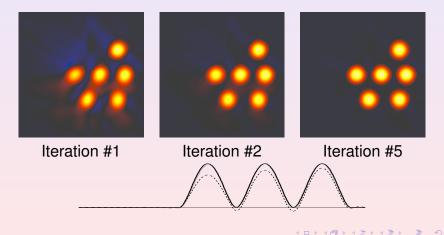
**Time reversal** 

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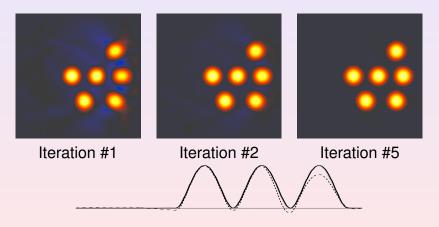
## Numerical Experiments (3/4)

Reconstruction for full boundary data with short time  $T = 1.6 > \sqrt{2}$ .



# Numerical Experiments (4/4)

#### Reconstruction for partial data with short time $T = 3 > 2\sqrt{2}$ .



### Questions

#### Thank you for your attention!



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**Time reversal**