



MAX PLANCK INSTITUTE  
FOR DYNAMICS OF COMPLEX  
TECHNICAL SYSTEMS  
MAGDEBURG



NUMERICAL LINEAR ALGEBRA  
FOR DYNAMICAL SYSTEMS

# Fast Iterative Solvers for a Coupled Cahn–Hilliard/Navier–Stokes System

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October 27, 2015

Thanks to  
Christian Kahle  
(Universität Hamburg).





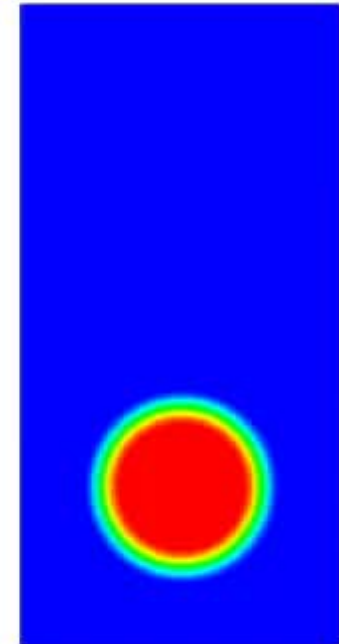
1. Motivation
2. Diffuse Interface Approach
3. Cahn–Hilliard/Navier–Stokes Model
4. Linear Systems and Preconditioning
5. Numerical Results



## Setting

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- two-phase flow
- two immiscible, incompressible viscous fluids
- different physical properties
  - densities  $\tilde{\rho}_1, \tilde{\rho}_2$
  - viscosities  $\tilde{\eta}_1, \tilde{\eta}_2$
- diffuse interface model
  - to cope with topological changes



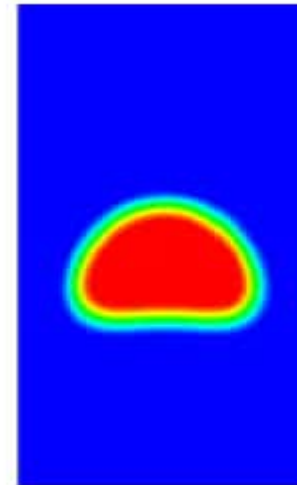
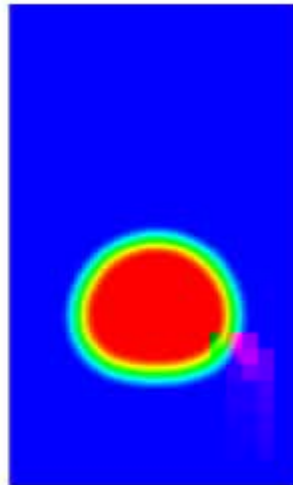
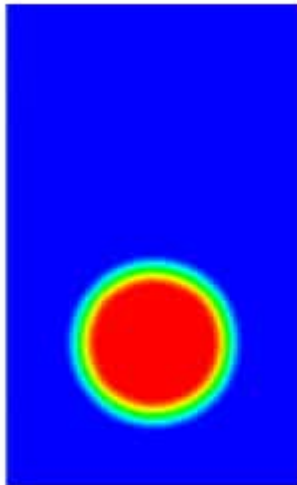


## Examples

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Model for the dynamics of multiphase fluids

- rising bubble / falling droplet



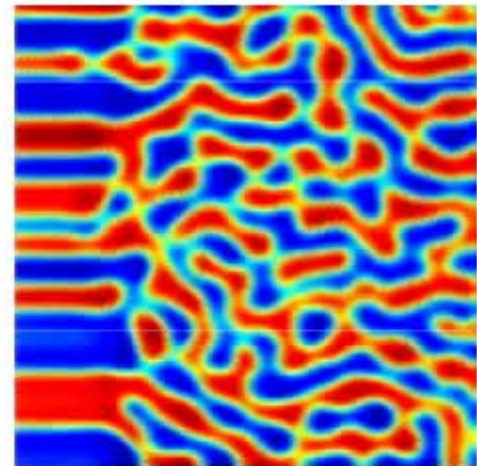
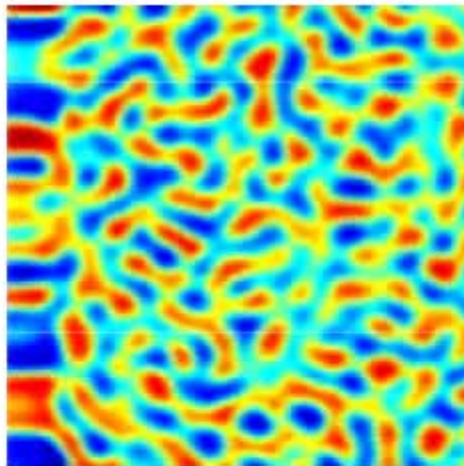
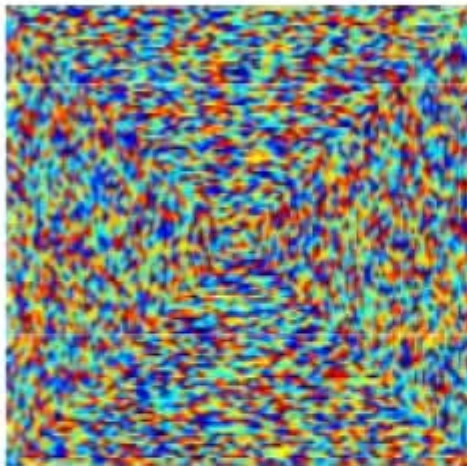


## Examples

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Model for the dynamics of multiphase fluids

- rising bubble / falling droplet
- mixing of fluids in a driven cavity
- spinodal decomposition under shear



[KAY/WELFORD '07]



## Examples

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Model for the dynamics of multiphase fluids

- rising bubble / falling droplet
- mixing of fluids in a driven cavity
- spinodal decomposition under shear
- Rayleigh–Taylor instability



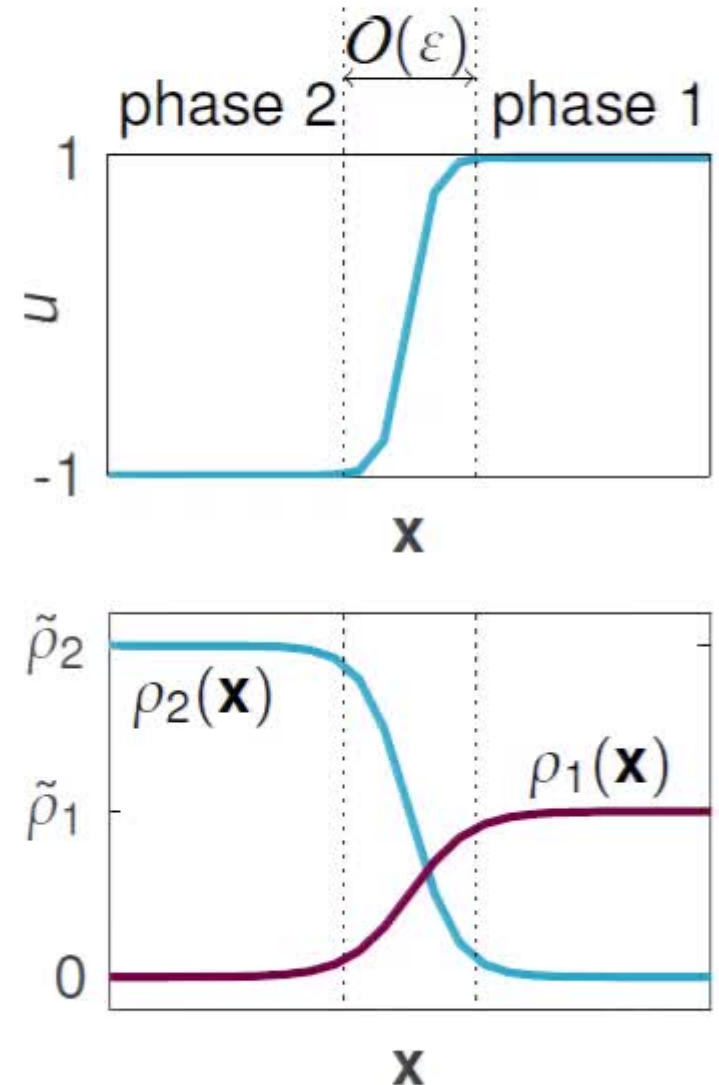
## Cahn–Hilliard Model

- continuous order-parameter  $u$  indicating the two phases

$$u(\mathbf{x}) = \frac{\rho_1(\mathbf{x})}{\tilde{\rho}_1} - \frac{\rho_2(\mathbf{x})}{\tilde{\rho}_2}$$

$$u \in \begin{cases} \{-1\} & \text{if } \mathbf{x} \text{ in phase 2} \\ (-1, 1) & \text{if } \mathbf{x} \text{ in interface} \\ \{1\} & \text{if } \mathbf{x} \text{ in phase 1} \end{cases}$$

- interface of small thickness  $O(\varepsilon)$
- mixing inside the interface





## Energy Functional

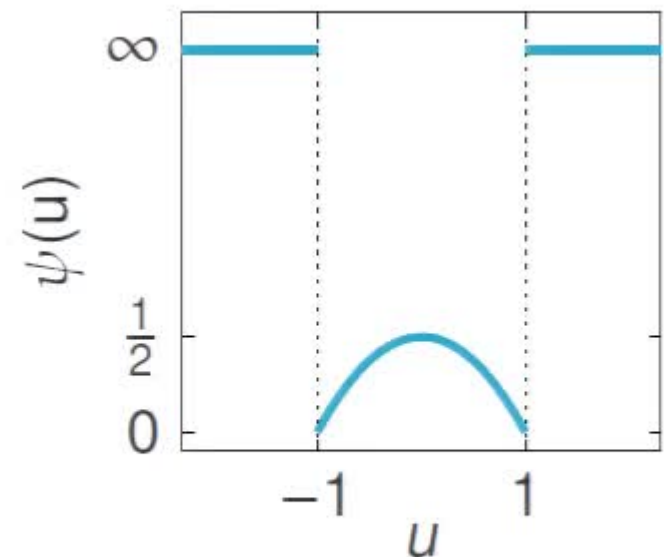
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- the evolution of  $u$  is driven by the minimization of

$$\mathcal{E}(u) = \int_{\Omega} \left[ \frac{\sigma \varepsilon}{2} |\nabla u|^2 + \frac{\sigma}{\varepsilon} \psi(u) \right] d\mathbf{x}$$

- double-obstacle potential

$$\psi(u) = \begin{cases} \frac{1}{2}(1 - u^2) & \text{if } u \in [-1, 1] \\ \infty & \text{else} \end{cases}$$







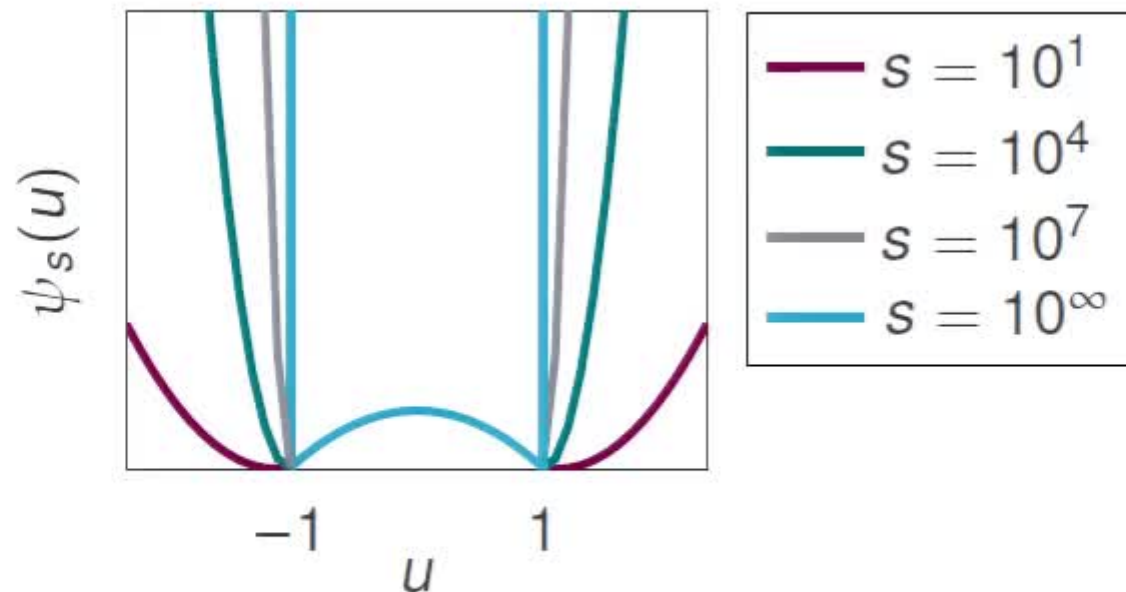
## Regularized Potential

[HINTERMÜLLER/HINZE/TBER '11]

Relaxation by Moreau–Yosida regularization:

$$\lambda(u) = \max(0, u - 1) + \min(0, u + 1)$$

$$\psi_s(u) = \frac{1}{2}(1 - u^2 + s\lambda^2(u)) \quad (s \rightarrow \infty)$$





## Cahn–Hilliard Part

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### Convective Cahn–Hilliard equations

$$\begin{cases} \mathbf{v} \cdot \nabla u + \partial_t u = \operatorname{div}(m \nabla w) \\ w = -\sigma \varepsilon \Delta u + \sigma \varepsilon^{-1} \psi'_s(u) \end{cases}$$

- $w$  chemical potential
- $m(u)$  mobility
- $\sigma$  surface tension
- $\varepsilon$  interfacial width
- $\mathbf{v}$  velocity

### Boundary conditions

- $\nabla w \cdot \mathbf{n} = 0 \rightsquigarrow$  no diffusion through boundary
- $\nabla u \cdot \mathbf{n} = 0 \rightsquigarrow$  interface is orthogonal to the boundary



## Navier–Stokes Part

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Navier–Stokes equation + surface tension force

$$\begin{cases} \rho \partial_t \mathbf{v} + ((\rho \mathbf{v} + \mathbf{J}) \cdot \nabla) \mathbf{v} - \operatorname{div}(2\eta D\mathbf{v}) + \nabla p = \rho \mathbf{g} + w \nabla u \\ \operatorname{div} \mathbf{v} = 0 \end{cases}$$

■  $\rho(u)$  density

■  $\eta(u)$  viscosity

■  $\mathbf{g}$  gravitation

■  $2D\mathbf{v} = \nabla \mathbf{v} + (\nabla \mathbf{v})^T$  strain tensor

$$\text{■ } \mathbf{J} = -\frac{\tilde{\rho}_2 - \tilde{\rho}_1}{2} m(u) \nabla w$$

## Boundary conditions

■  $\mathbf{v} = \mathbf{f}$  with  $\mathbf{f} \cdot \mathbf{n} = 0$



## Coupled System

[ABELS/GARCKE/GRÜN '12]

## Cahn–Hilliard/Navier–Stokes system

$$\left\{ \begin{array}{l} \rho \partial_t \mathbf{v} + ((\rho \mathbf{v} + \mathbf{J}) \cdot \nabla) \mathbf{v} - \operatorname{div}(2\eta D\mathbf{v}) + \nabla p = \rho \mathbf{g} + w \nabla u \\ \operatorname{div} \mathbf{v} = 0 \\ \mathbf{v} \cdot \nabla u + \partial_t u = \operatorname{div}(m \nabla w) \\ w = -\sigma \varepsilon \Delta u + \sigma \varepsilon^{-1} \psi'_s(u) \end{array} \right.$$

- $w \nabla u$  capillary force
- $\mathbf{v} \cdot \nabla u$  convection term
- $\mathbf{J} = -\frac{\tilde{\rho}_2 - \tilde{\rho}_1}{2} m(u) \nabla w$  crucial for thermodynamical consistency



- thermodynamical consistent time discrete scheme
- adaptive spatial discretization  
     $\rightsquigarrow$  thermodynamical consistency is conserved
- Taylor-Hood LBB-stable  $P2 - P1$  finite elements for the velocity-pressure field
- $P1$  finite elements for the phase field and chemical potential



## Fully Coupled CH/NS System

$$\mathcal{A}\mathbf{x} = \left( \begin{array}{ccc|cc} F_{11} & F_{12} & B_1^T & 0 & I_1 \\ F_{21} & F_{22} & B_2^T & 0 & I_2 \\ B_1 & B_2 & 0 & 0 & 0 \\ \hline T_1 & T_2 & 0 & C_{11} & C_{12} \\ 0 & 0 & 0 & C_{21} & C_{22} \end{array} \right) \begin{pmatrix} \delta v_1 \\ \delta v_2 \\ \delta p \\ \delta u \\ \delta w \end{pmatrix} = \begin{pmatrix} r_{v_1} \\ r_{v_2} \\ r_p \\ r_u \\ r_w \end{pmatrix}$$

- $F$  discrete convection-diffusion operator
- $B$  negative discrete divergence operator
- $I$  interfacial force coupling
- $T$  interfacial transport coupling
- $C_{12}$  mobility dependent diffusion operator
- $C_{21}$  diffusion operator + regularization



## Optimal Preconditioning

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Non-singular (1, 1) block (as in Navier–Stokes)

$$\mathcal{A} = \begin{pmatrix} F & B^T \\ B & 0 \end{pmatrix} \quad \mathcal{P} = \begin{pmatrix} F & B^T \\ 0 & BF^{-1}B^T \end{pmatrix}$$

[MURPHY/GOLUB/WATHEN '00] showed exactly two eigenvalues  $\pm 1$ .

Non-singular (1, 1) block

$$\mathcal{A} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \quad \mathcal{P} = \begin{pmatrix} A & B \\ 0 & D - CA^{-1}B^T \end{pmatrix}$$

[IPSEN '01] showed exactly two eigenvalues  $\pm 1$ .



## Navier–Stokes Subproblem

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$$\mathcal{A}_{NS} = \begin{pmatrix} F & B^T \\ B & 0 \end{pmatrix}$$

- $F$  discrete convection-diffusion operator
- $B$  negative discrete divergence operator

[ELMAN/SILVESTER/WATHEN '05] showed

$$\mathcal{P}_{NS} = \begin{pmatrix} F & B^T \\ 0 & \hat{S}_{NS} \end{pmatrix} \quad \hat{S}_{NS} = A_p F_p^{-1} Q_p$$

as a good approximation to the Schur complement preconditioner.

- $A_p$  pressure space Laplacian
- $F_p$  pressure space convection-diffusion operator
- $Q_p$  pressure space mass matrix





## Cahn–Hilliard Subproblem

$$\mathcal{A}_{CH} = \begin{pmatrix} M & -\varepsilon K - sGMG + \frac{1}{\varepsilon}M \\ \tau K & M \end{pmatrix} = \begin{pmatrix} M & -L \\ \tau K & M \end{pmatrix}$$

- $M$  mass matrix
- $K$  stiffness matrix
- $G$  diagonal penalization matrix
- $\tau$  time step

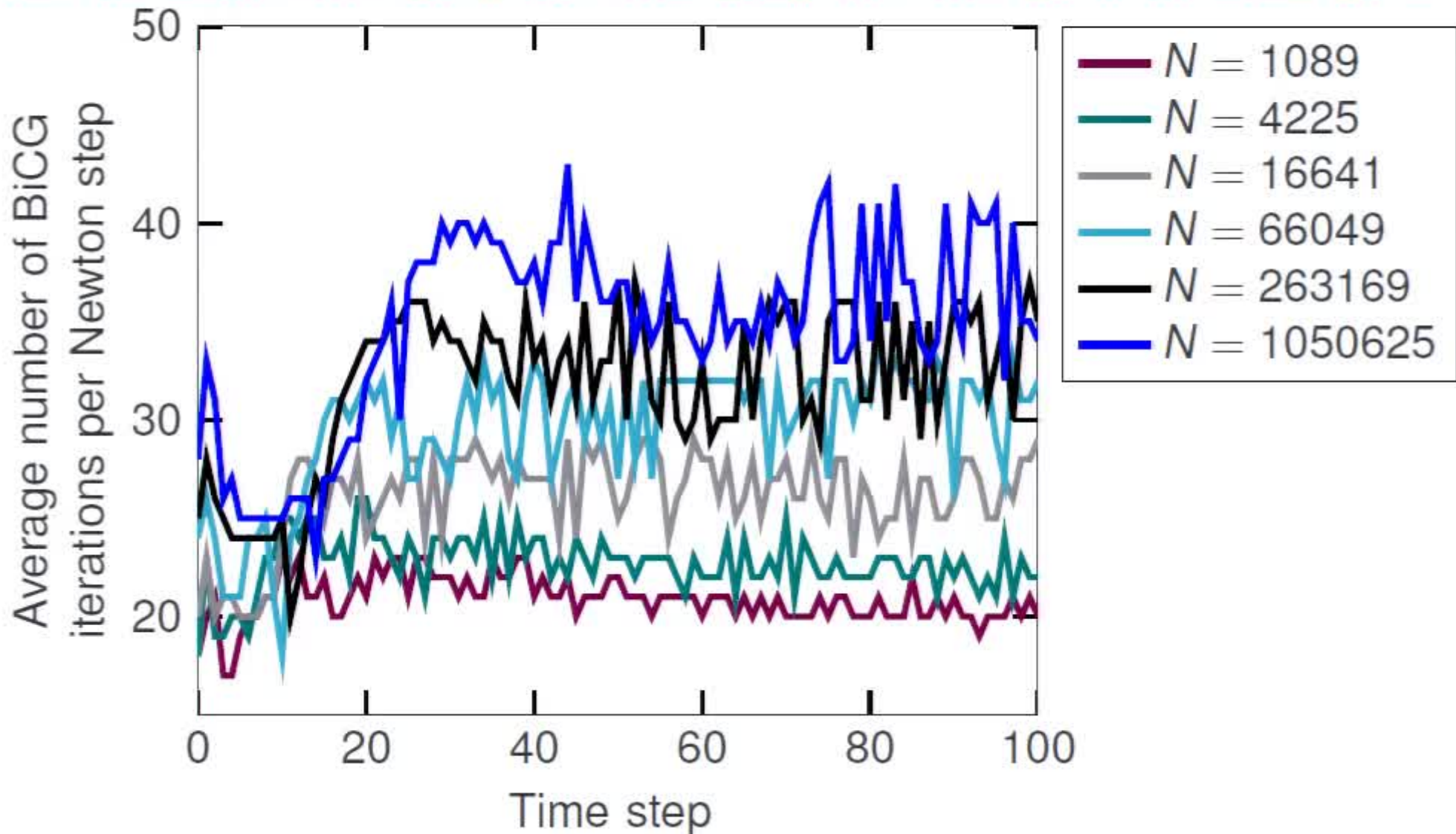
[BOSCH/STOLL/BENNER '14] showed

$$\mathcal{P}_{NS} = \begin{pmatrix} M & 0 \\ \tau K & -\hat{S}_{CH} \end{pmatrix} \quad \hat{S}_{CH} = (M + \sqrt{\tau}K)M^{-1}(M + \sqrt{\tau}L)$$

as a good approximation to the Schur complement preconditioner.



## BiCG Iteration Numbers for Different Uniform Meshes





## CH/NS Preconditioner

$$\mathcal{A} = \left( \begin{array}{ccc|cc} F_{11} & F_{12} & B_1^T & 0 & I_1 \\ F_{21} & F_{22} & B_2^T & 0 & I_2 \\ B_1 & B_2 & 0 & 0 & 0 \\ \hline T_1 & T_2 & 0 & C_{11} & C_{12} \\ 0 & 0 & 0 & C_{21} & C_{22} \end{array} \right)$$

$$\mathcal{P} = \left( \begin{array}{ccc|cc} F_{11} & 0 & B_1^T & 0 & I_1 \\ 0 & F_{22} & B_2^T & 0 & I_2 \\ 0 & 0 & -\hat{S}_{NS} & 0 & 0 \\ \hline 0 & 0 & 0 & C_{21} & C_{22} \\ 0 & 0 & 0 & 0 & -\hat{S}_{CH} \end{array} \right)$$



## CH/NS Preconditioner

$$\mathcal{P} = \left( \begin{array}{ccc|cc} F_{11} & 0 & B_1^T & 0 & I_1 \\ 0 & F_{22} & B_2^T & 0 & I_2 \\ 0 & 0 & -\hat{S}_{NS} & 0 & 0 \\ \hline 0 & 0 & 0 & C_{21} & C_{22} \\ 0 & 0 & 0 & 0 & -\hat{S}_{CH} \end{array} \right)$$

$$S_{NS} = BF^{-1}B^T$$

$$S_{CH} = C_{12} - C_{11}C_{21}^{-1}C_{22}$$

$$\hat{S}_{NS} = A_p F_p^{-1} Q_p$$

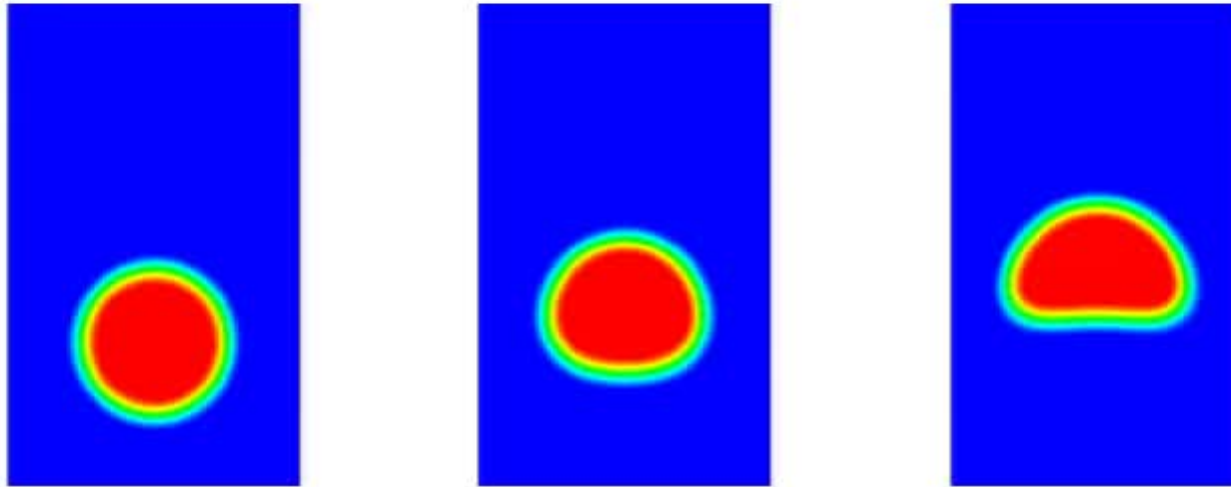
$$\hat{S}_{CH} = S_A C_{21}^{-1} S_B$$

$$= (\tau^{-1} C_{12} + C_{11}) C_{21}^{-1} (\tau C_{21} - C_{22})$$

- $F_{11}^{-1}, F_{22}^{-1}, A_p^{-1}, C_{21}^{-1}, S_A^{-1}, S_B^{-1}$  via AMG (library HSL\_MI20)
- $Q_p^{-1}$  via lumping



## Rising Bubble



We fix:

$$\tilde{\rho}_2 = 100, \tilde{\eta}_1 = 10, \tilde{\eta}_2 = 1, s = 10^6, \sigma = 15.6$$

We vary:

$$h, \tilde{\rho}_1 \rightsquigarrow \text{Re}, \epsilon, \tau, m$$



## Rising Bubble - Vary Mesh

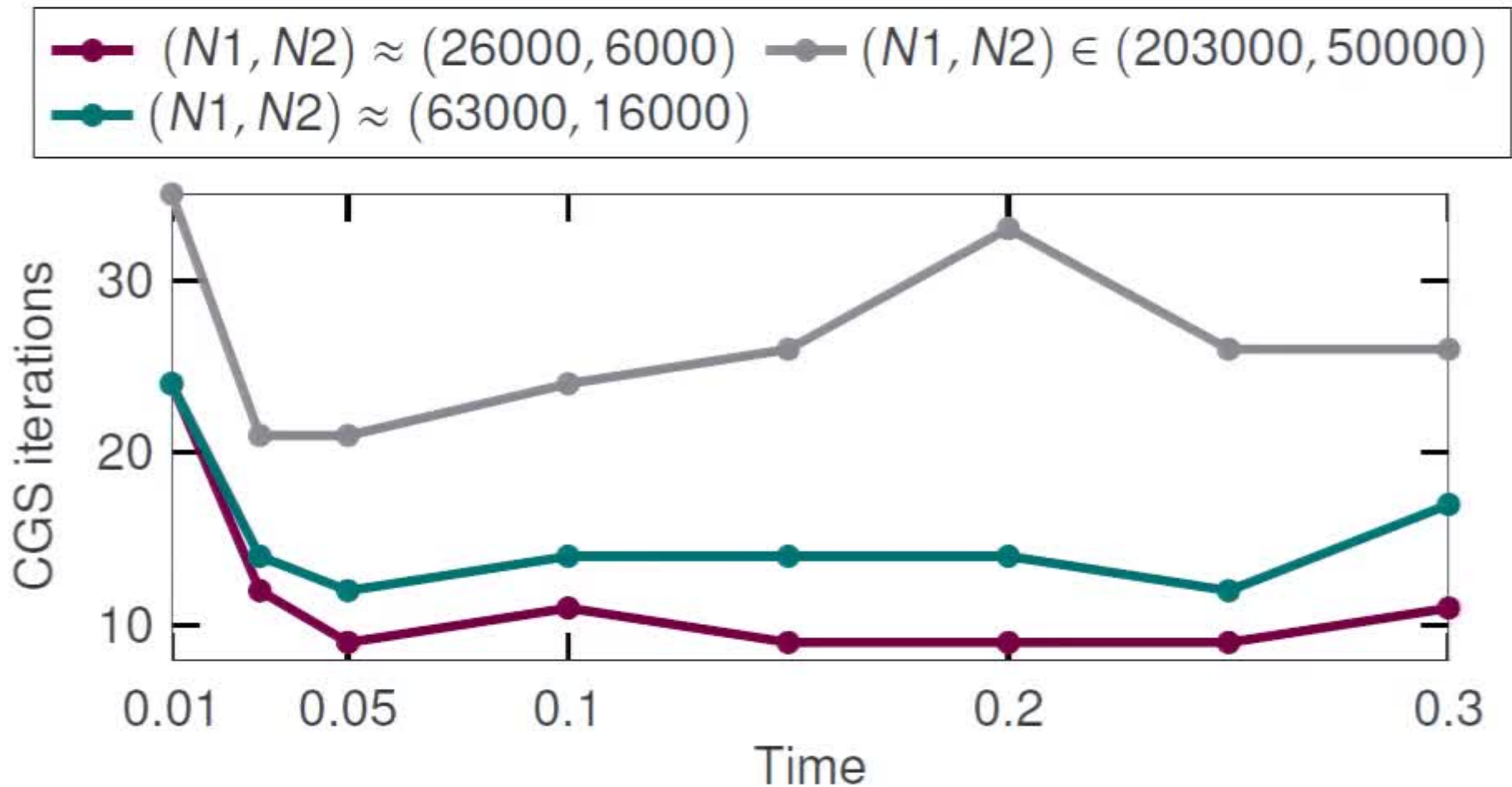
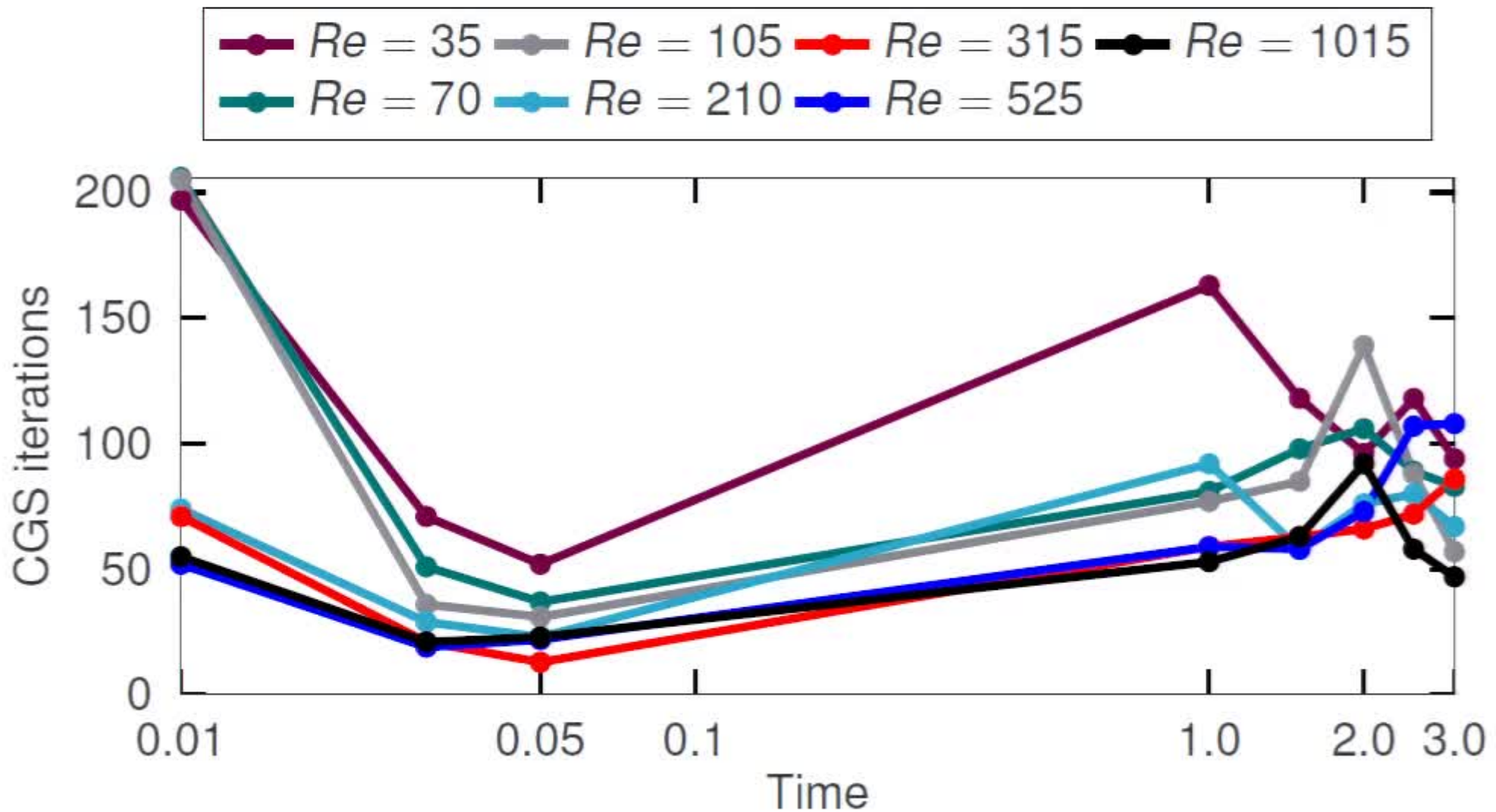


Figure : Vary the mesh, fix  $\varepsilon = 0.04$ ,  $\tau = 1.25 \cdot 10^{-4}$ ,  $m = 4 \cdot 10^{-5}$ ,  $\tilde{\rho}_1 = 1000$ .



## Rising Bubble - Vary Density / Reynolds Number





## Rising Bubble - Vary Mesh, Interface and Time Step Simultaneously

●  $(N1, N2) \approx (3.0 \cdot 10^4, 7.0 \cdot 10^3)$  ●  $(N1, N2) \approx (8.0 \cdot 10^4, 2.0 \cdot 10^4)$   
●  $(N1, N2) \approx (5.5 \cdot 10^4, 1.2 \cdot 10^4)$  ●  $(N1, N2) \approx (1.4 \cdot 10^5, 3.5 \cdot 10^4)$

